

Dimensional BCS-BEC crossover

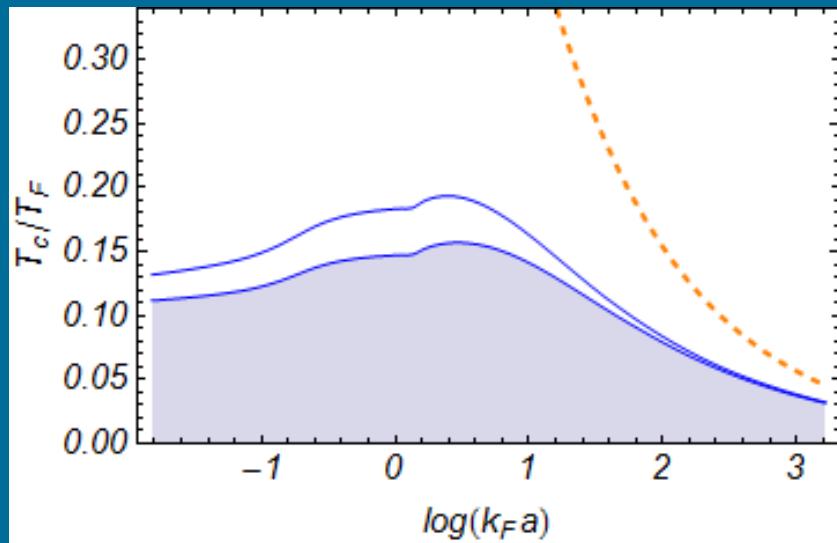
Igor Boettcher

Institute for Theoretical Physics,
Heidelberg University

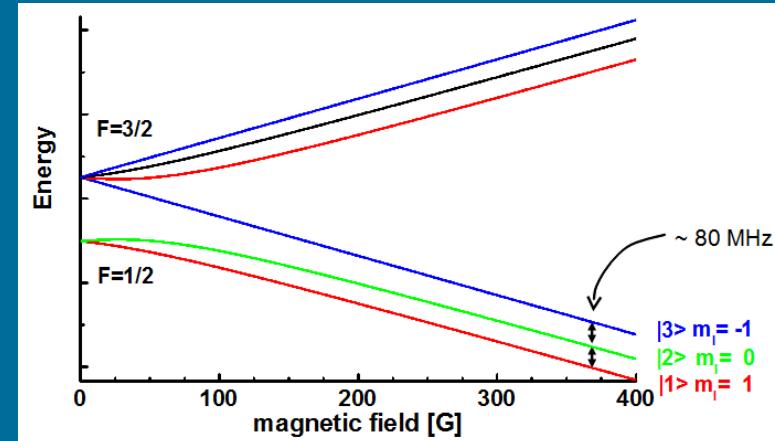
ERG 2014

Outline

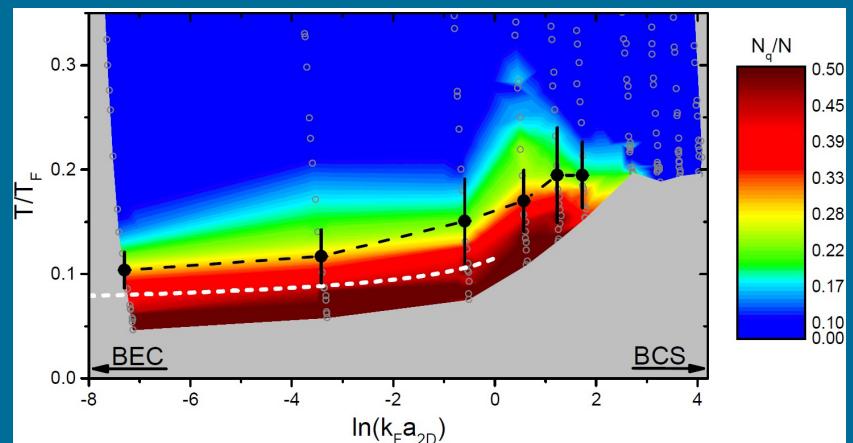
Ultracold atoms



2D Experiments



BCS-BEC Crossover



with...

Theory

C. Wetterich
J. M. Pawlowski

T. K. Herbst
N. Strodthoff
J. Braun (Darmstadt)
D. Roscher (Darmstadt)
L. von Smekal (Giessen/Darmstadt)
T. Enss
M. Scherer
S. Lammers
F. Golibrzuch
S. Diehl (Dresden)
S. Floerchinger (CERN)

Experiment

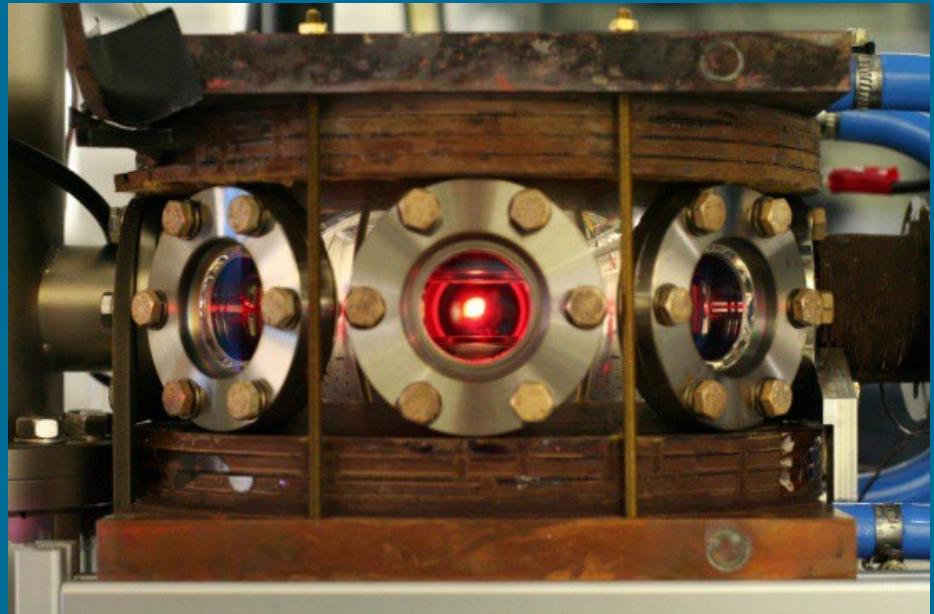
S. Jochim

M. Ries
A. Wenz
G. Zuern
P. Murthy
M. Neidig
T. Lompe
D. Kedar
L. Bayha

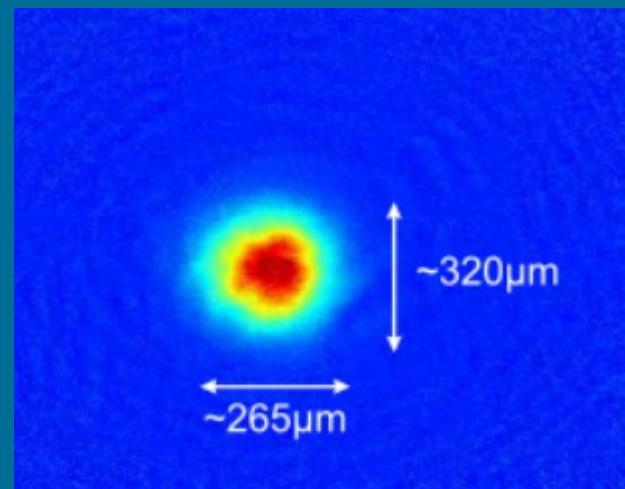
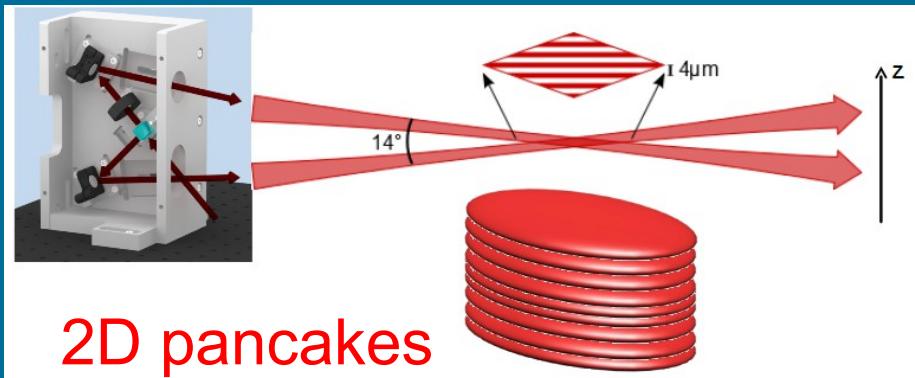
Many-body physics with cold atoms

Ultracold atoms:

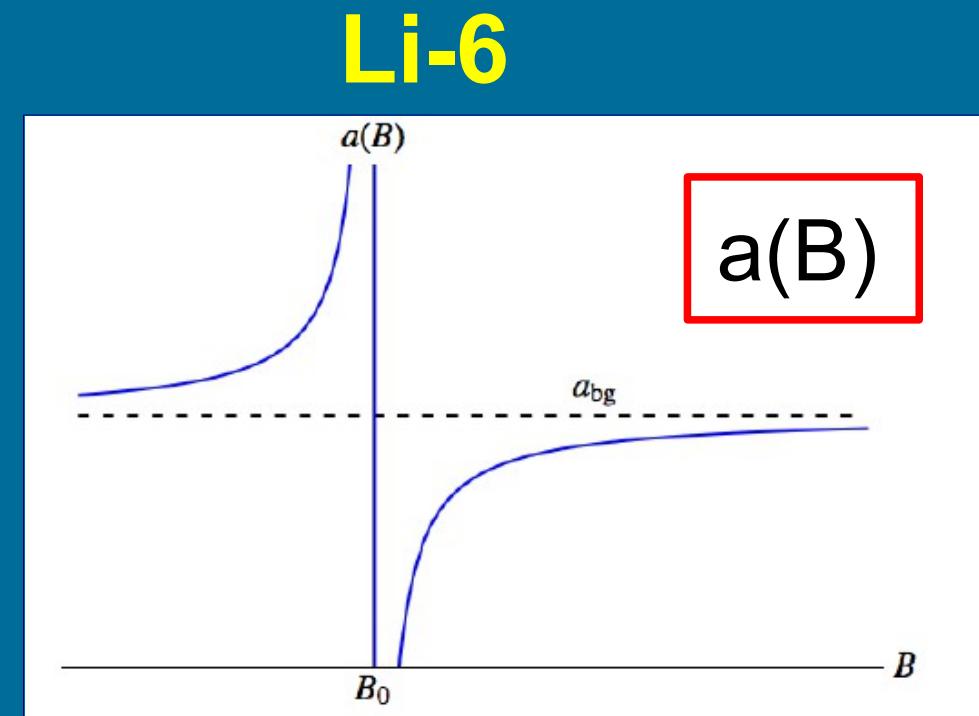
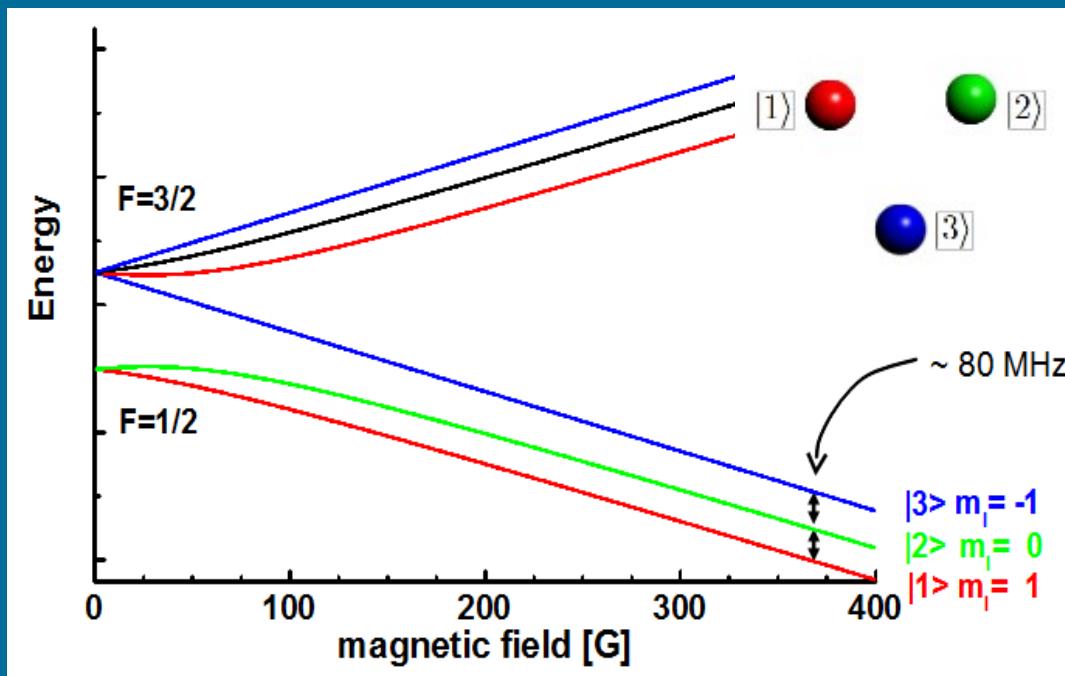
- 10^6 atoms trapped by light (Hz-kHz)
- $100 \text{ nK}, T/T_F = O(1)$



Clean realizations of many-body Hamiltonians.
High experimental control & accessibility.



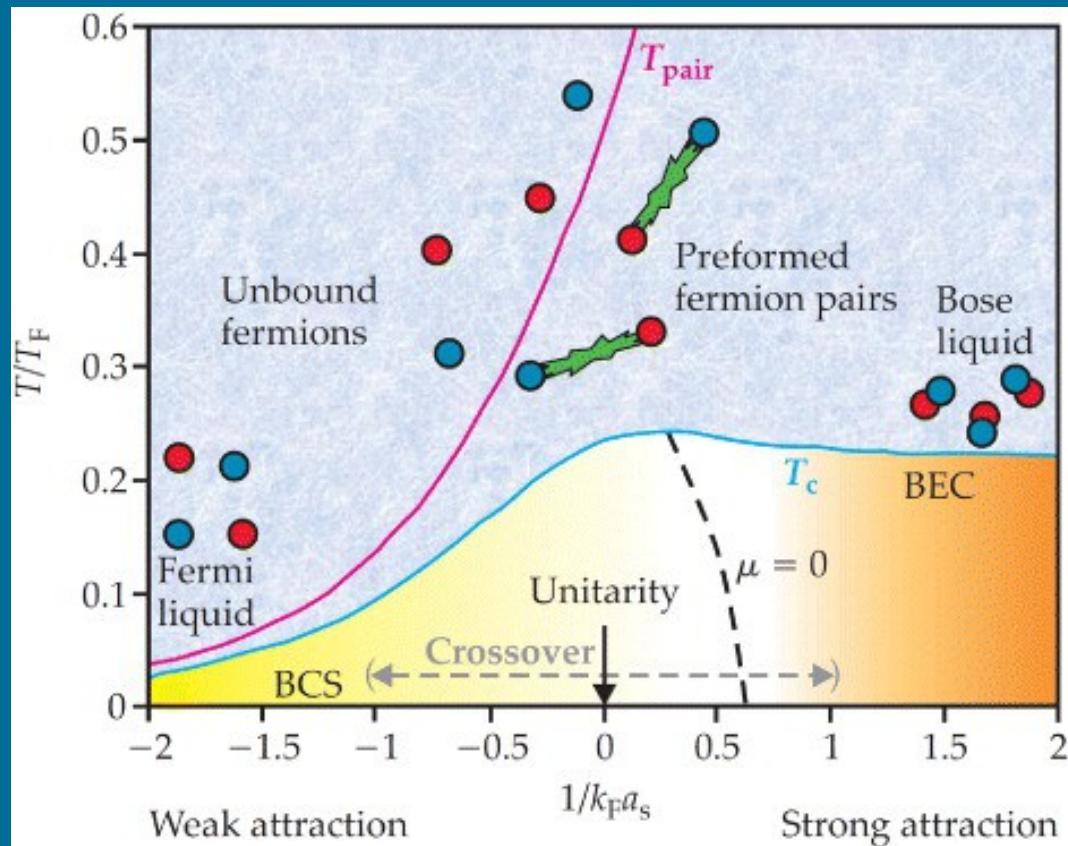
Many-body physics with cold atoms



Interactions between $|1\rangle$ and $|2\rangle$ atoms encoded in the s-wave scattering length $a=a(B)$

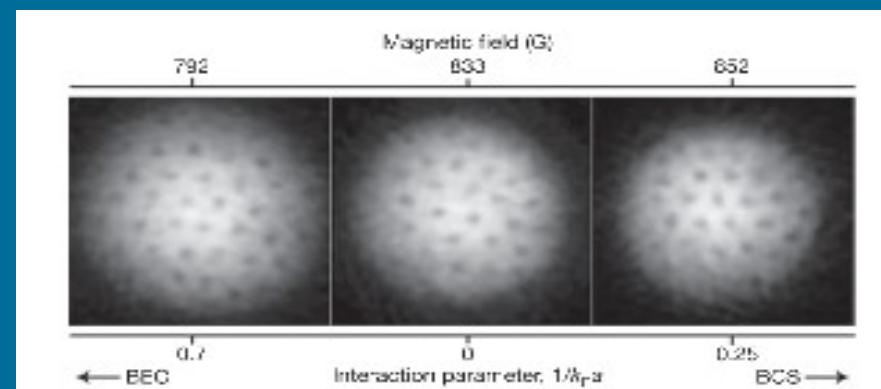
For small a : $\sigma = 4\pi a^2$

BCS-BEC Crossover



C.A.R. Sà de Melo, Physics Today 61(10), 45 (2008)

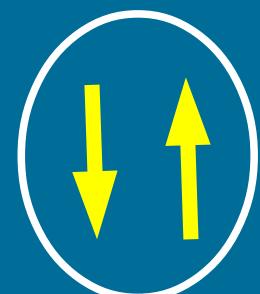
superfluidity: vortices



Zwierlein et al., Nature 435 , 1047-1051 (2005)

$$\sigma = 4\pi a^2$$

$$(k_F a)^{-1}$$



BCS-BEC Crossover

Feshbach resonance physics described by two-channel model

$$S = \int_X \left[\psi_1^* \left(\partial_\tau - \nabla^2 - \mu_1 \right) \psi_1 + \psi_2^* \left(\partial_\tau - \nabla^2 - \mu_2 \right) \psi_2 \right. \\ \left. + \phi^* \left(\partial_\tau - \frac{1}{2} \nabla^2 + m_\phi^2 \right) \phi - h \left(\phi^* \psi_1 \psi_2 + \phi \psi_2^* \psi_1^* \right) \right]$$

- Yukawa-type model:
- Fermion propagator
 - Boson propagator
 - Effective potential

- Key observables:
- Phase diagram
 - Equation of state
 - Momentum distribution
- $$P(\mu, T, a) = \frac{k_B T}{V} \log Z_{k=0}$$

Functional Renormalization Group

FRG setup $\partial_k \Gamma_k = \frac{1}{2} S \text{Tr} \left(\frac{1}{\Gamma_k^{(2)} + R_k} \partial_k R_k \right)$

Effective potential: Taylor expansion or resolution on a grid

$$U_k(\rho) = m_k^2(\rho - \rho_{0k}) + \frac{\lambda_k}{2}(\rho - \rho_{0k})^2 + \dots$$

Boson propagator: scale-dependent derivative expansion

$$\Sigma_{\phi k}(P) = Z_k i p_0 + V_k p_0^2 + A_k p^2$$

Fermion propagator: iterative, frequency and momentum-resolved*

$$\Sigma_\psi(P) \simeq \frac{4C}{-i p_0 + p^2 - \mu} - \delta\mu$$

*IB, S. Diehl, J. M Pawłowski, C. Wetterich, PRA 87, 023606 (2013)

Functional Renormalization Group

q^2 -opt:

$$R_{\phi,k}(Q) = \left(k^2 - \frac{q^2}{2}\right) \theta\left(k^2 - \frac{q^2}{2}\right),$$

$$R_{\psi,k}(Q) = \left[\operatorname{sgn}(q^2 - \mu)k^2 - (q^2 - \mu)\right] \theta\left(k^2 - |q^2 - \mu|\right),$$

Q -exp:

$$R_{\phi,k}(Q) = \left(iq_0 + \frac{q^2}{2}\right) r\left(\frac{q_0^2 + q^4/4}{c_\phi k^4}\right),$$

$$R_{\psi,k}(Q) = \left(iq_0 + q^2 - \mu\right) r\left(\frac{q_0^2 + (q^2 - \mu)^2}{k^4}\right),$$

$$r(X) = \frac{1}{e^X - 1}$$

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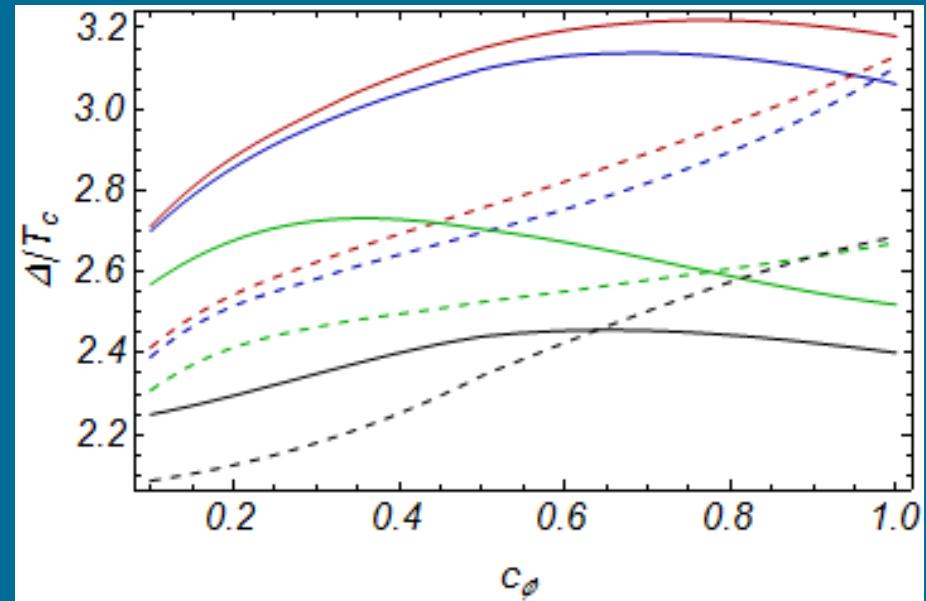
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Functional Renormalization Group

Critical temperature & superfluid gap of the Unitary Fermi Gas

	Δ/μ		T_c/μ	
Truncation	q^2 -opt	Q -exp	q^2 -opt	Q -exp
FB ₀ , ϕ^4	1.09	1.244(5)	0.441	0.399(2)
FB ₀ , ϕ^8	1.13	1.227	0.424	0.380
FB ₀ , ϕ^{10}	1.16	-	0.427	0.394
FB, ϕ^4	1.05	1.228(10)	0.405	0.389(2)
FB, ϕ^8	-	1.240	-	0.380
FB, ϕ^{10}	-	-	-	0.386



	T_c/μ				
Truncation	F	FB ₀	FB	FBM ₀	FBM
ϕ^4	0.664	0.381	0.381	0.385	0.383

$$\frac{T_c}{\mu} = 0.38(2), \quad \frac{\Delta}{\mu} = 1.04(15)$$

Dimensional BCS-BEC Crossover

$$\delta\mu = h = (\mu_1 - \mu_2)/2$$

d



Dimensional BCS-BEC Crossover

$$\delta\mu = h = (\mu_1 - \mu_2)/2$$

Phase structure of spin-imbalanced unitary Fermi gases
IB, J. Braun, T. K. Herbst, J. M. Pawłowski, D. Roscher,
C. Wetterich - arXiv:1409.5070

see talk by Dietrich Roscher today

Sarma phase in relativistic and non-relativistic systems
IB, T. K. Herbst, J. M. Pawłowski, N. Strodthoff,
L. von Smekal, C. Wetterich - arXiv:1409.5232

see talk by Tina K. Herbst today

Dimensional BCS-BEC Crossover

$$\delta\mu = h = (\mu_1 - \mu_2)/2$$

Observation of pair condensation in a strongly interacting two-dimensional Fermi gas

M. G. Ries, A. N. Wenz, G. Zuern, L Bayha, IB, D. Kedar, P. A. Murthy, M. Neidig, T. Lompe, S. Jochim - arXiv:1409.5373

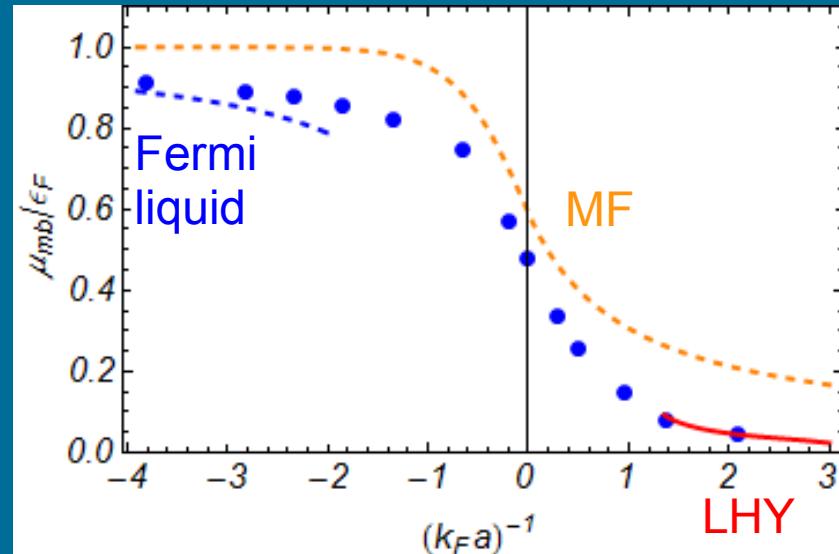
Dimensional BCS-BEC crossover in ultracold Fermi gases
IB, PHD thesis (2014)

IB, J. M. Pawłowski, C. Wetterich (in preparation, 2014)

d

Dimensional BCS-BEC Crossover

Equation of state at T=0



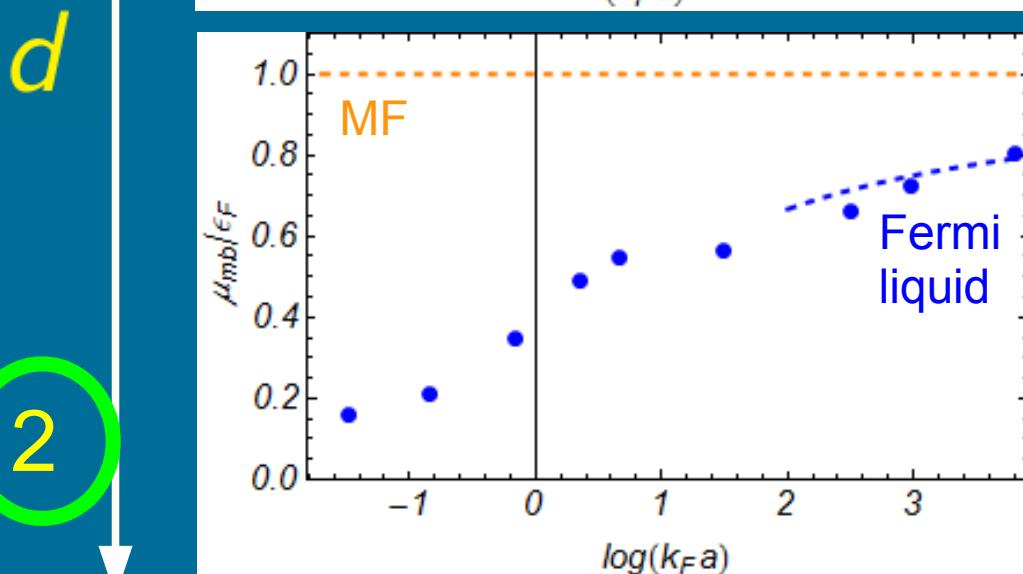
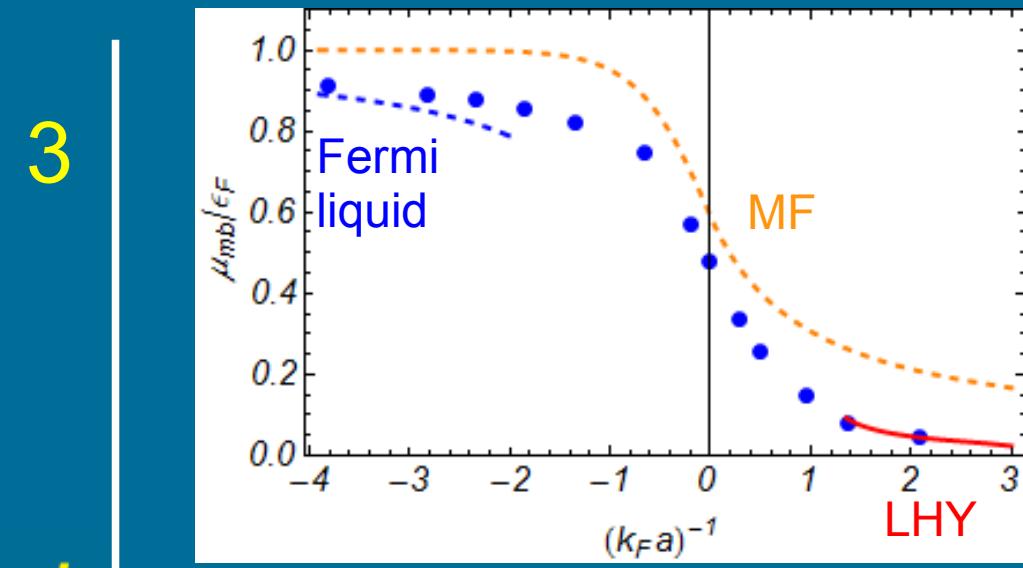
Crossover parameter 3D:

$$(k_F a)^{-1} \sim (\sqrt{\mu_{mb}} a)^{-1}$$

$$\mu_{mb} = \mu - \frac{\varepsilon_B}{2} \geq 0$$

Dimensional BCS-BEC Crossover

Equation of state at T=0



Crossover parameter 3D:

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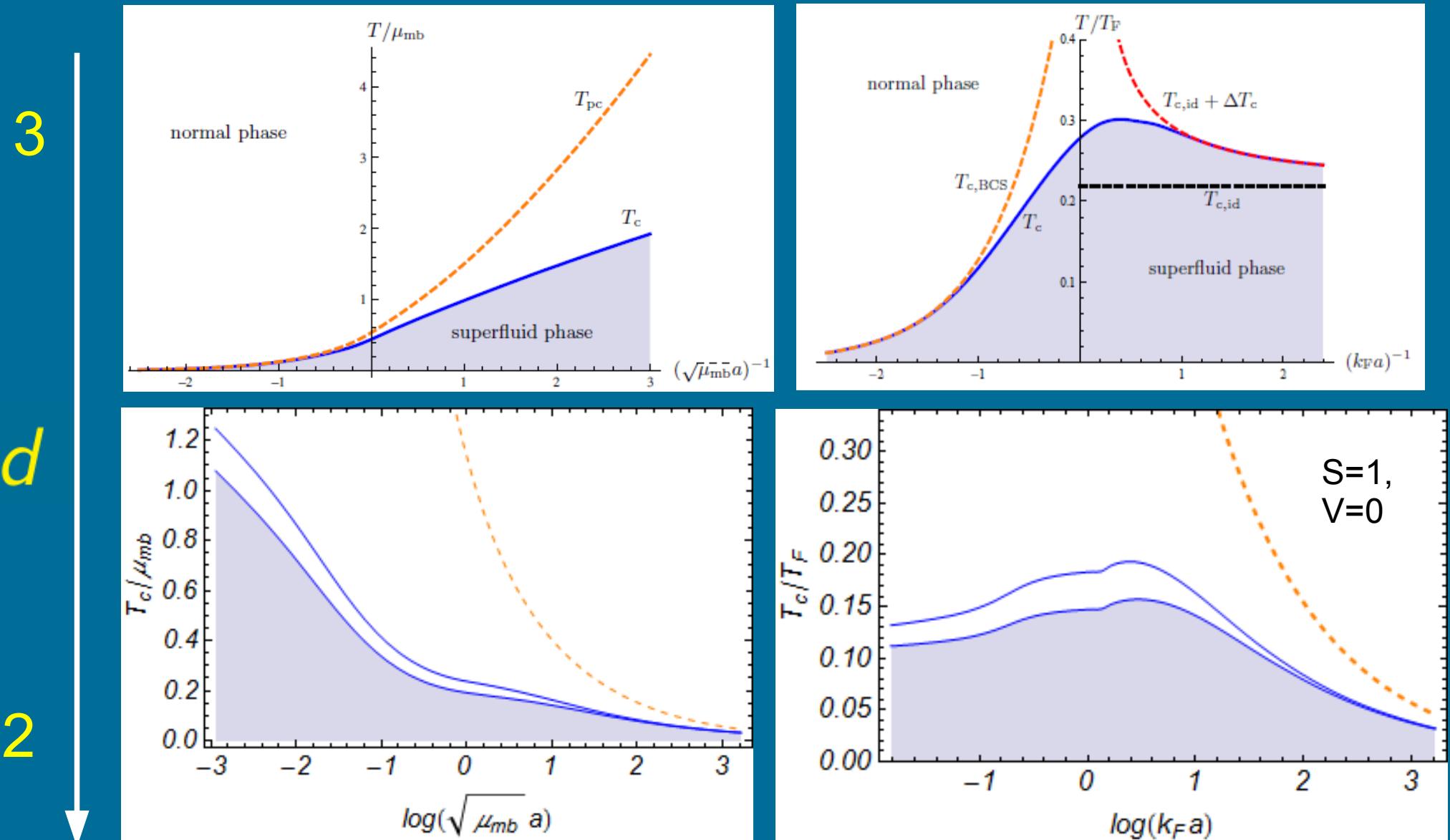
2D:

$$\log(k_F a) \sim \log(\sqrt{\mu_{mb}} a)$$

$$\mu_{mb} = \mu - \frac{\varepsilon_B}{2} \geq 0$$

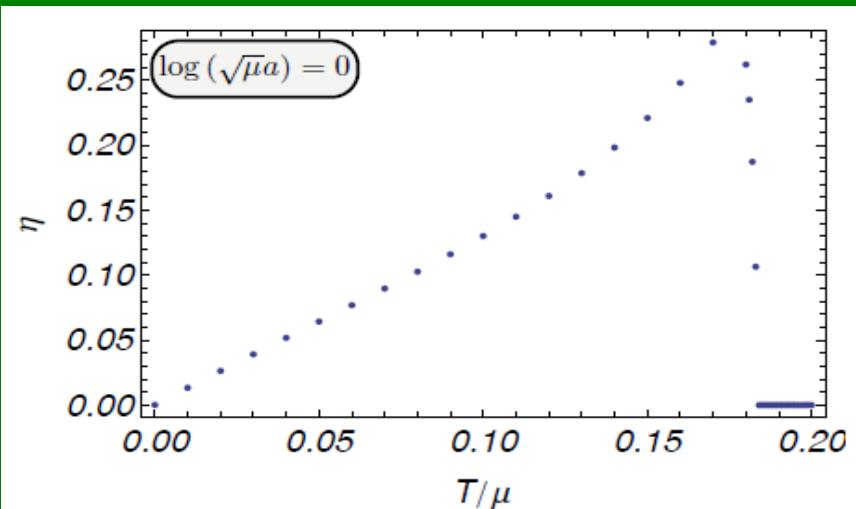
Dimensional BCS-BEC Crossover

Phase diagram

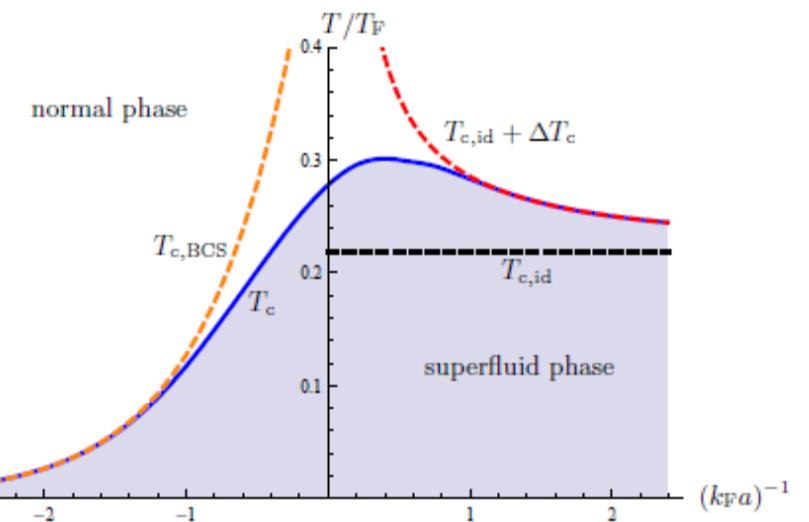


Dimensional BCS-BEC Crossover

BKT scaling is resolved from FRG
for all $\log(kFa)$:

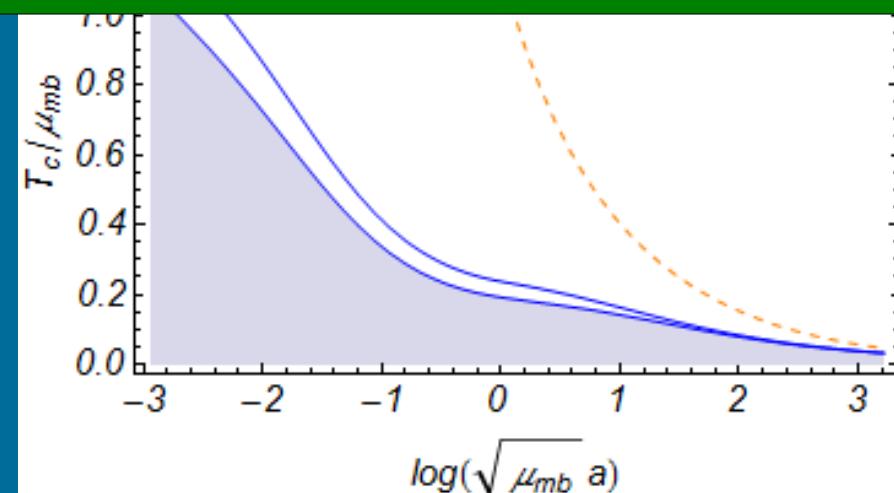


am

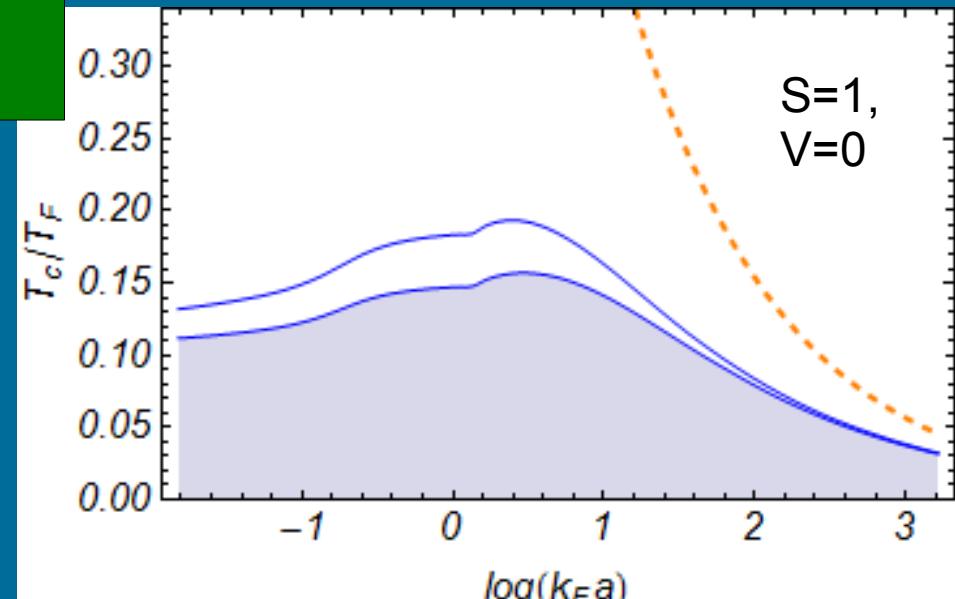


3

d

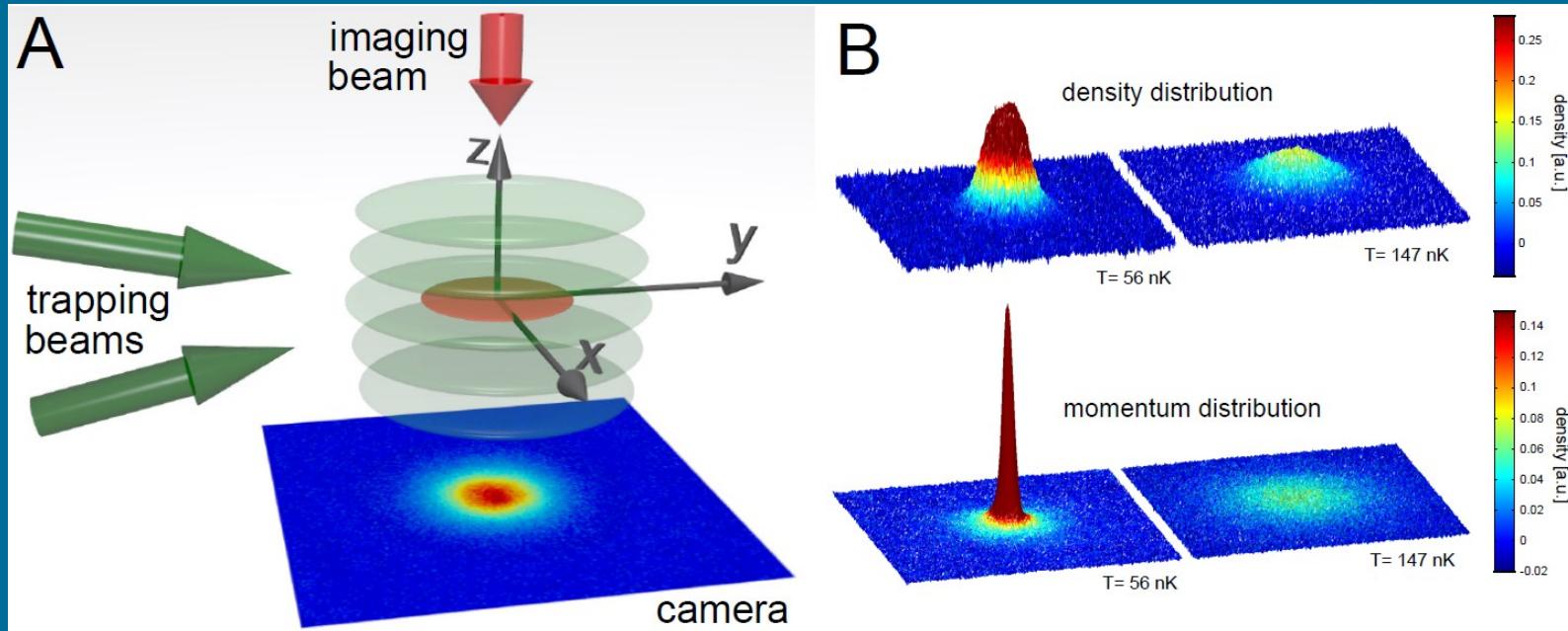


2



$S=1,$
 $V=0$

2D experiment

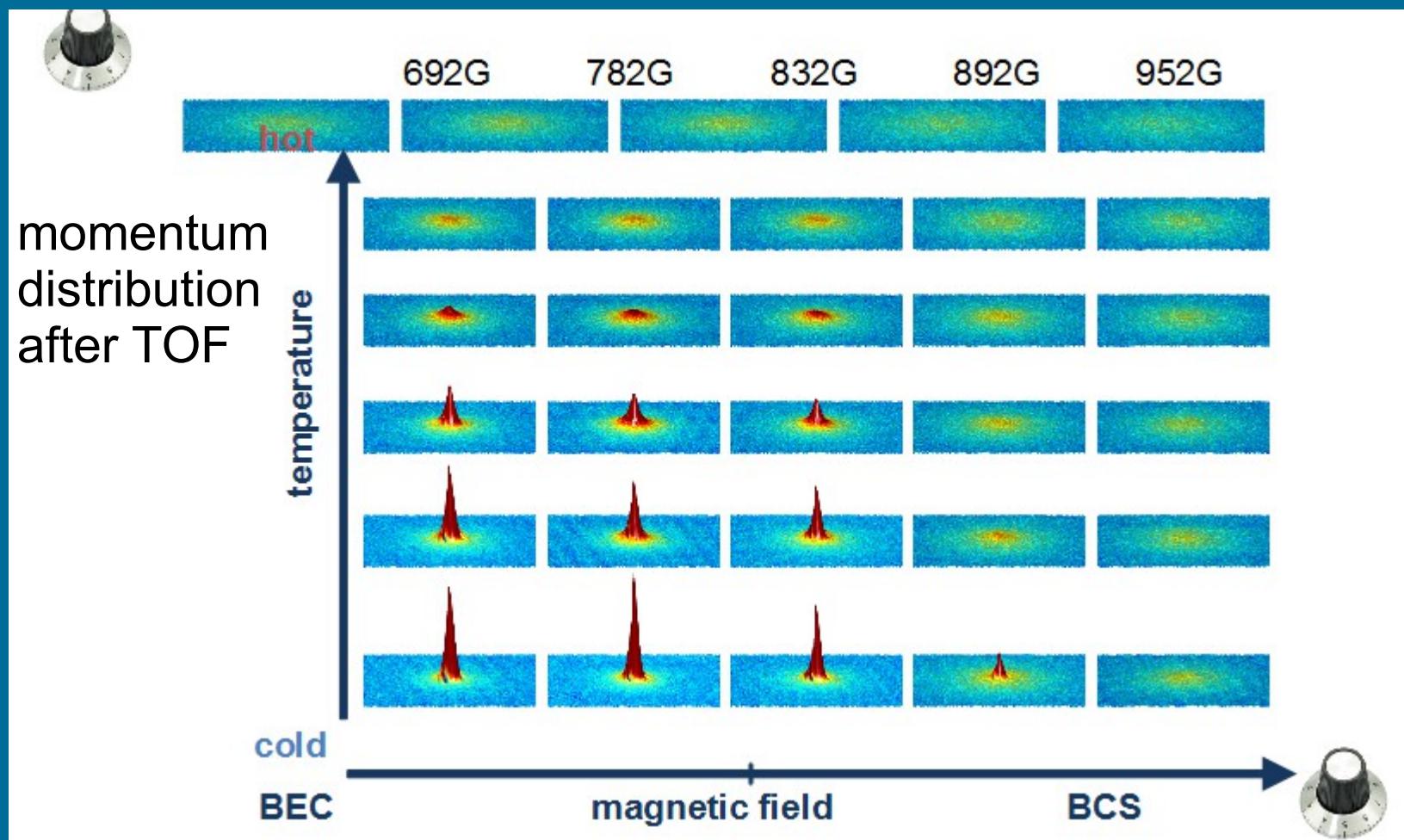


Jochim group @ PI Heidelberg

- **in-situ profiles** and **momentum distribution** after TOF expansion
- trap aspect ratio $\omega_x/\omega_z=1/300 \Rightarrow$ effectively 2D

2D experiment

Superfluid/condensation transition from enhanced occupation of low-momentum states

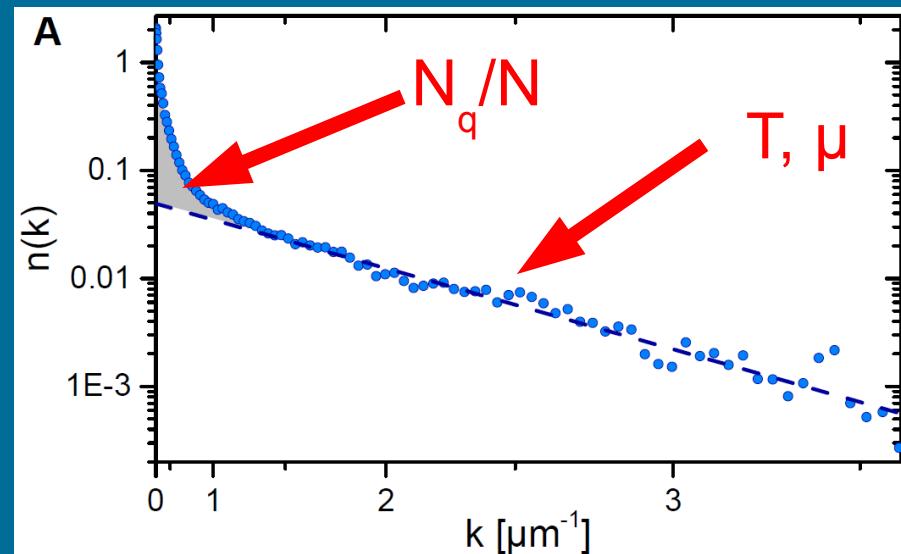


2D experiment

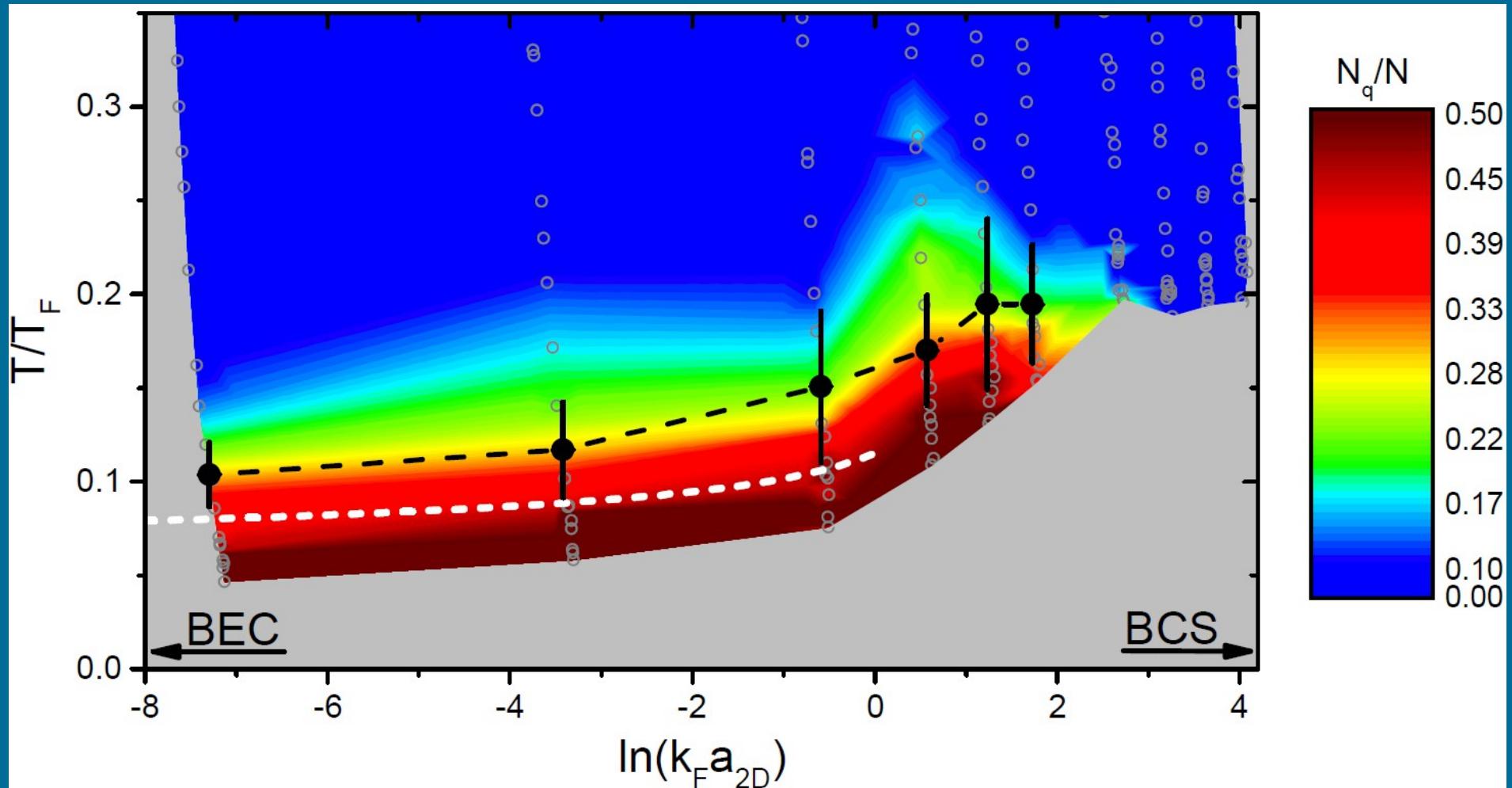
- 2D scattering length a from magnetic field B

$$a_{2D} = \ell_z \sqrt{\frac{\pi}{0.905}} \exp\left(-\sqrt{\frac{\pi}{2}} \frac{\ell_z}{a_{3D}(B)}\right)$$

- Boltzmann fit gives T and μ
- Density n from in-situ image
=> Full thermodynamic characterization $n(\mu, T, a)$!

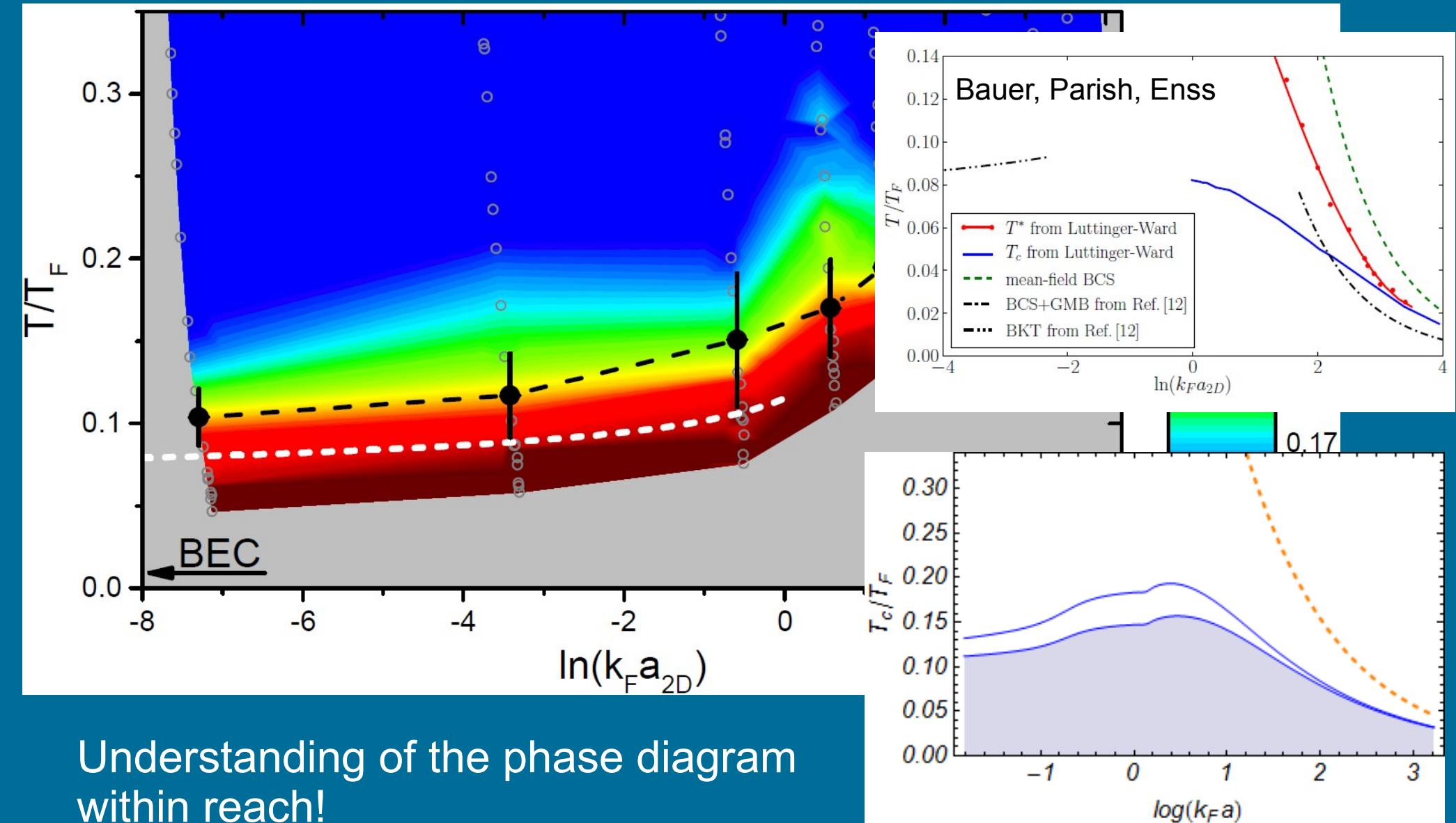


2D experiment



BKT verified from „infinite“ correlation length and algebraic decay of correlations below T_c

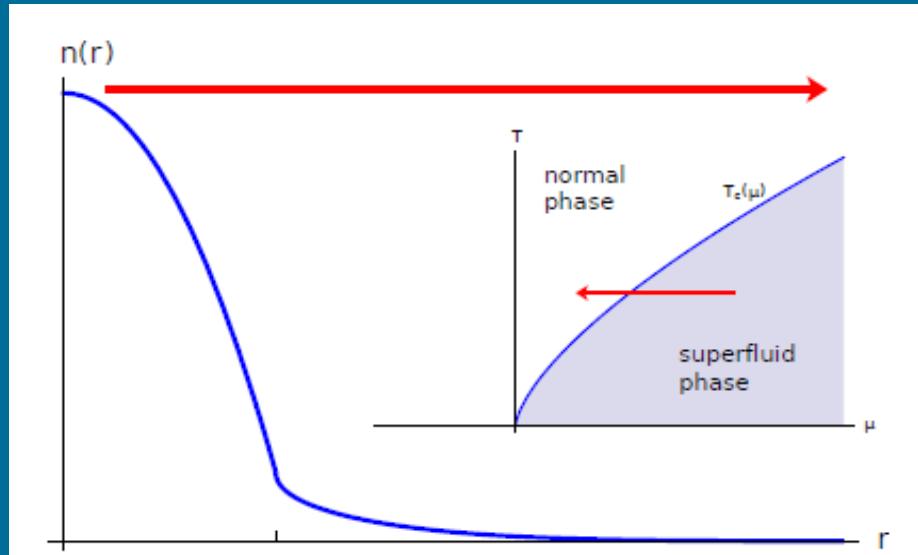
2D experiment



Equation of state determination

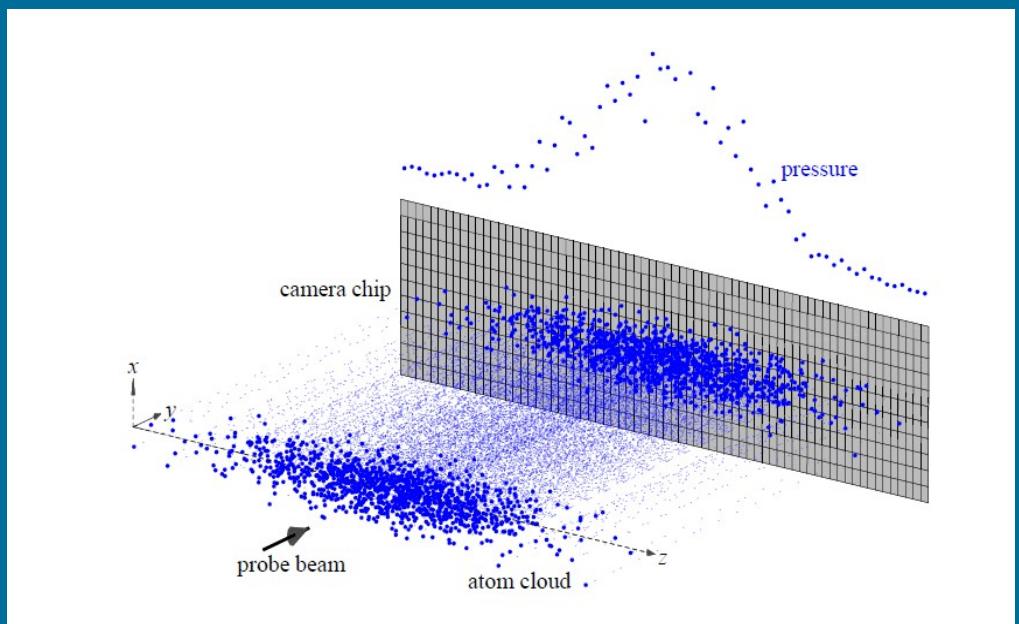
Equation of state $n(\mu, T, a)$
from in-situ images

$$P(\mu, T) \rightarrow P(\mu - V_{\text{ext}}(\vec{x}), T)$$



$$P(\mu_z, T) = \frac{m\omega_r^2}{2\pi} \bar{n}(z),$$

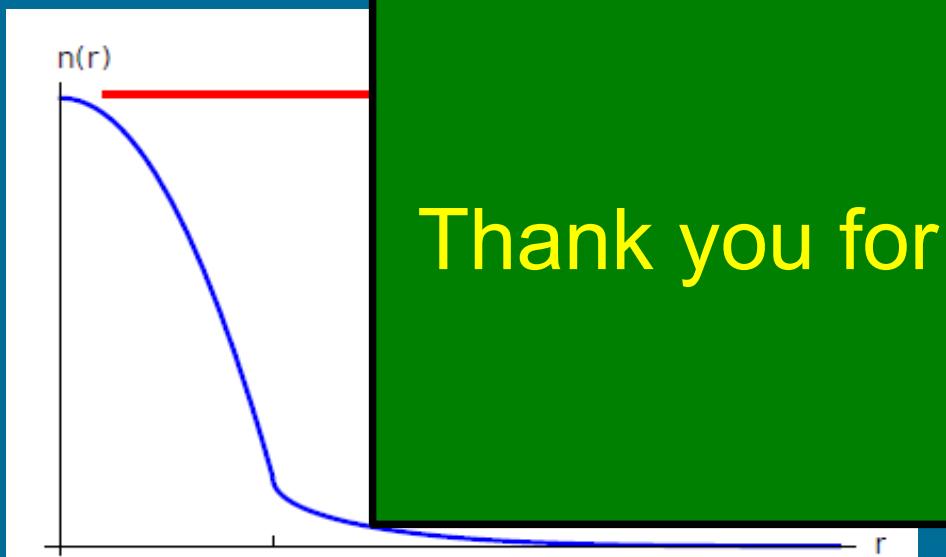
local density approximation



Equation of state determination

Equation of state $n(\mu, T, a)$
from in-situ images

$$P(\mu, T) \rightarrow P(\mu - V_{\text{ext}}(\vec{x}), T)$$



Thank you for your attention!

$$P(\mu_z, T) = \frac{m\omega_r^2}{2\pi} \bar{n}(z),$$

