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# **Long-range and anisotropic generalizations of the Kardar-Parisi-Zhang equation**

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## Experiment



Observations:

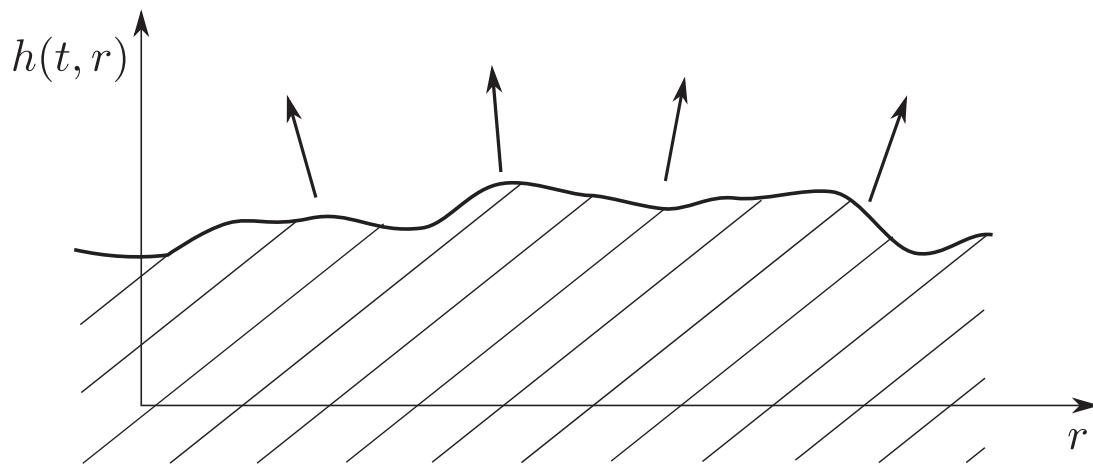
stochastic growth process

rough structure of the interface

scale invariance

Can we develop a theoretical understanding of random growth processes having the same features?

## Theoretical description



Dynamical scaling

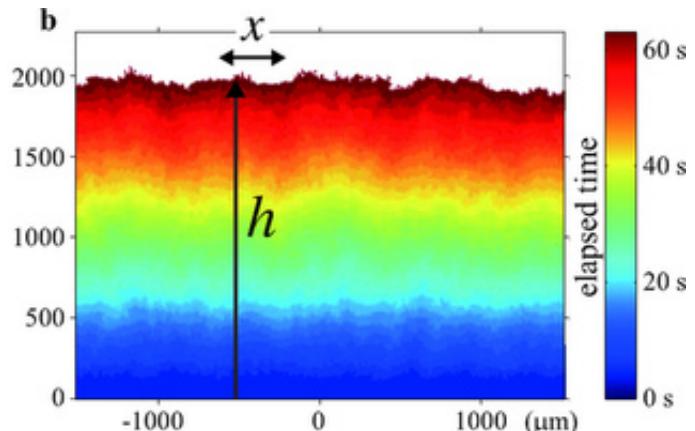
$$C(t, \vec{r}) = \langle h(t, \vec{r}) h(0, 0) \rangle_c \sim r^{2\chi} F(t/r^z)$$

## Model

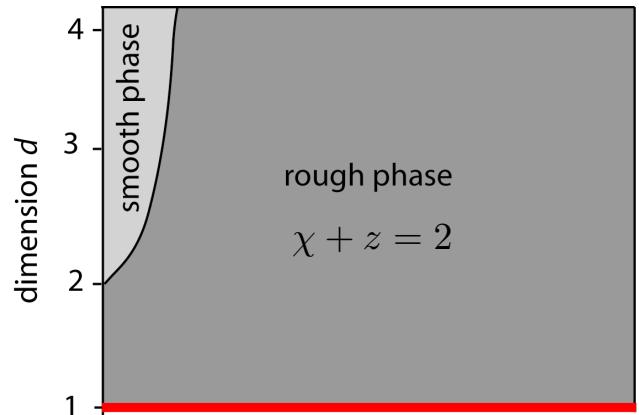
$$\frac{\partial h(t, \vec{x})}{\partial t} = \nu \nabla^2 h(t, \vec{x}) + \frac{\lambda}{2} (\nabla h(t, \vec{x}))^2 + \eta(t, \vec{x}) \quad \text{KPZ equation}$$

$$\langle \eta(t, \vec{x}) \eta(t', \vec{x}') \rangle = 2 D \delta^d(\vec{x} - \vec{x}') \delta(t - t')$$

Kardar, Parisi, Zhang, PRL (1986).



Takeuchi, Sano, Sasamoto, Spohn, Sci. Rep. (2011).



exactly solvable (since 2010)  
see e.g. Corwin, Random Matr. (2012).

## The upper critical dimension

$d_c$  finite

Mézard, Parisi, J Phys. I (1991):  $d_c = 2$

Le Doussal, Wiese, PRB (2003):  $d_c = 2.5$

Garel, Orland, PRB (1997):  $d_c = 2$

Colaiori, Moore, PRL (2001):  $d_c = 4$

Halpin-Healy, PRL (1989):  $d_c = 4$

Fogedby, PRL (2005):  $d_c = 4$

$d_c$  infinite

Marinari, Pagnani, Parisi, PRE (2013)

Tang, Forrest, Wolf, PRA (1992)

Ala-Nissila, et al., J Stat Phys (1992)

Castellano, et al., PRL (1998)

Oliveira, Alves, Ferreira., PRE (2013)



Study variants of the KPZ equation

- long-range correlated noise:  $d_c = 4$       Janssen, Täuber, Frey, Eur. Phys. J. B (1999)
- anisotropy:  $d_c \geq 3$                           Täuber, Frey, Europhys. Lett. (2002)

## Method

nonperturbative renormalization group (NPRG)

$$\partial_\kappa \Gamma_\kappa = \frac{1}{2} \text{Tr} \left[ \partial_\kappa R_\kappa (\Gamma_\kappa^{(2)} + R_\kappa)^{-1} \right]$$

Wetterich, Phys. Lett. B (1993).

ansatz

$$\Gamma_\kappa[\varphi, \tilde{\varphi}] = \int_{\mathbf{x}} \left\{ \tilde{\varphi} f_\kappa^\lambda(-\tilde{D}_t^2, -\nabla^2) D_t \varphi - \frac{1}{2} \left[ \nabla^2 \varphi f_\kappa^\nu(-\tilde{D}_t^2, -\nabla^2) \tilde{\varphi} + \tilde{\varphi} f_\kappa^\nu(-\tilde{D}_t^2, -\nabla^2) \nabla^2 \varphi \right] - \tilde{\varphi} f_\kappa^D(-\tilde{D}_t^2, -\nabla^2) \tilde{\varphi} \right\}$$

$$\tilde{D}_t = \partial_t - \lambda(\nabla \varphi) \cdot \nabla$$

form strongly constrained by symmetries: Galilei (gauged), shift symmetry (time gauged)

truncated at quadratic order in the response field

arbitrary momentum/frequency dependence

Blaizot, Méndez-Galain, Wschebor, Phys Lett B (2006)

$$R_\kappa(\mathbf{q}) = r \begin{pmatrix} \frac{q^2}{\kappa^2} \end{pmatrix} \begin{pmatrix} 0 & \nu_\kappa q^2 \\ \nu_\kappa q^2 & -2D_\kappa \end{pmatrix}$$

$$r(x) = \alpha / (\exp(x) - 1)$$

correlation and scaling functions

$$C(\varpi, p) = \langle h(\varpi, p) h(-\varpi, -p) \rangle_c \sim \frac{1}{p^{d+2+\chi}} \mathring{F}(\tau)$$

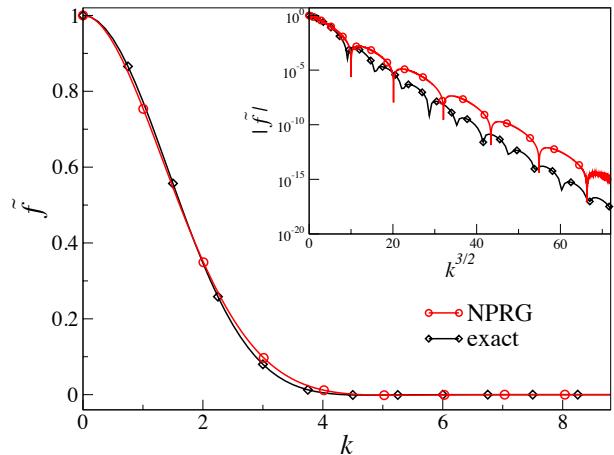
$$\mathring{F}(\tau) = \frac{\hat{\zeta}^D(\tau)}{(\tau \hat{\zeta}^\lambda(\tau))^2 + \hat{\zeta}^\nu(\tau)^2},$$

$$f_*^x(\varpi, p) \sim p^{-\eta_*^x} \zeta^x(\varpi/p^z)$$

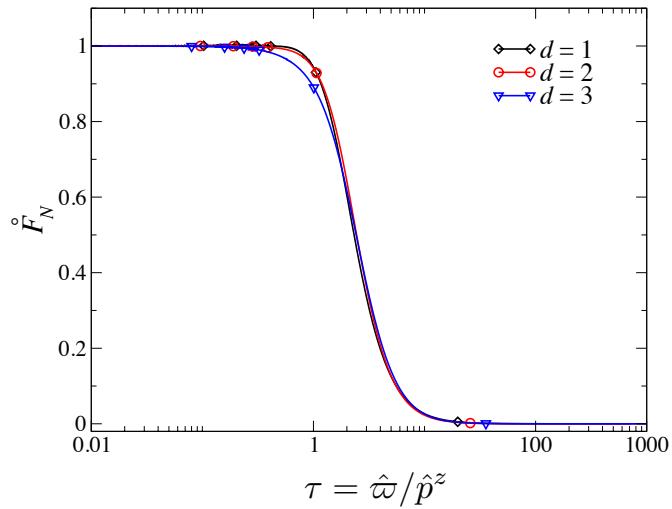
Canet, Delamotte, Wschebor PRE (2011).  
TK, Canet, Wschebor PRE (2012).

## Scaling functions and critical exponents

$d = 1$  vs. exact result



$d \geq 1$



exact result from Prähofer, Spohn J. Stat. Phys. (2004).

roughness exponent  $\chi$  for different dimensions  $d$  and comparison to literature results

$d$	1	2	3	4
$\chi$ (NPRG)	1/2	0.373(1)	0.179(4)	–
$\chi$ (Literature)*	1/2	0.379(15)	0.300(12)	0.246(7)

\*average over simulation and RG studies

Canet, Delamotte, Wschebor PRE (2011).  
TK, Canet, Wschebor PRE (2012).

## Long-range correlated noise

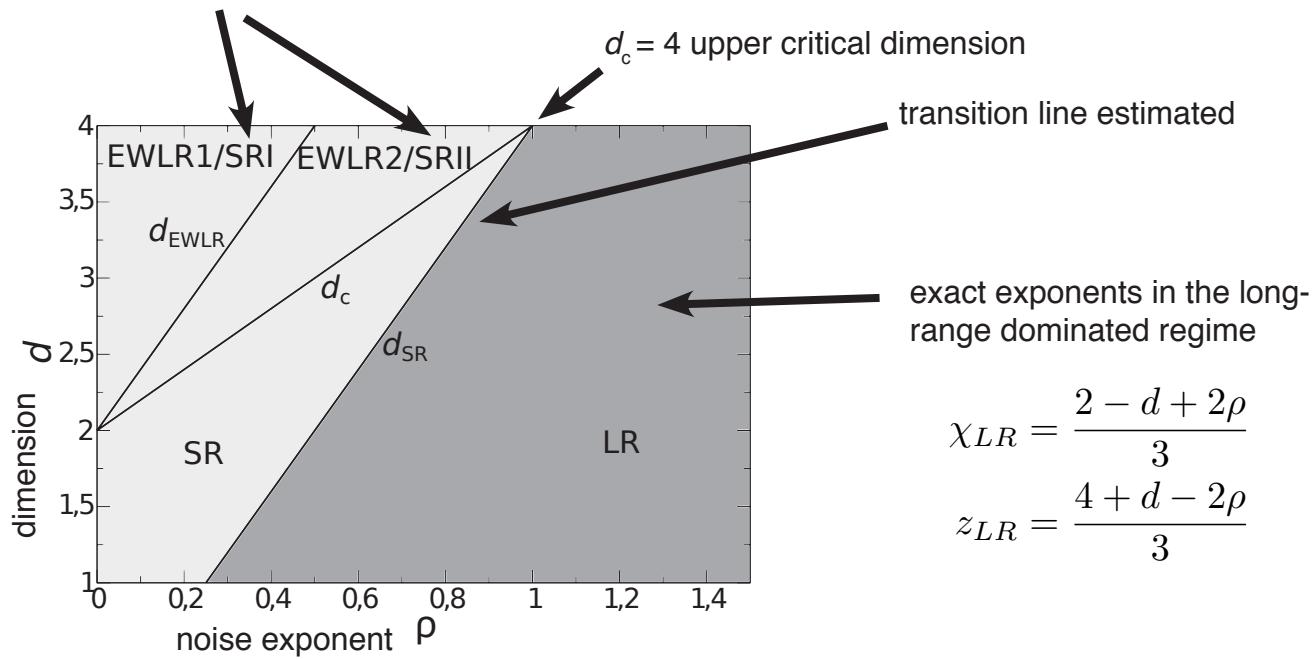
$$\langle \eta(t, \vec{x}) \eta(t', \vec{x}') \rangle = 2 R^D(\vec{x} - \vec{x}') \delta(t - t')$$

$$R^D(\vec{x} - \vec{x}') \sim |\vec{x} - \vec{x}'|^{2\rho-d}, \quad \rho \leq d/2 \quad \text{spacially correlated noise}$$

Weak-coupling phase diagram

Janssen, Täuber, Frey, Eur. Phys. J. B (1999).

different short range phases?



$$\chi_{LR} = \frac{2 - d + 2\rho}{3}$$

$$z_{LR} = \frac{4 + d - 2\rho}{3}$$

## Long-range correlated noise (NPRG)

Approximations

$$\hat{f}_\kappa^x(\omega, q) \rightarrow \hat{f}_\kappa^x(q) \quad (\text{NLO})$$

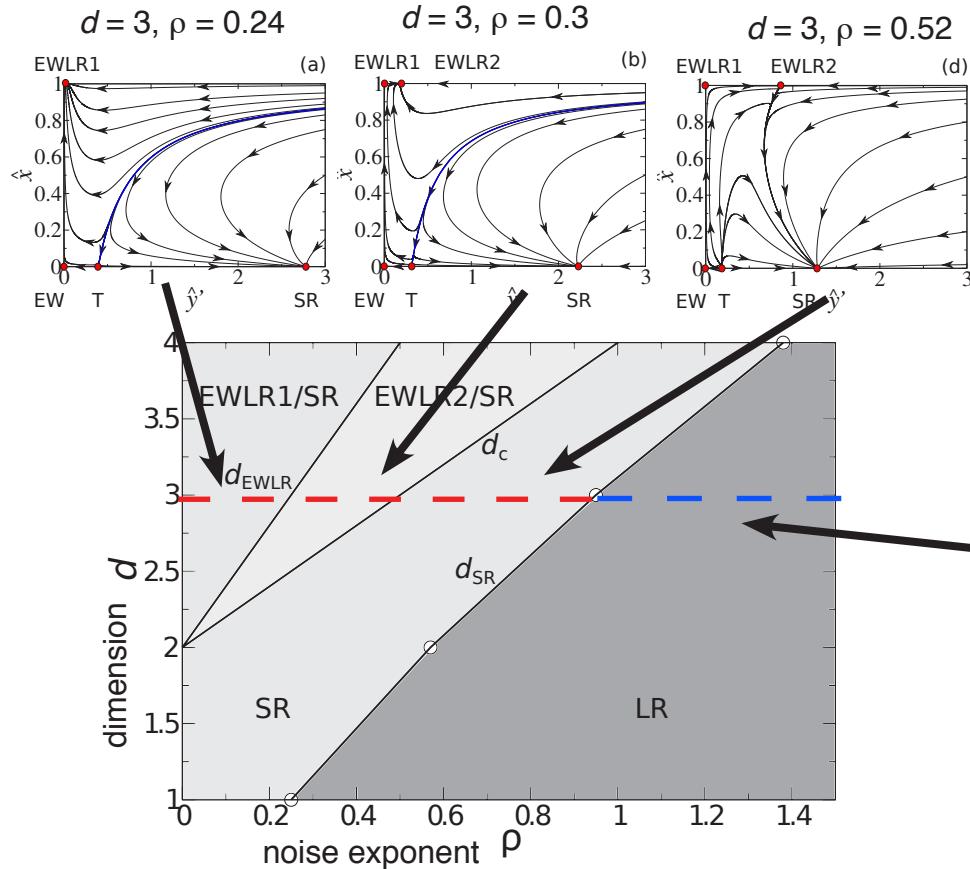
$$\hat{f}_\kappa^x(\omega, q) \rightarrow 1 \quad (\text{LPA}')$$

long-range noise

$$\partial_s \hat{x}_\kappa = \hat{x}_\kappa (1 - \hat{x}_\kappa) (\eta_\kappa^D - 2\rho)$$

KPZ coupling

$$\partial_s \hat{y}_\kappa = \hat{y}_\kappa (2\hat{x}_\kappa (\eta_\kappa^D - 2\rho) + d - 2 + 3\eta_\kappa^\nu - \eta_\kappa^D)$$



TK, Canet, Delamotte, Wschebor, PRE (2014).

## Anisotropy

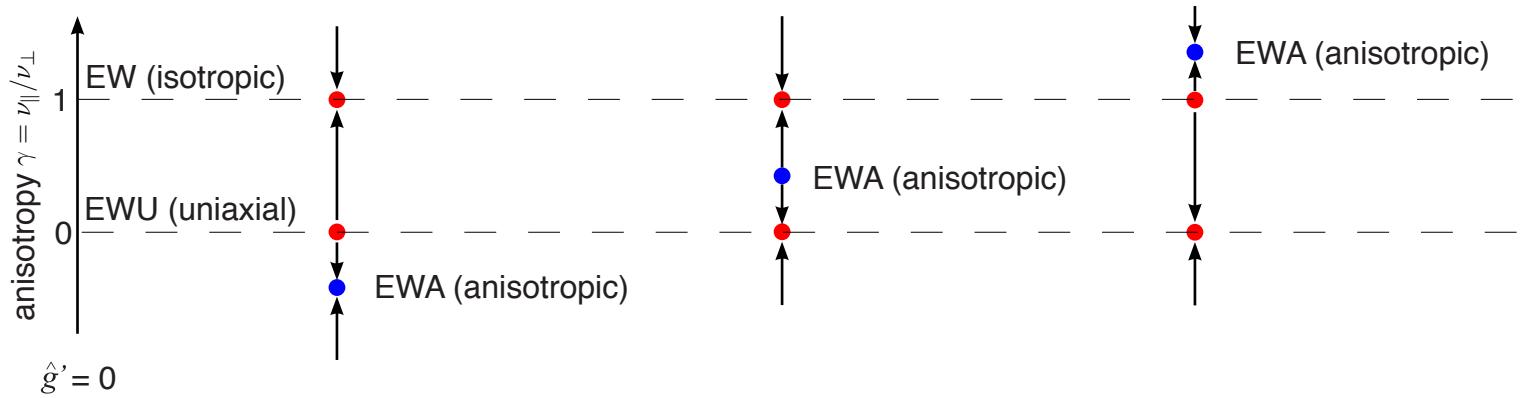
Anisotropic Kardar-Parisi-Zhang equation

$$\frac{\partial h(t, \vec{x})}{\partial t} = \nu_{\parallel} \nabla_{\parallel}^2 h(t, \vec{x}) + \nu_{\perp} \nabla_{\perp}^2 h(t, \vec{x}) + \frac{\lambda_{\parallel}}{2} (\nabla_{\parallel} h(t, \vec{x}))^2 + \frac{\lambda_{\perp}}{2} (\nabla_{\perp} h(t, \vec{x}))^2 + \eta(t, \vec{x})$$

$$\vec{x} = (\vec{x}_{\parallel}, \vec{x}_{\perp})^T, \quad d = d_{\parallel} + d_{\perp}$$

Wolf, PRL (1991)

Weak-coupling fixed points



Depending on the total dimension and the splitting  $d = d_{\parallel} + d_{\perp}$ , the isotropic KPZ equation may become unstable against anisotropic perturbations.

Täuber, Frey, Europhys. Lett. (2002)

## Anisotropy (NPRG)

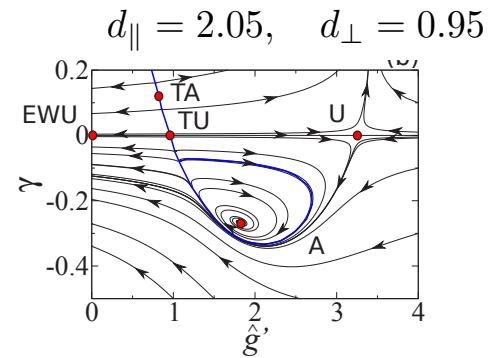
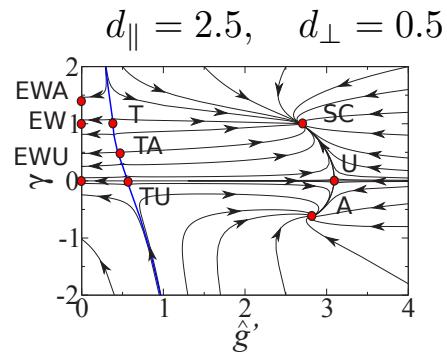
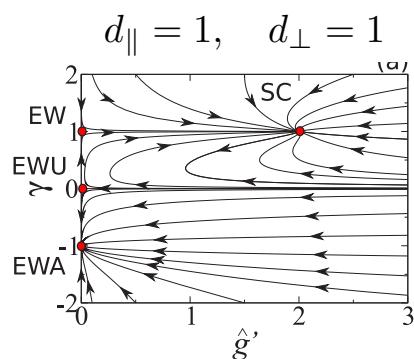
Ansatz

$$\Gamma_\kappa[\varphi, \tilde{\varphi}] = \int_{\mathbf{x}} \left\{ \tilde{\varphi} f_\kappa^\lambda(\nabla) D_t \varphi - (\nabla_\parallel^2 \varphi) f_\kappa^{\nu\parallel}(\nabla) \tilde{\varphi} - (\nabla_\perp^2 \varphi) f_\kappa^{\nu\perp}(\nabla) \tilde{\varphi} - \tilde{\varphi} f_\kappa^D(\nabla) \tilde{\varphi} \right\}$$

$$\begin{aligned} \hat{f}_\kappa^x(\hat{q}_\parallel, \hat{q}_\perp) & \quad (\text{NLO}) \\ \hat{f}_\kappa^x(\hat{q}_\parallel, \hat{q}_\perp) & \rightarrow 1 \quad (\text{LPA'}) \end{aligned}$$

$$D_t \varphi = \partial_t \varphi - \frac{\lambda_\parallel}{2} (\nabla_\parallel \varphi)^2 - \frac{\lambda_\perp}{2} (\nabla_\perp \varphi)^2$$

RG trajectories (LPA')



New anisotropic (A) and uniaxial (U) fixed points

The isotropic strong coupling (SC) fixed point is always stable.

## Conclusion

NPRG approach to the KPZ problem (including long-range noise and anisotropy)

naive extrapolation from weak- to strong-coupling behavior not possible

strong-coupling behavior changes qualitatively the phase diagram

no hint for an upper critical dimension

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