

# NUCLEAR CHIRAL THERMODYNAMICS and the FUNCTIONAL RENORMALIZATION GROUP



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ECT\* Trento and Technische Universität München

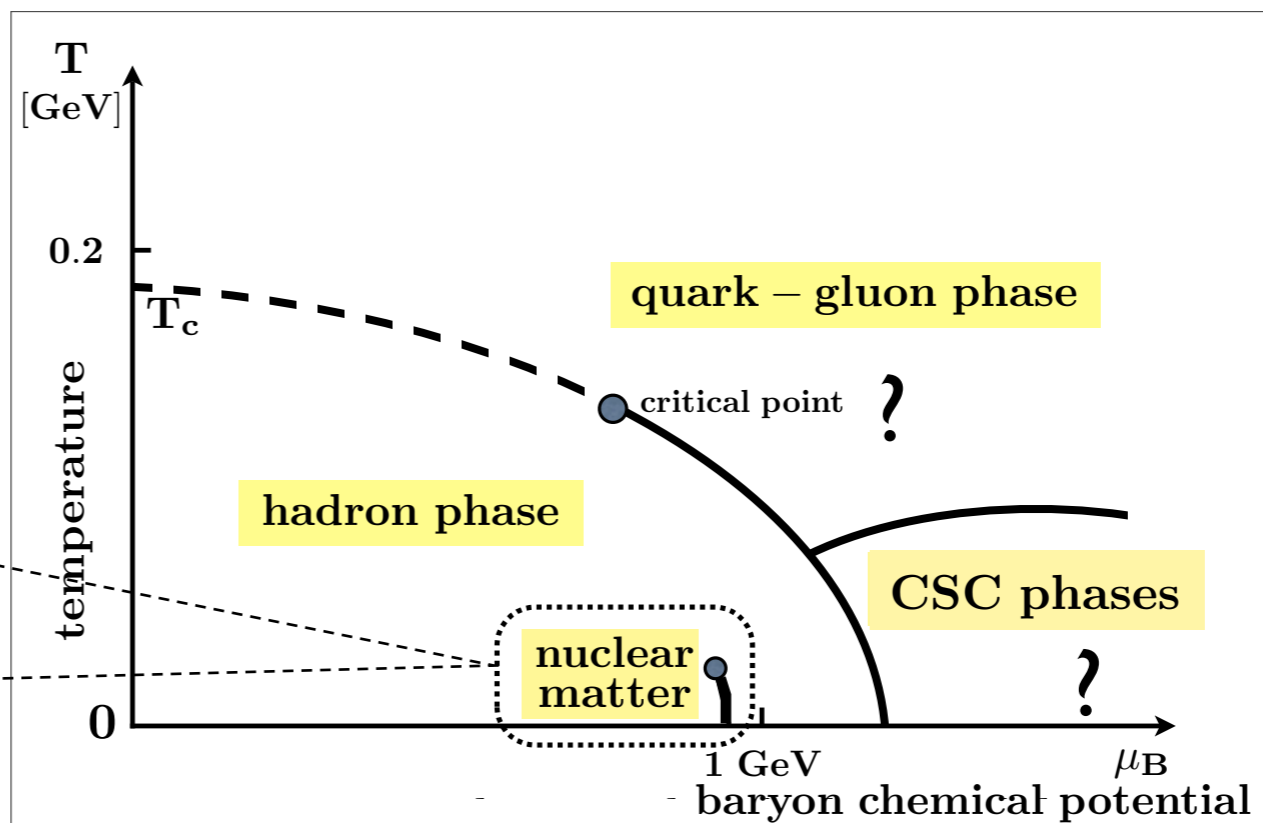
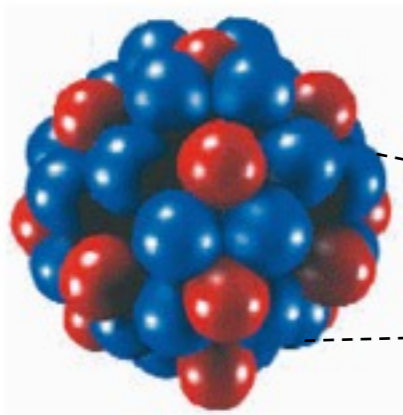


- Introductory glance at the **QCD phase diagram**
- **Chiral models** of the **nuclear equation of state**
- Beyond mean field:  
**non-perturbative fluctuations** and **FRG** approach
- Symmetric and asymmetric **nuclear matter**
- **Neutron matter** and **neutron stars**
- Density & temperature dependence of **chiral order parameter**

(based on doctoral thesis work of **Matthias Drews**)

# NUCLEAR MATTER and QCD PHASES

nuclei



## Scales in $N = Z$ nuclear matter:

- momentum scale:  
**Fermi momentum**
- NN distance:
- energy per nucleon:
- compression modulus:

$$k_F \simeq 1.4 \text{ fm}^{-1} \sim 2m_\pi$$

$$d_{NN} \simeq 1.8 \text{ fm} \simeq 1.3 m_\pi^{-1}$$

$$E/A \simeq -16 \text{ MeV}$$

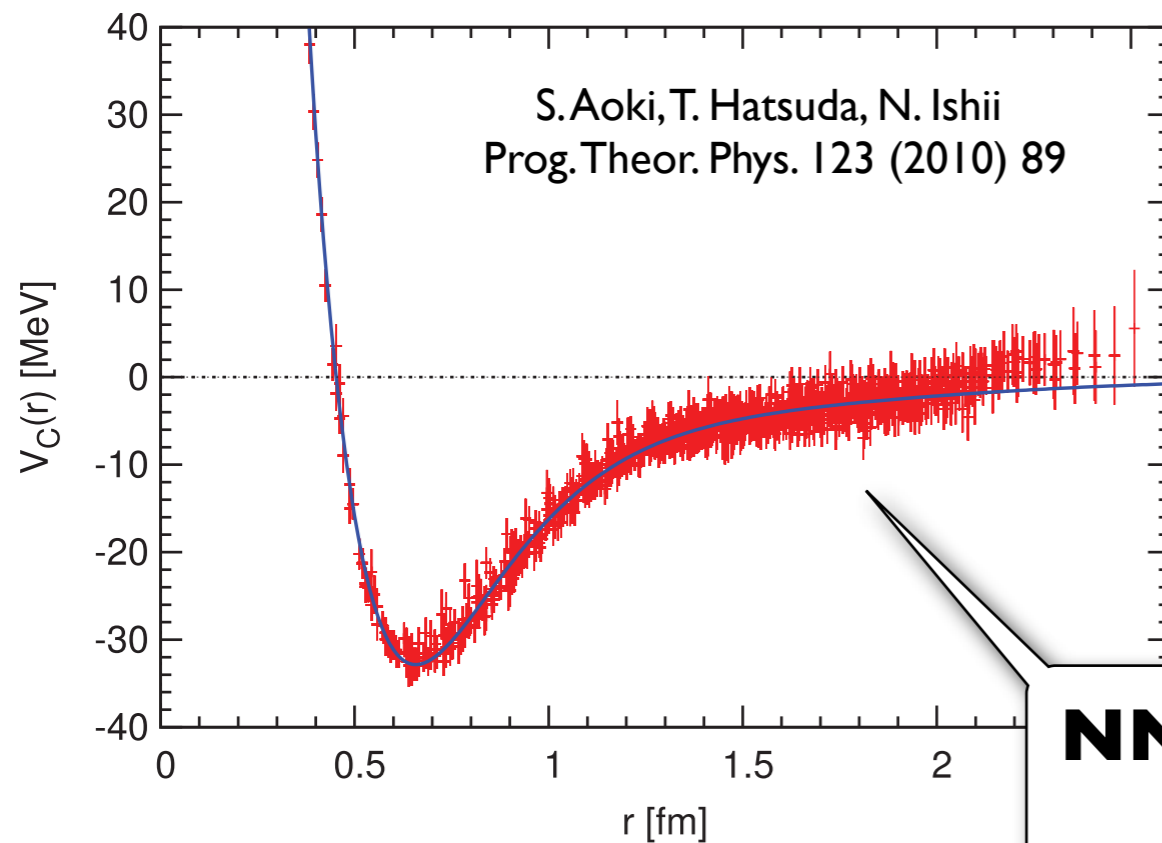
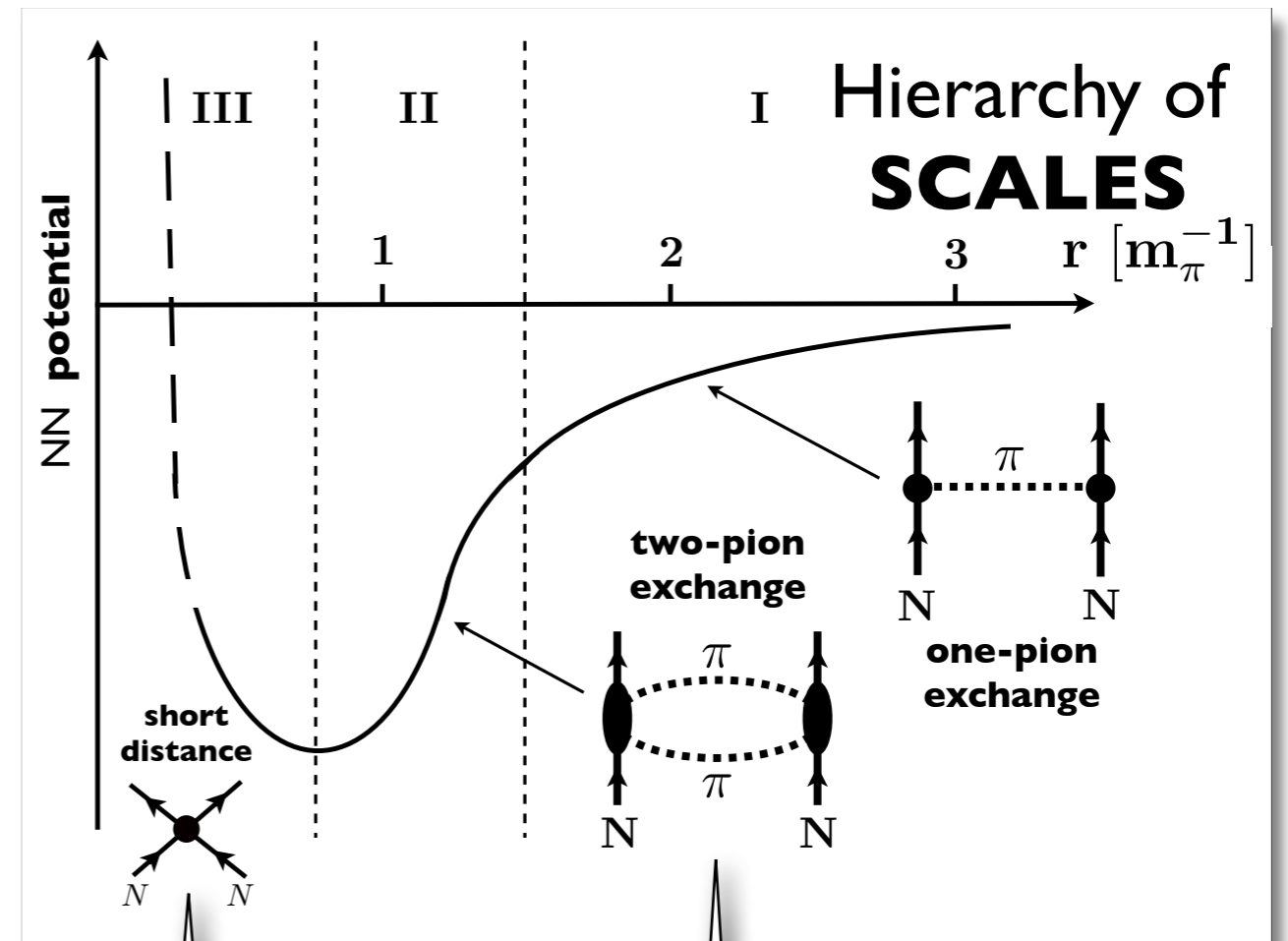
$$K = (260 \pm 30) \text{ MeV} \sim 2m_\pi$$

# Nuclear Forces

contemporary approaches:

**Chiral Effective  
Field Theory  
+  
Lattice QCD**

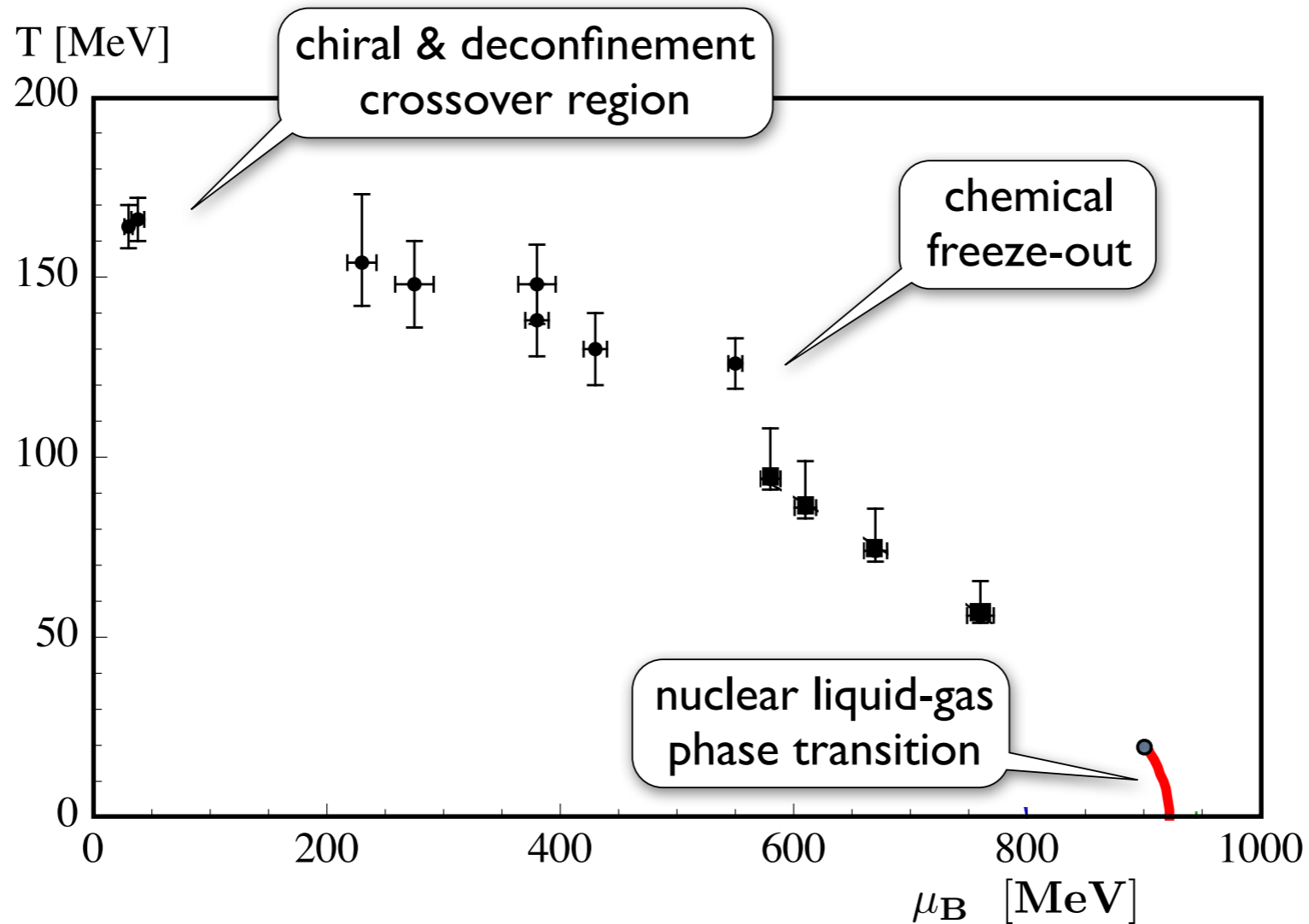
Early history: M. Taketani et al. (1951)



**NN Central Potential  
from Lattice QCD**

# PHASE DIAGRAM and CHEMICAL FREEZE-OUT

- Empirical freeze-out systematics: thermal hadron production



A. Andronic,  
P. Braun-Munzinger,  
J. Stachel

Phys. Lett.  
B 673 (2009) 142  
B 678 (2009) 516

# **PIONS, NUCLEONS and NUCLEI** in the context of **LOW-ENERGY QCD**

- **CONFINEMENT** of quarks and gluons in hadrons
  - Spontaneously broken **CHIRAL SYMMETRY**
  - **LOW-ENERGY QCD** with light (u,d) quarks:  
**Effective Field Theory** of (weakly) interacting  
**Nambu-Goldstone Bosons** (pions)
- **Chiral EFT** represents QCD at energy/momentum scales  
$$Q \ll 4\pi f_\pi \sim 1 \text{ GeV}$$
  - **Strategies** at the interface between QCD and **nuclear physics** :

In-medium **Chiral Perturbation Theory**  
based on **non-linear sigma model**  
(with inclusion of nucleons)

**expansion of free energy density  
in powers of Fermi momentum**

**Chiral Nucleon-Meson** model  
based on **linear sigma model**

**non-perturbative  
Renormalization Group approach**

# *PART I:*

## In-medium Chiral Perturbation Theory and the nuclear many-body problem

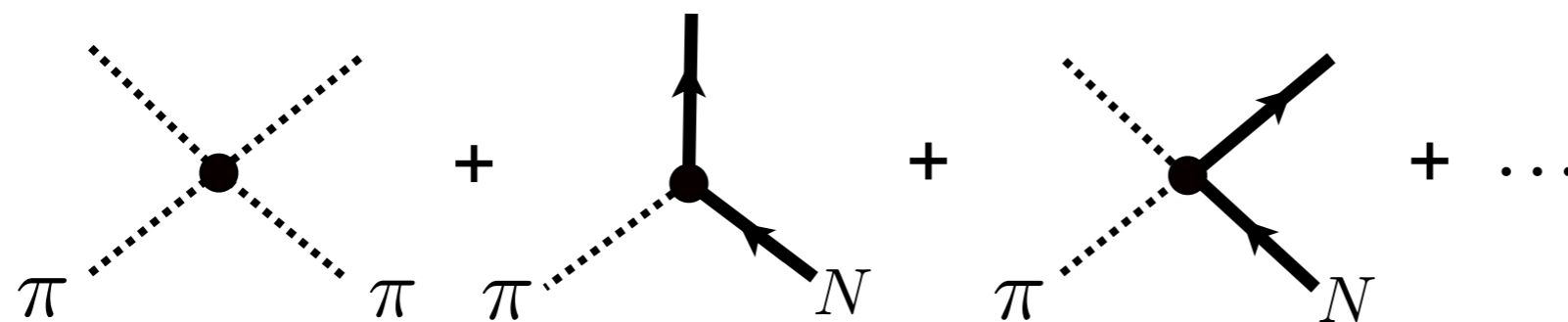
# CHIRAL EFFECTIVE FIELD THEORY

- Interacting systems of **PIONS** (light / fast) and **NUCLEONS** (heavy / slow):

$$\mathcal{L}_{eff} = \mathcal{L}_\pi(U, \partial U) + \mathcal{L}_N(\Psi_N, U, \dots)$$

$$U(x) = \exp[i\tau_a \pi_a(x) / f_\pi]$$

- Construction of Effective Lagrangian: **Symmetries**



**short  
distance  
dynamics:  
contact terms**

# NUCLEAR INTERACTIONS from CHIRAL EFFECTIVE FIELD THEORY

Weinberg

Bedaque & van Kolck

Bernard, Epelbaum, Kaiser, Meißner; ...

	Two-nucleon force	Three-nucleon force	Four-nucleon force
$\mathcal{O}\left(\frac{Q^0}{\Lambda^0}\right)$		—	—
$\mathcal{O}\left(\frac{Q^2}{\Lambda^2}\right)$		—	—
$\mathcal{O}\left(\frac{Q^3}{\Lambda^3}\right)$			—
$\mathcal{O}\left(\frac{Q^4}{\Lambda^4}\right)$			



**Systematically organized HIERARCHY**





# IN-MEDIUM CHIRAL PERTURBATION THEORY

- **Small scales:**

energy, momentum,  $m_\pi$ ,  $k_F \ll 4\pi f_\pi \sim 1 \text{ GeV}$

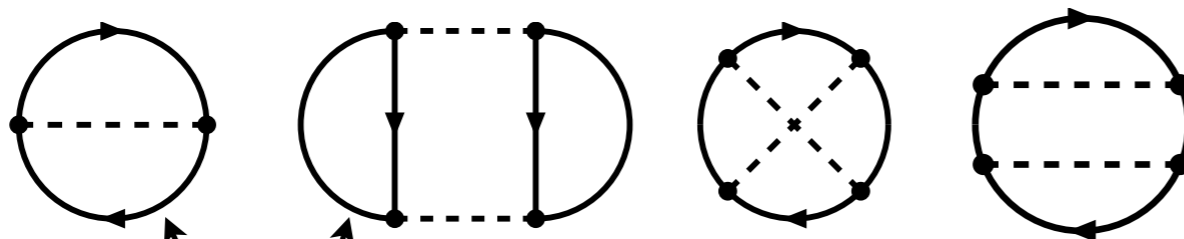
- **Loop expansion of (In-Medium) Chiral Perturbation Theory**



Systematic expansion of **ENERGY DENSITY**  $\mathcal{E}(k_F)$  in **powers of Fermi momentum** [modulo functions  $f_n(k_F/m_\pi)$ ]

(works for  $k_F \ll 4\pi f_\pi \sim 1 \text{ GeV}$ )

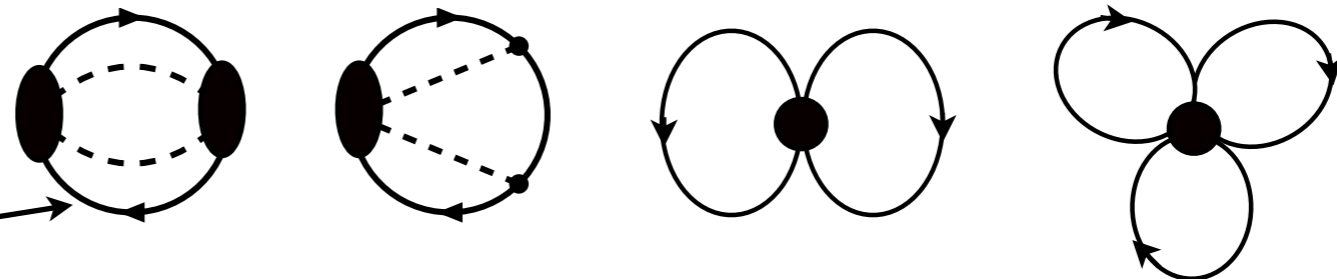
- Nuclear **thermodynamics**: compute **free energy density**



(3-loop order)

N. Kaiser, S. Fritsch, W.W.  
(2002-2004)

**in-medium**  
nucleon propagators  
incl. Pauli blocking



# NUCLEAR MATTER

- **In-medium ChPT**

3-loop  $(\pi, N, \Delta)$

- **Input** parameters:  
two contact terms

- basically:  
analytic calculation

- **Output:**

- ▶ Binding & saturation

$$E_0/A = -16 \text{ MeV} , \rho_0 = 0.16 \text{ fm}^{-3} , K = 290 \text{ MeV}$$

- ▶ Realistic (complex, momentum dependent) single-particle potential  
... satisfying Hugenholtz - van Hove and Luttinger theorems (!)

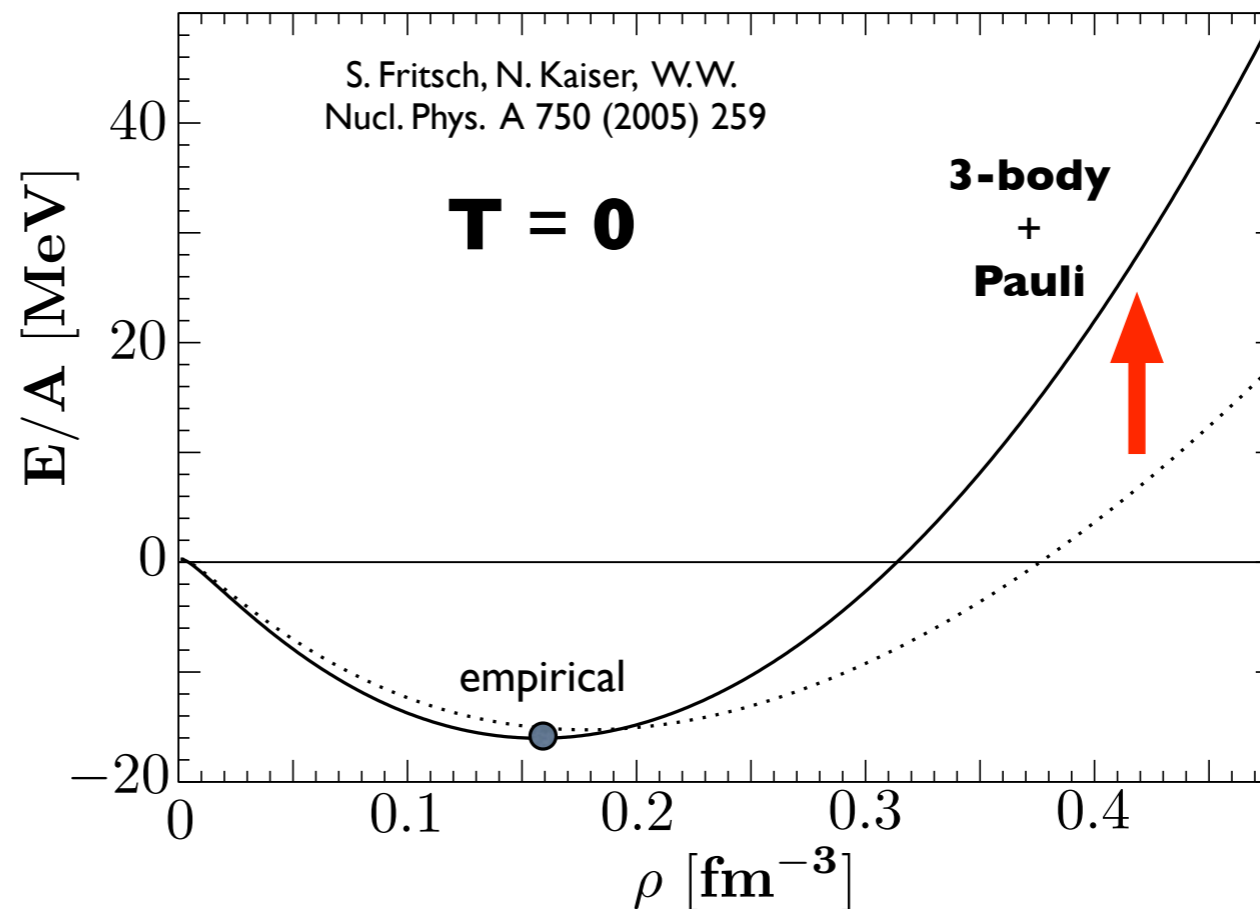
J.W. Holt, N. Kaiser, W.W.  
Nucl. Phys. A 870 (2011) 1,  
Nucl. Phys. A 876 (2012) 61,  
Phys. Rev. C 87 (2013) 014338

- ▶ Asymmetry energy:  $A(k_F^0) = 34 \text{ MeV}$

- ▶ Fermi Liquid Theory:

Quasiparticle interaction and Landau parameters

C.Wellenhofer, J.W. Holt, N. Kaiser, W.W.  
Phys. Rev. C 89 (2014) 064009

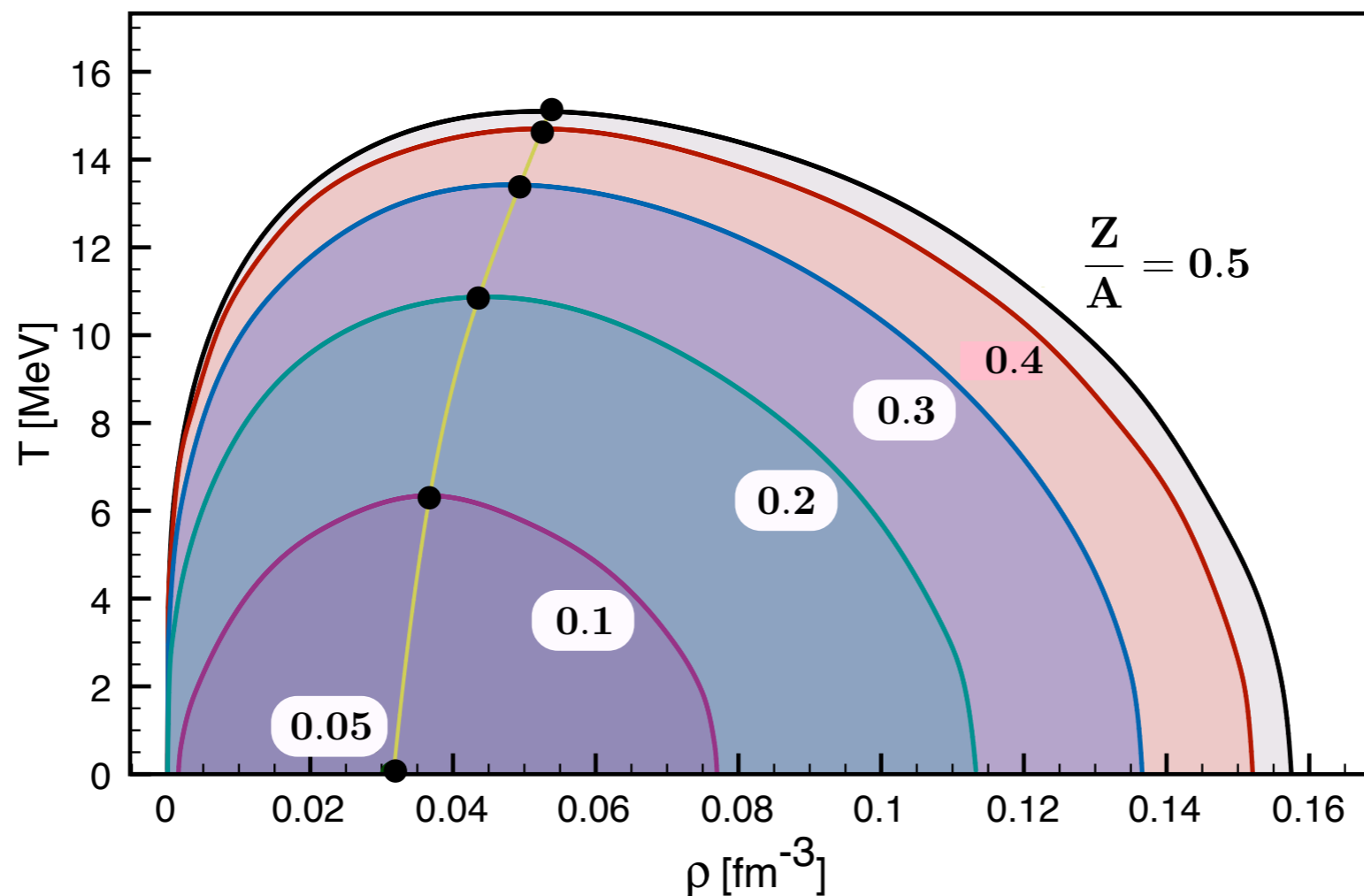


Recent review: J.W. Holt, N. Kaiser, W.W. Prog. Part. Nucl. Phys. 73 (2013) 35



# PHASE DIAGRAM of NUCLEAR MATTER

- Trajectory of **CRITICAL POINT** for **asymmetric matter** as function of proton fraction  $Z/A$

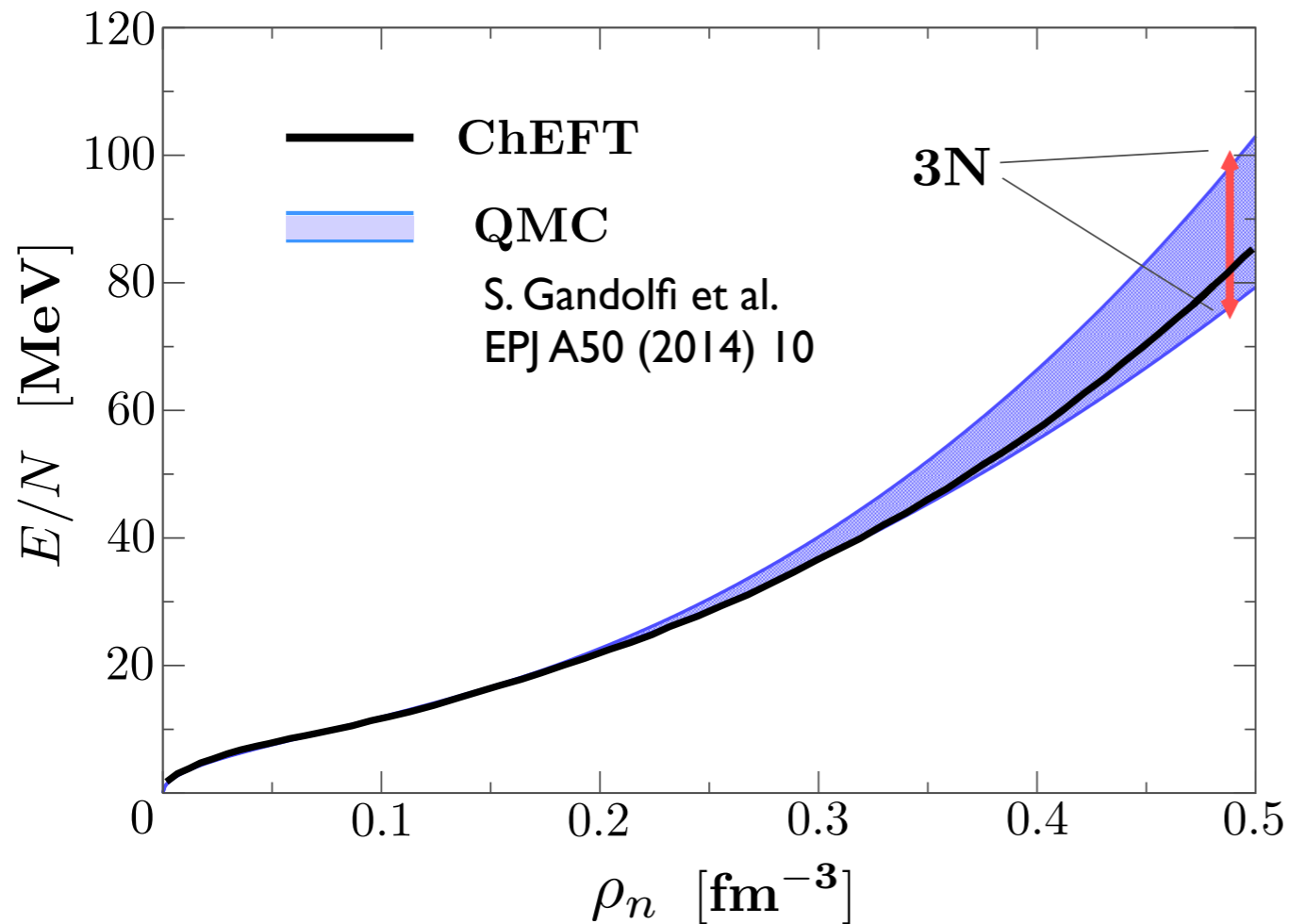


S. Fiorilla, N. Kaiser, W.W.  
Nucl. Phys. A880 (2012) 65

... determined almost completely by  
**isospin** dependent (one- and two-) **pion** exchange dynamics

# NEUTRON MATTER

- In-medium chiral effective field theory (3-loop) with resummation of short distance contact terms (large nn scattering length,  $a_s = 19$  fm)



- Neutron matter behaves almost (but not quite) like a unitary Fermi gas

- Bertsch parameter

$$\xi = \frac{\bar{E}}{E_{\text{Fermi gas}}} \simeq 0.5$$

J.W.Holt, N.Kaiser, W.W.  
 Phys. Rev. C 87 (2013) 014338

- agreement with sophisticated many-body calculations (e.g. recent **Q**uantum **M**onte **C**arlo computations)



# *PART II:*

## Chiral Nucleon-Meson Model and Functional Renormalization Group

# Mesons, Nucleons, Nuclear Matter and Functional Renormalization Group

- **Chiral nucleon - meson model**  $\psi = (\psi_p, \psi_n)^T$

$$\begin{aligned} \mathcal{L} = & \bar{\psi} i \gamma_\mu \partial^\mu \psi + \frac{1}{2} \partial_\mu \sigma \partial^\mu \sigma + \frac{1}{2} \partial_\mu \boldsymbol{\pi} \cdot \partial^\mu \boldsymbol{\pi} \\ & - \bar{\psi} \left[ g(\sigma + i \gamma_5 \boldsymbol{\tau} \cdot \boldsymbol{\pi}) + \gamma_\mu (g_\omega \omega^\mu + g_\rho \boldsymbol{\tau} \cdot \boldsymbol{\rho}^\mu) \right] \psi \\ & - \frac{1}{4} F_{\mu\nu}^{(\omega)} F^{(\omega)\mu\nu} - \frac{1}{4} \mathbf{F}_{\mu\nu}^{(\rho)} \cdot \mathbf{F}^{(\rho)\mu\nu} \\ & + \frac{1}{2} m_V^2 (\omega_\mu \omega^\mu + \boldsymbol{\rho}_\mu \cdot \boldsymbol{\rho}^\mu) - \mathcal{U}(\sigma, \boldsymbol{\pi}) \end{aligned}$$

- **Effective potential** constructed to reproduce standard nuclear thermodynamics around equilibrium
- **Mean field** calculations  
S. Floerchinger, Ch. Wetterich : Nucl. Phys. A 890-891 (2012) 11
- **Mesonic and nucleonic particle-hole** fluctuations treated non-perturbatively using **FRG**

M. Drews, T. Hell, B. Klein, W.W.

Phys. Rev. D 88 (2013) 096011

M. Drews, W.W.

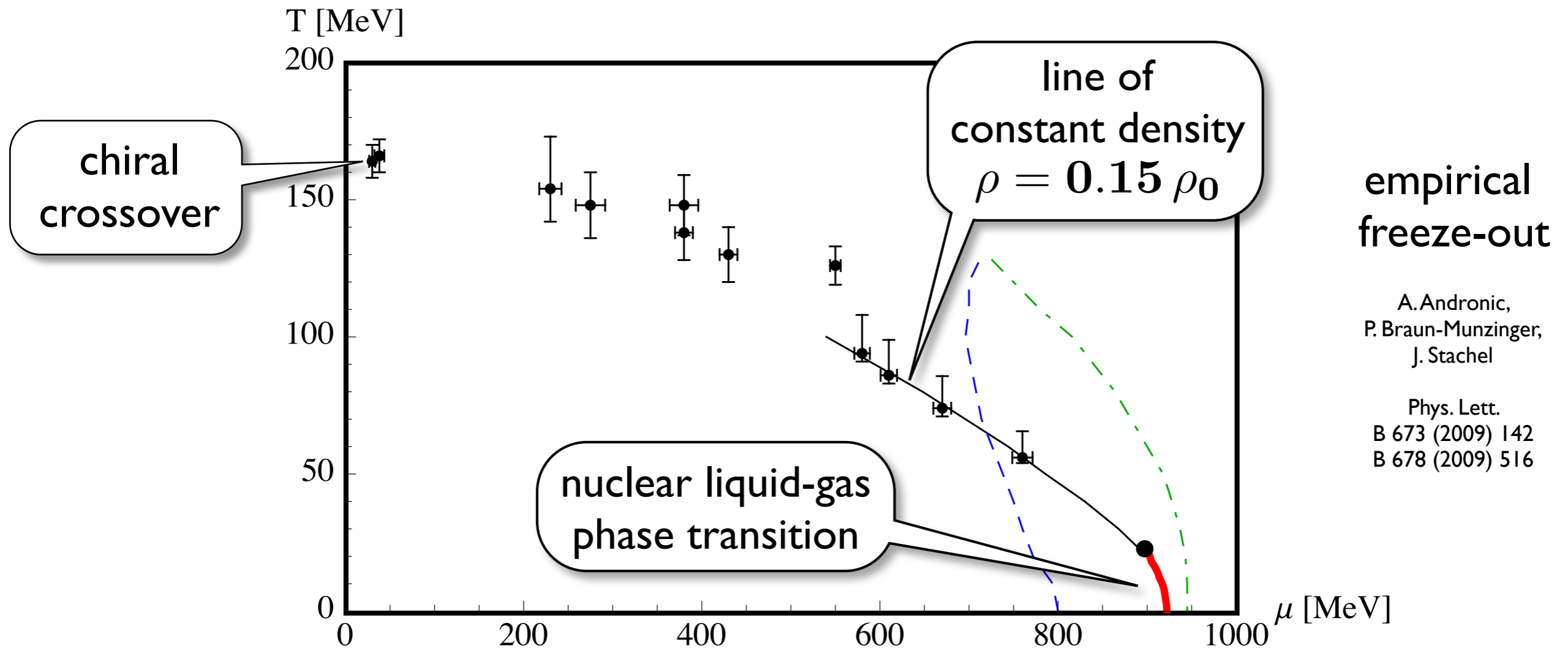
arXiv:1404.0882



# CHEMICAL FREEZE-OUT

S. Floerchinger, Ch. Wetterich : Nucl. Phys. A 890-891 (2012) 11

- **Chiral nucleon - meson model** in mean-field approximation



- **Chemical freeze-out** in **baryonic matter** at  $T < 100$  MeV is **not** associated with (**chiral**) phase transition or rapid crossover

# Fixing the input: some comments

- **Potential**  $\mathcal{U}(\sigma, \pi) = \mathcal{U}_0(\chi) - m_\pi^2 f_\pi (\sigma - f_\pi)$

chiral invariant part  
parametrized in powers of  
 $\chi = \frac{1}{2}(\sigma^2 + \pi^2)$

explicit chiral  
symmetry breaking

- **Scalar (“sigma”) field**

has mean-field (chiral **order parameter**) and fluctuating pieces.

$\sigma$  **mass**: NOT to be confused with pole in  $l = 0$  s-wave pion-pion T matrix.

**Nucleon mass**:  $m_N^2 = 2g \chi$  ... in vacuum:  $m_N = g f_\pi$

- **Vector fields** encode short-distance NN dynamics,  
treated as self-consistently determined background mean fields (non-fluctuating)  
(not to be identified with physical  $\omega$  and  $\rho$  mesons)

**Effective chemical potentials**  $\mu_{n,p}^{\text{eff}} = \mu_{n,p} - g_\omega \omega_0 \pm g_\rho \rho_0^3$

Relevant quantities:  $G_\rho = \frac{g_\rho^2}{m_V^2}$ ,  $G_\omega = \frac{g_\omega^2}{m_V^2}$   $\longleftrightarrow$  contact terms in ChEFT

- **Parameters**: 2 coefficients in  $\mathcal{U}_0$ ,  $m_\sigma \simeq 0.8 \text{ GeV}$ ,  $G_\rho \sim G_\omega/4 \simeq 1 \text{ fm}^2$   
determined by tuning nuclear matter properties and symmetry energy





# Chiral nucleon - meson model beyond mean-field - Renormalization Group strategies -

M. Drews, T. Hell, B. Klein, W.W. Phys. Rev. D 88 (2013) 096011

## Fluctuations: Wetterich's RG flow equations

**effective action**

**full propagator**

$$k \frac{\partial \Gamma_k}{\partial k} = \text{Diagram} = \frac{1}{2} \text{Tr} \frac{k \frac{\partial R_k}{\partial k}}{\Gamma_k^{(2)} + R_k}$$

The diagram shows a circle with a cross in a square on the left and a solid dot on the right.

**regulator:**

$$R_k(p^2) = (k^2 - p^2) \theta(k^2 - p^2)$$

C. Wetterich:  
Phys. Lett. B 301 (1993) 90

D.F. Litim, J.M. Pawłowski:  
JHEP 0611 (2006) 026  
J.-P. Blaizot, A. Ipp,  
R. Mendez-Galain, N. Wschebor:  
NP A784 (2007) 376

## Thermodynamics

**nucleons**

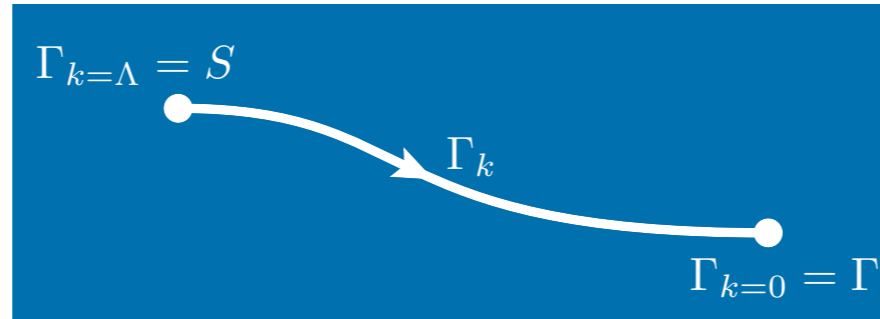
**pions**

$$k \partial_k \bar{\Gamma}_k(T, \mu) = \left( \text{Diagram}_1 + \text{Diagram}_2 \right) \Big|_{T, \mu_p, \mu_n} - \left( \text{Diagram}_1 + \text{Diagram}_2 \right) \Big|_{T=0, \mu = \mu_0 (= m_N - E_0/A)}$$

The diagrams are circles with a cross in a square on the left and a solid dot on the right. The first two diagrams have solid lines, while the last two have dashed lines.

# Flow equations in practice

UV scale:  
 $\Lambda = 1.4 \text{ GeV}$



“full” effective action &  
 effective potential  $U_{k=0}$

$$k \frac{\partial U_{k,\chi}}{\partial k} (T, \mu_p, \mu_n, \chi, \omega_0, \rho_0^3) = \text{[Diagram 1]} + \text{[Diagram 2]}$$

$$= \frac{k^5}{12\pi^2} \left\{ \frac{1 + 2n_B(E_\sigma)}{E_\sigma} + \frac{3[1 + 2n_B(E_\pi)]}{E_\pi} - 4 \sum_{i=n,p} \frac{1 - n_F(E_N - \mu_{i,\text{eff}})}{E_N} \right\}$$

$$E_\pi^2 = k^2 + U'_k(\chi), \quad E_\sigma^2 = k^2 + U'_k(\chi) + 2\chi U''_k(\chi),$$

$$U'_k(\chi) = \frac{\partial U_k(\chi)}{\partial \chi}, \quad E_N^2 = k^2 + 2g^2 \chi,$$

$$\mu_{n,p}^{\text{eff}}(k) = \mu_{n,p} - g_\omega \omega_0(k) \pm g_\rho \rho_0^3(k),$$

$$n_B(E) = \frac{1}{e^{E/T} - 1},$$

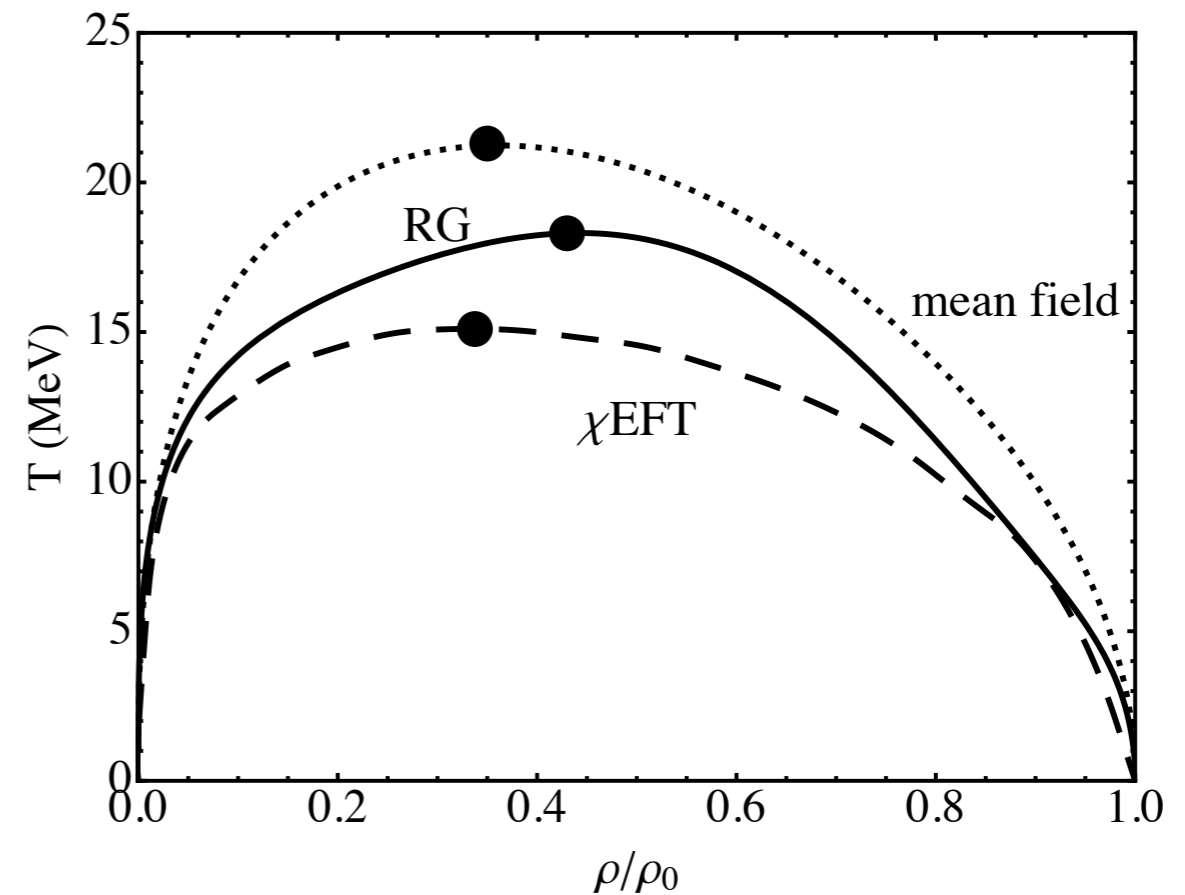
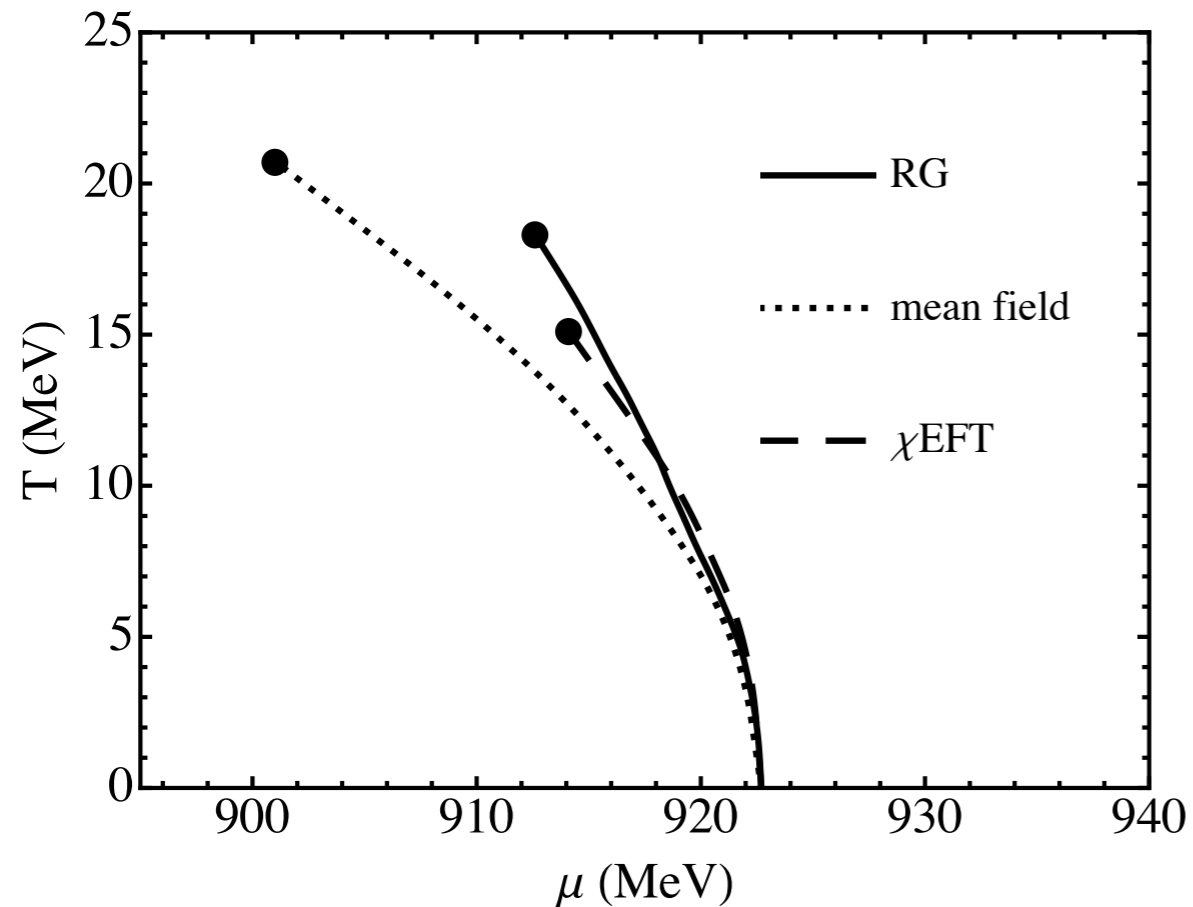
$$n_F(E) = \frac{1}{e^{E/T} + 1}.$$

... plus vector field equations, then full system of equations solved on a grid.

# Results : Liquid - Gas Transition

- symmetric nuclear matter -

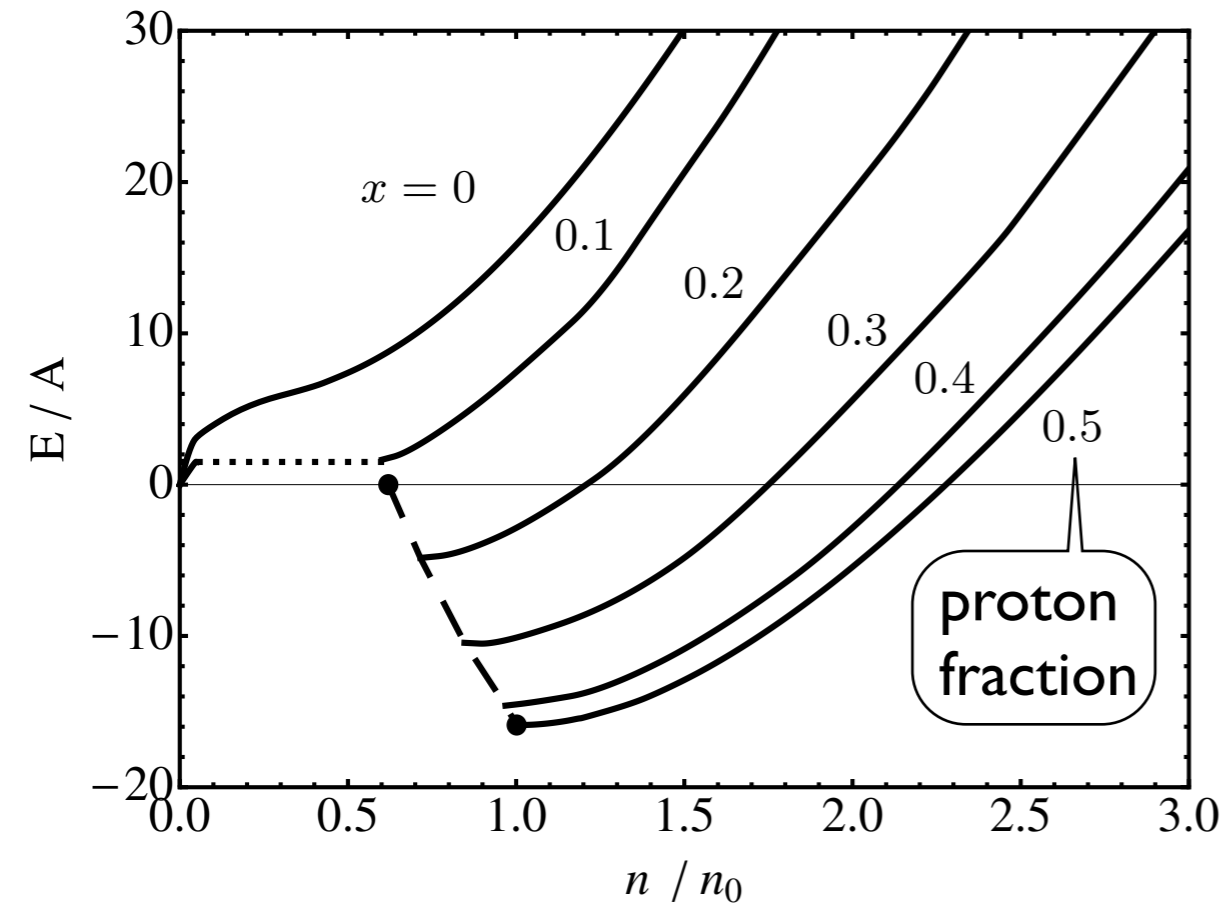
M. Drews, T. Hell, B. Klein, W.W.  
Phys. Rev. D 88 (2013) 096011



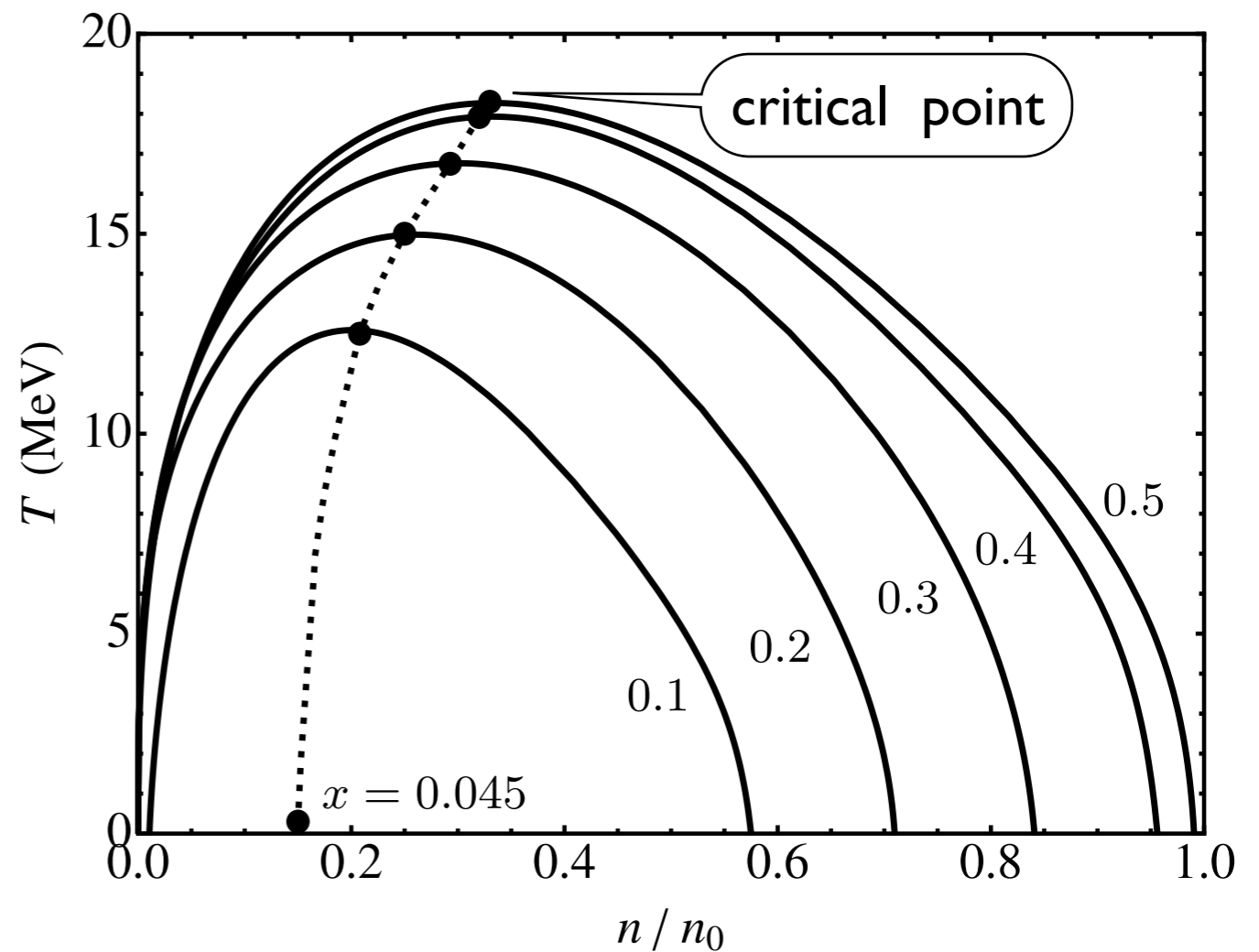
- note: close correspondence between (perturbative) in-medium ChEFT and (non-perturbative) FRG results

# Asymmetric nuclear matter in the **chiral RG** approach

M. Drews, W.W. (2014)



**Liquid-gas phase transition:**  
evolution of coexistence regions  
from symmetric to asymmetric nuclear matter

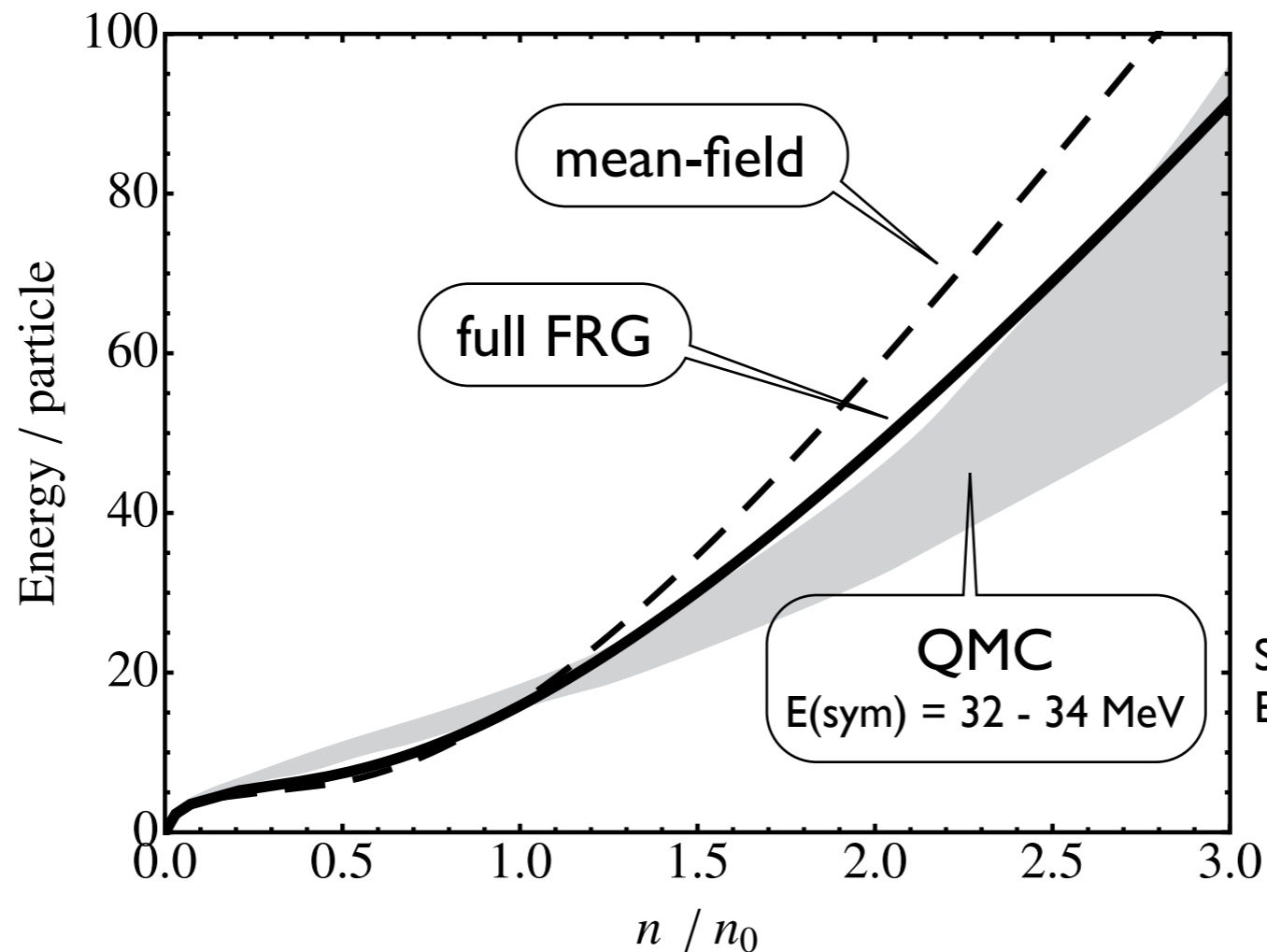


- **FRG results** (non-perturbative) are remarkably similar to in-medium **Chiral EFT calculations** (perturbative)



# Neutron matter in the chiral RG approach

M. Drews, W.W.  
arXiv:1404.0882

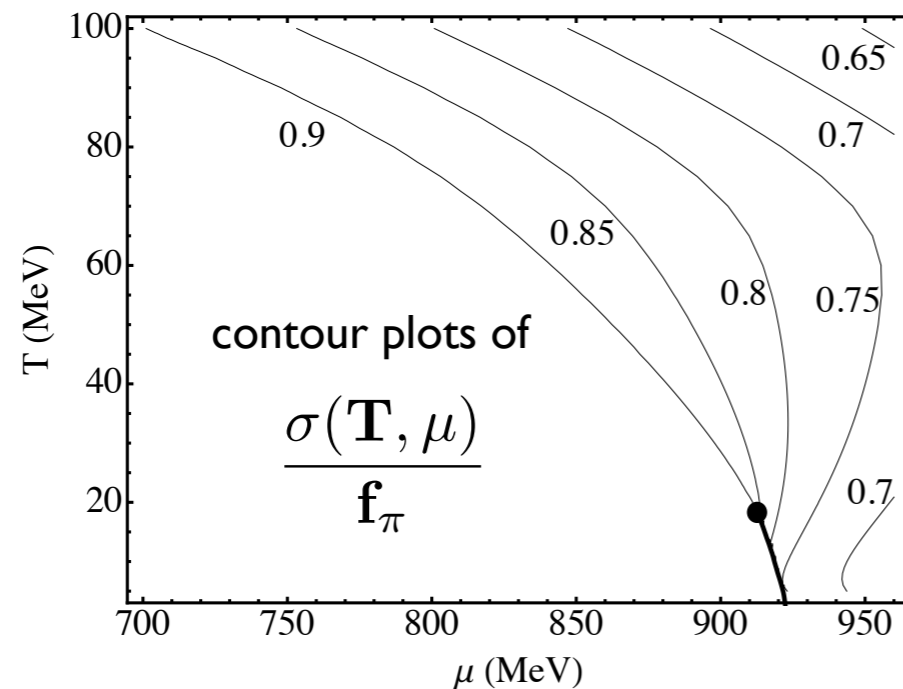
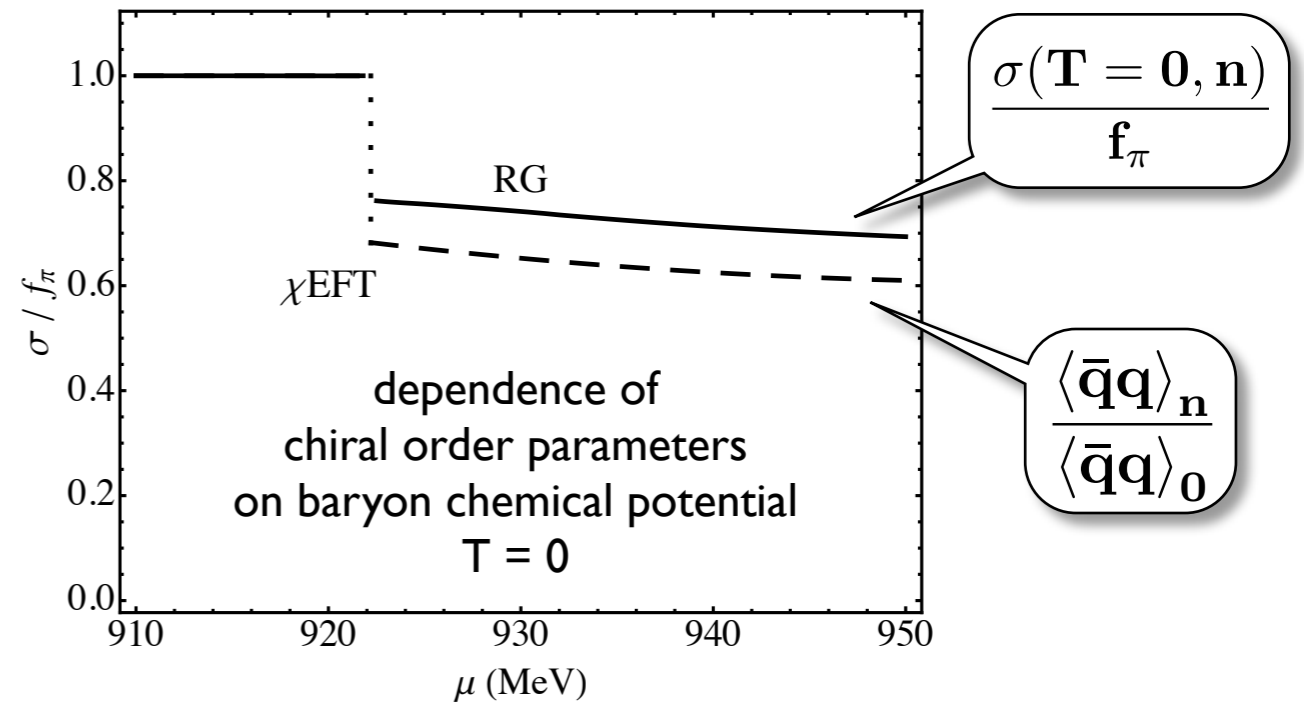
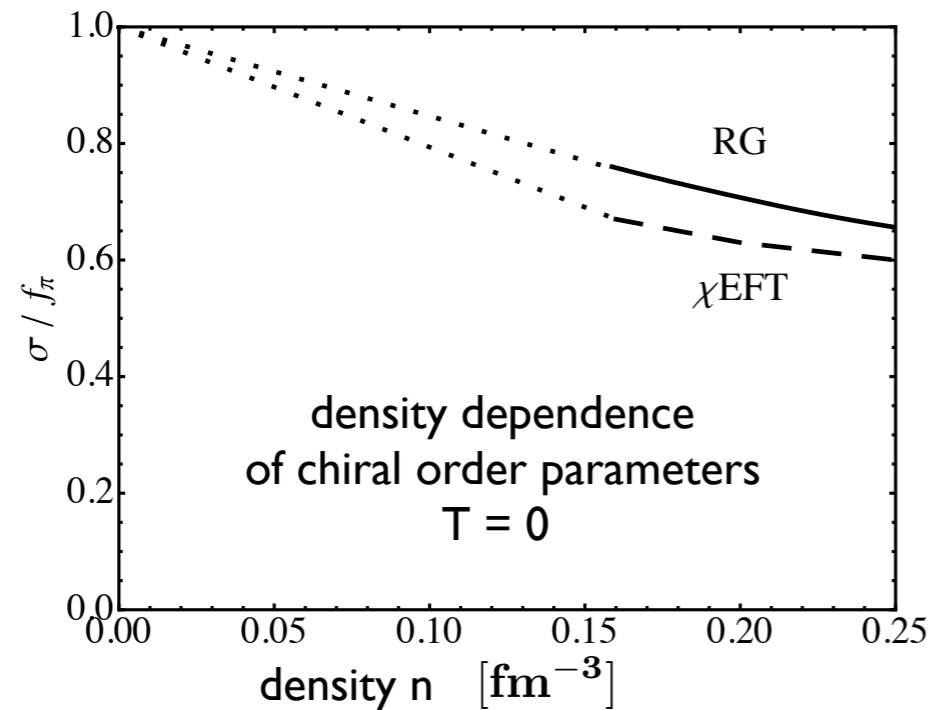


S. Gandolfi et al.  
EPJ A50 (2014) 10

- Single additional parameter:  
coupling strength of isovector-vector field / contact term  
fixed by symmetry energy  $E(\text{sym}) = 32 \text{ MeV}$

# Chiral Order Parameters

- Comparison of chiral effective field theory and model FRG results



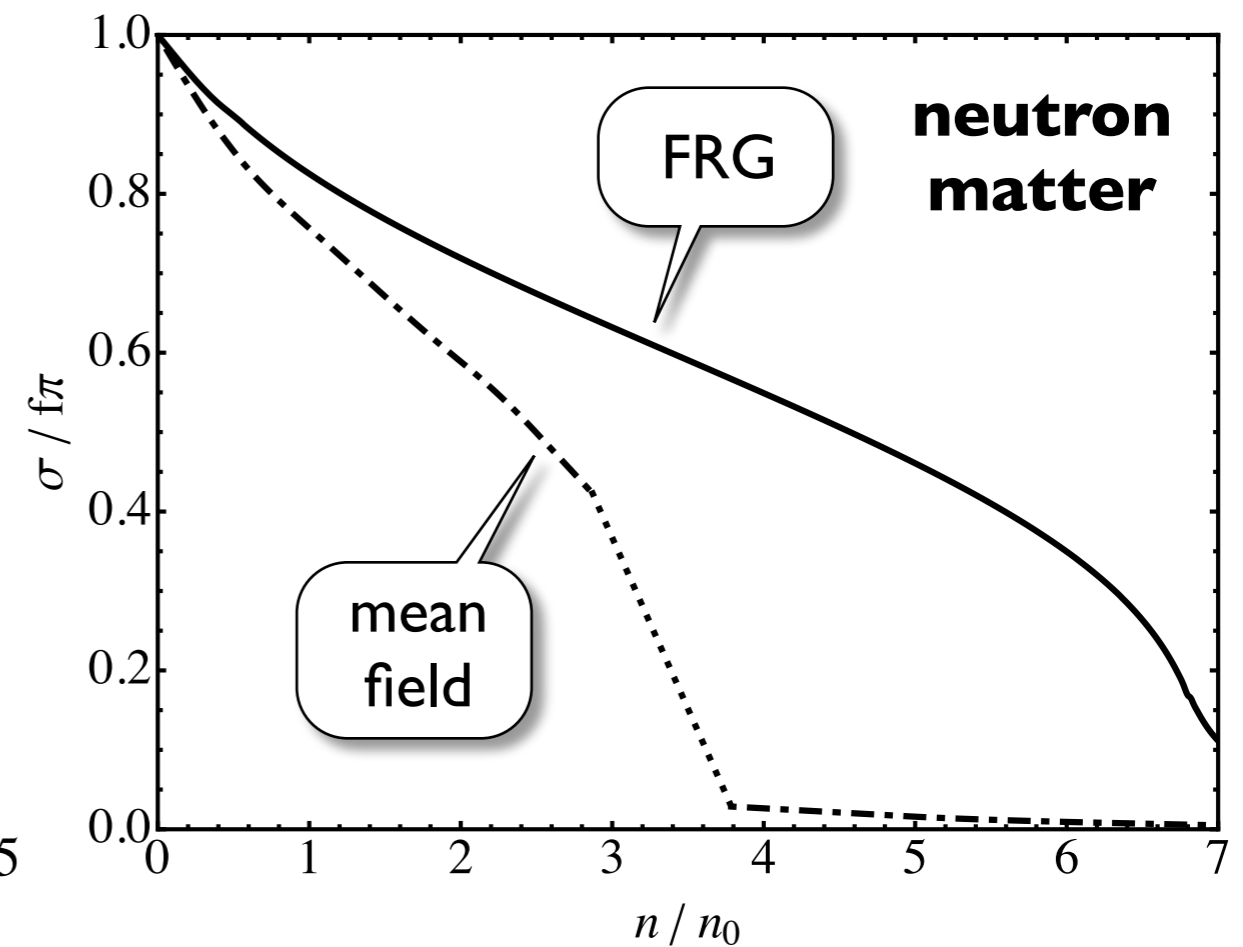
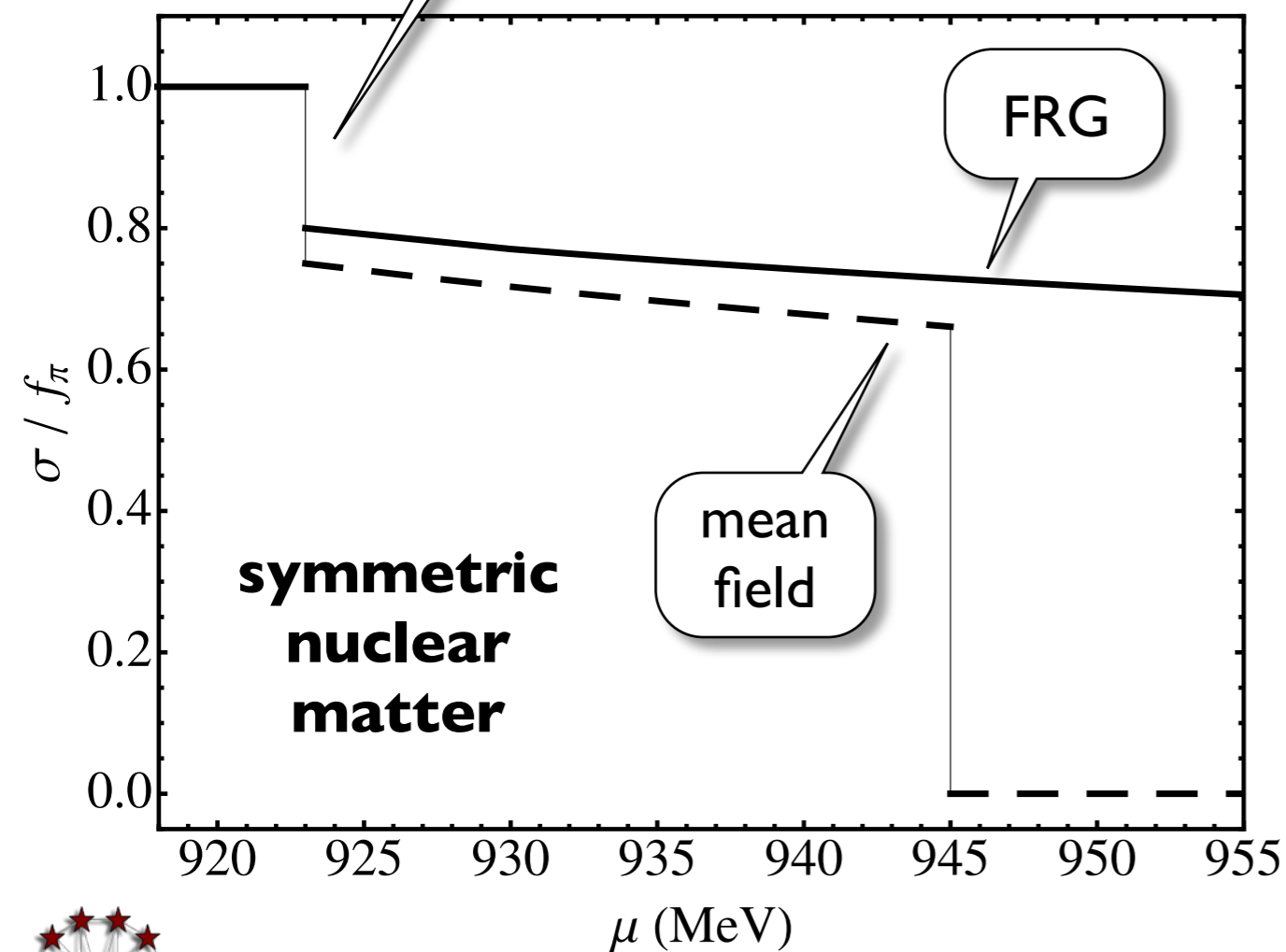
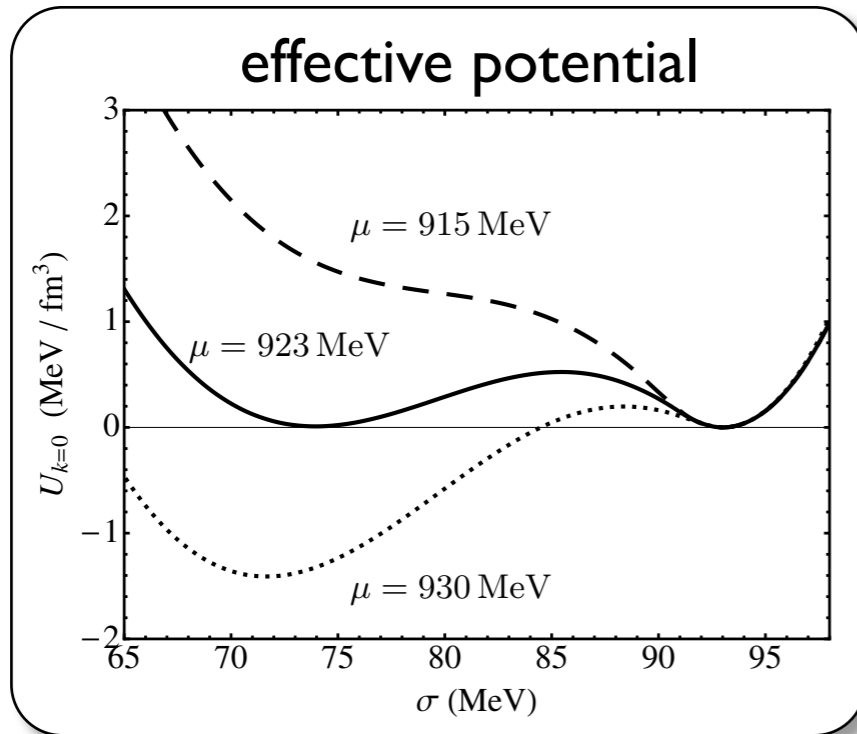
**No** tendency towards **chiral** phase transition for baryon chemical potentials  $\mu \lesssim 1 \text{ GeV}$  and temperatures  $T \lesssim 100 \text{ MeV}$

M. Drews, T. Hell, B. Klein, W.W. Phys. Rev. D 88 (2013) 096011



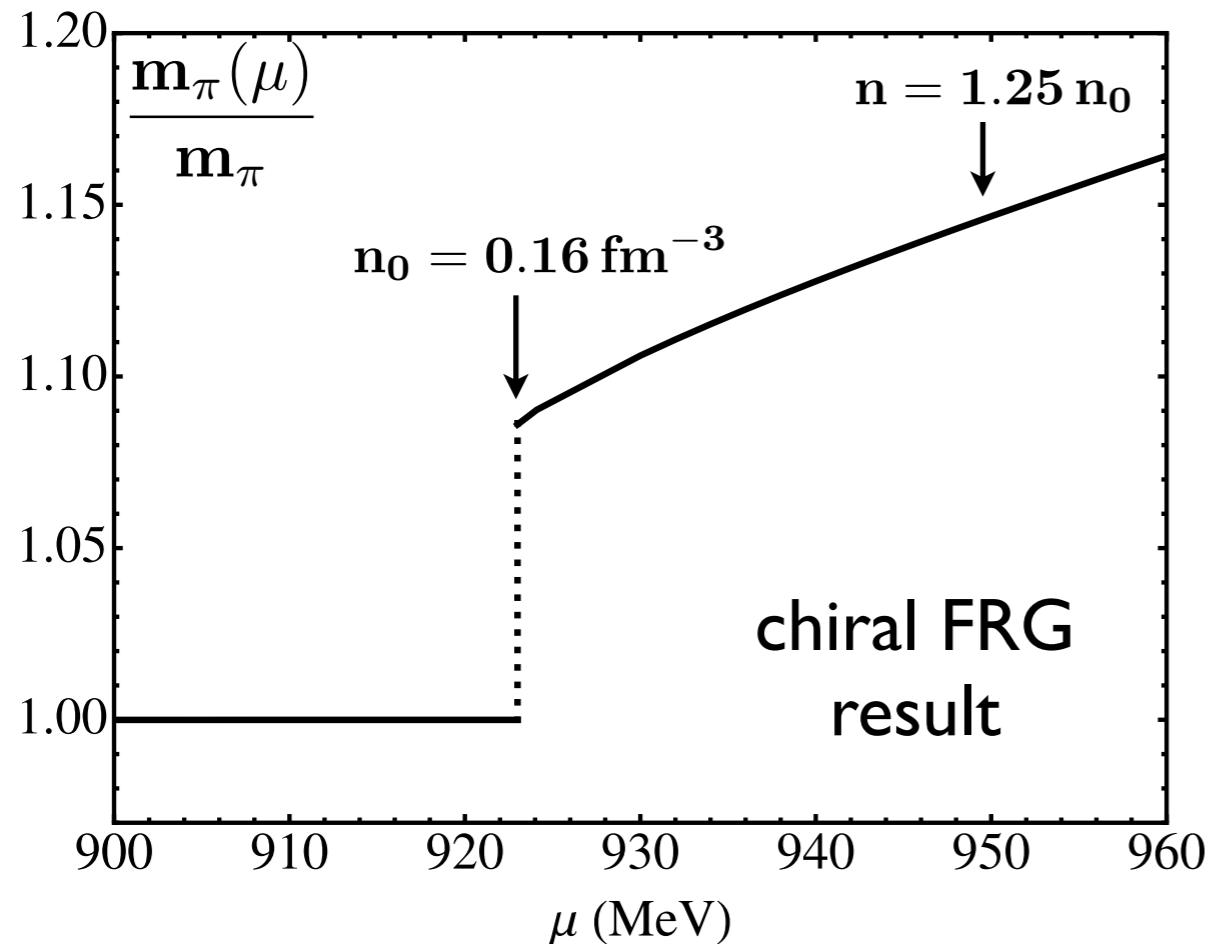
# Chiral Order Parameter

- important role of **fluctuations**:  
**disappearance of first-order chiral transition**  
 seen in mean-field approximation



# In-medium pion mass

- Contact with phenomenology :  
compare with s-wave pion-nuclear optical potential from pionic atoms



small

dominant

$$U(\mathbf{n}) = -\frac{2\pi}{m_\pi} \left[ b_0 - (b_0^2 + 2b_1^2) \left\langle \frac{1}{r} \right\rangle \right] \cdot \mathbf{n}$$

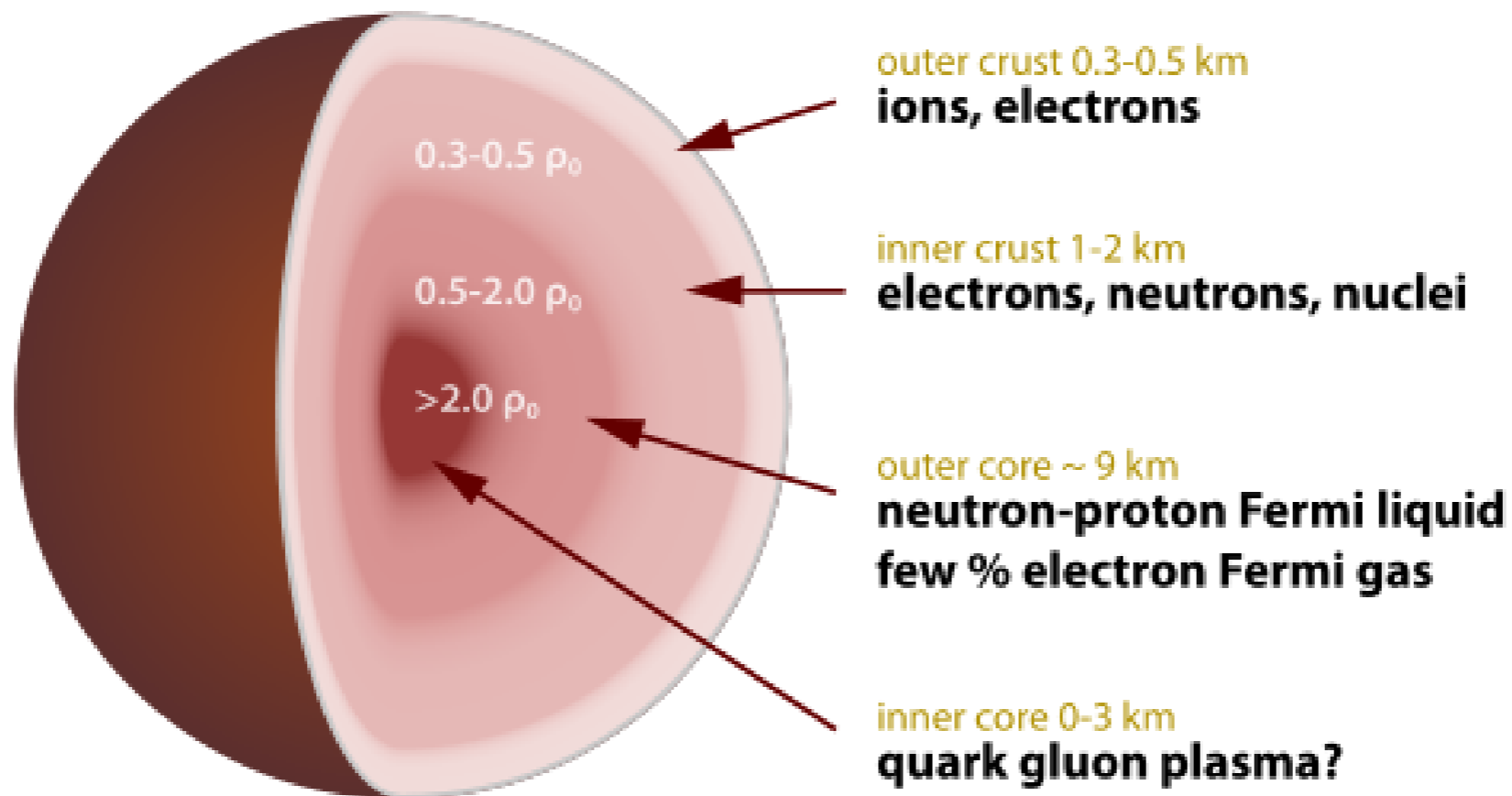
$$\frac{m_\pi(\mathbf{n})}{m_\pi} \simeq 1 + \frac{U(\mathbf{n})}{m_\pi} \simeq 1.1 \frac{\mathbf{n}}{n_0}$$

- Good agreement of **FRG** calculation with empirical **in-medium pion mass shift**, both in sign and magnitude



# *PART III:*

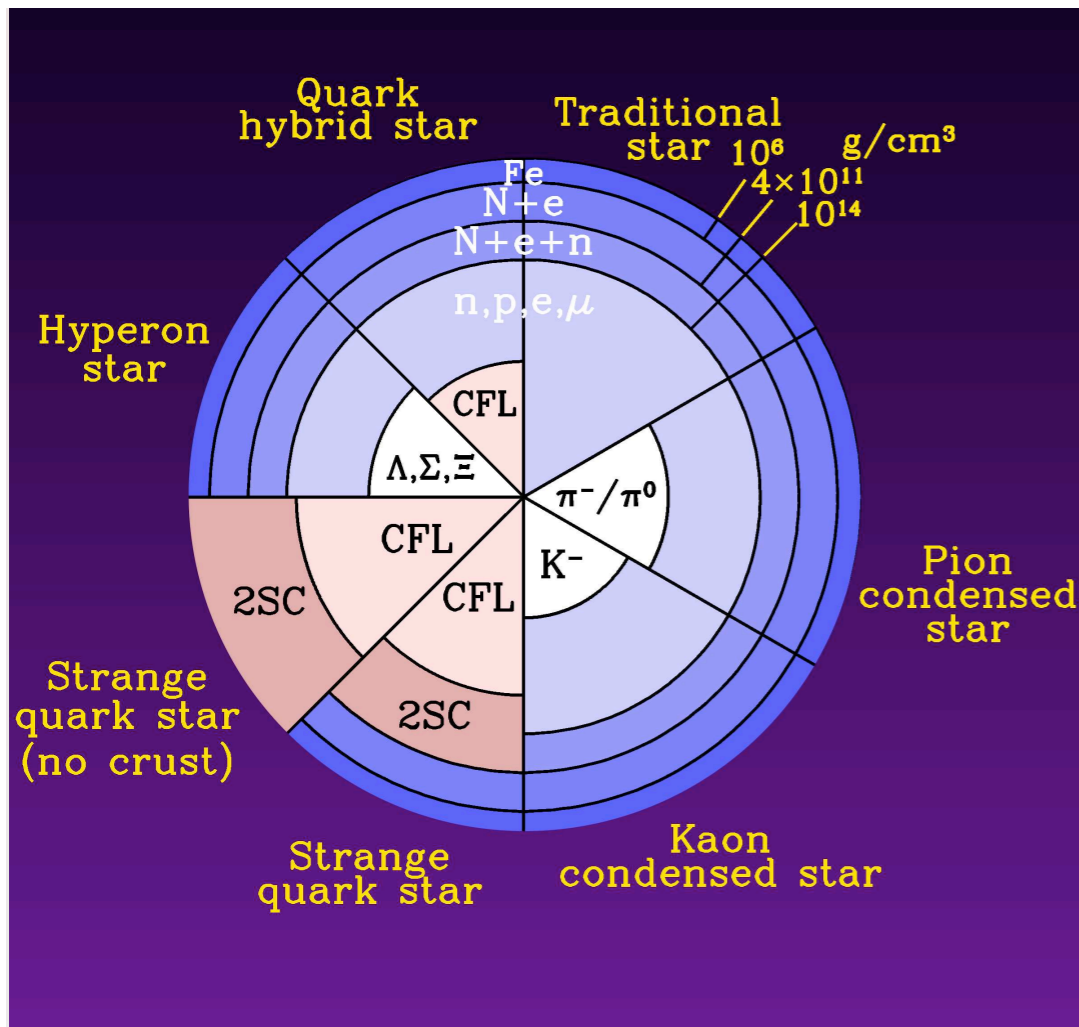
## Short digression on Neutron Stars



# NEUTRON STARS and the EQUATION OF STATE of DENSE BARYONIC MATTER

J. Lattimer, M. Prakash: *Astrophys. J.* 550 (2001) 426  
*Phys. Reports* 442 (2007) 109

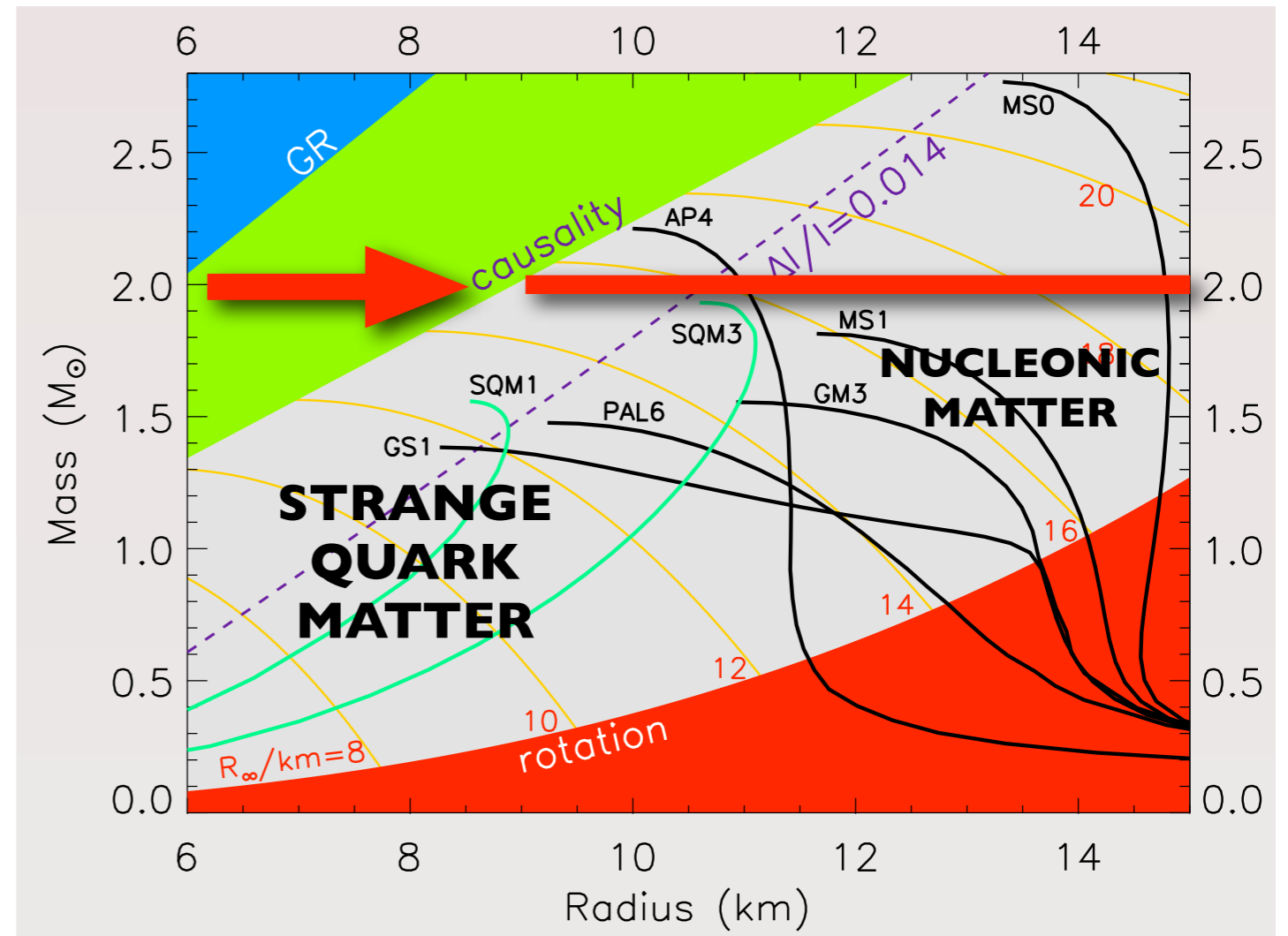
## ● Mass-Radius Relation



## Neutron Star Scenarios Tolman-Oppenheimer-Volkov Equations

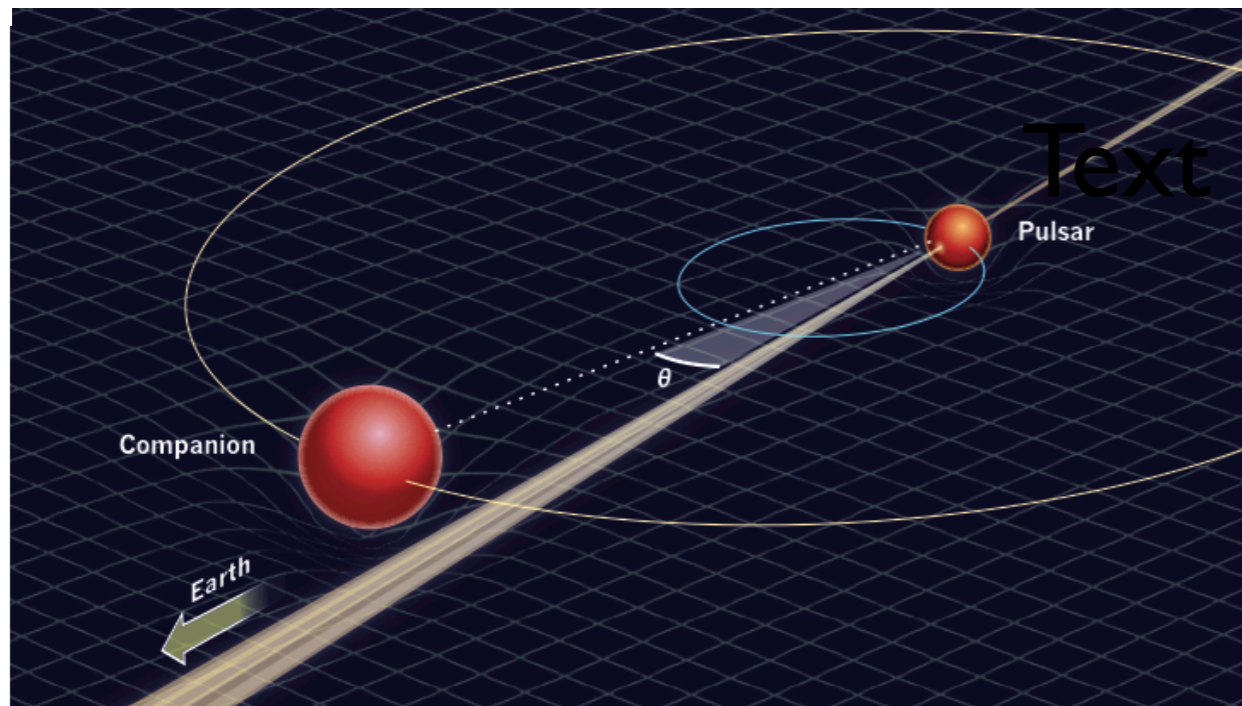
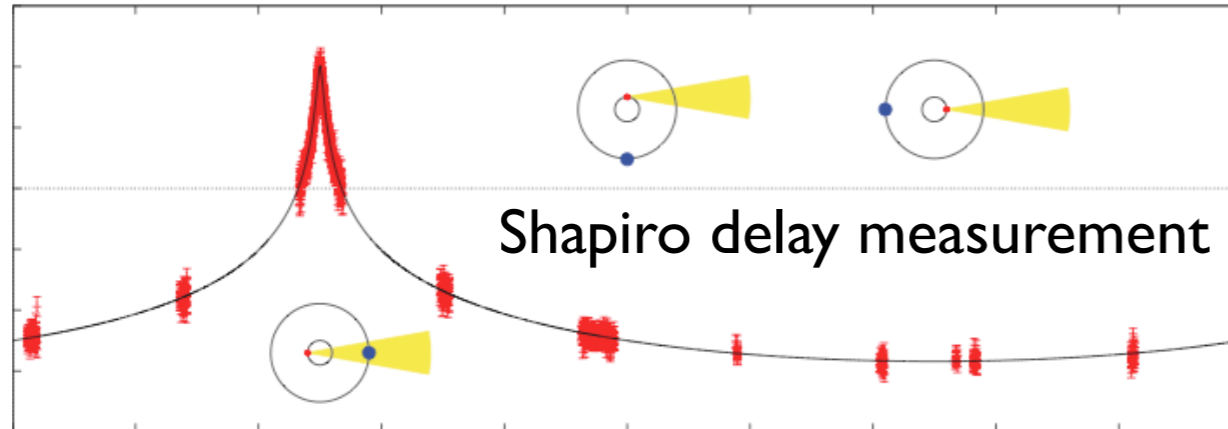
$$\frac{dP}{dr} = -\frac{G}{c^2} \frac{(M + 4\pi Pr^3)(\mathcal{E} + P)}{r(r - GM/c^2)}$$

$$\frac{dM}{dr} = 4\pi r^2 \frac{\mathcal{E}}{c^2}$$



# New constraints from NEUTRON STARS

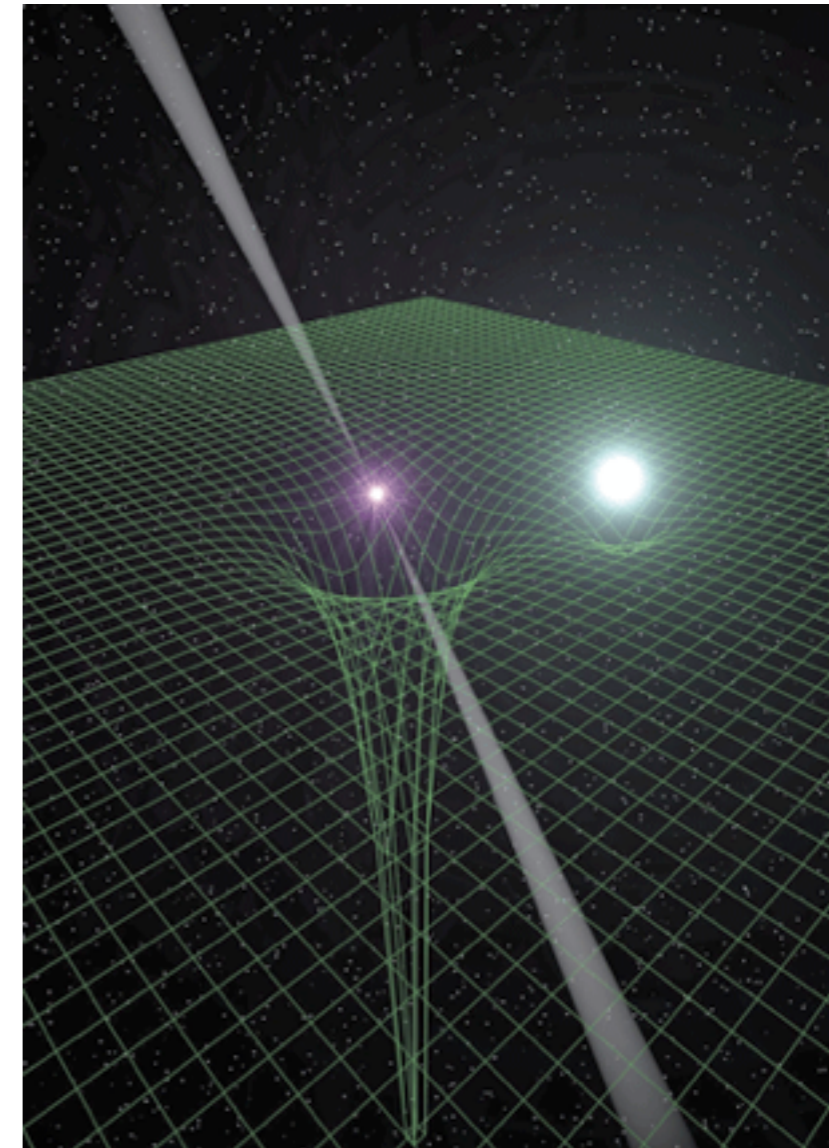
P.B. Demorest et al.  
Nature 467 (2010) 1081



PSR J1614+2230

$$M = 1.97 \pm 0.04 M_{\odot}$$

J. Antoniadis et al.  
Science 340 (2013) 6131



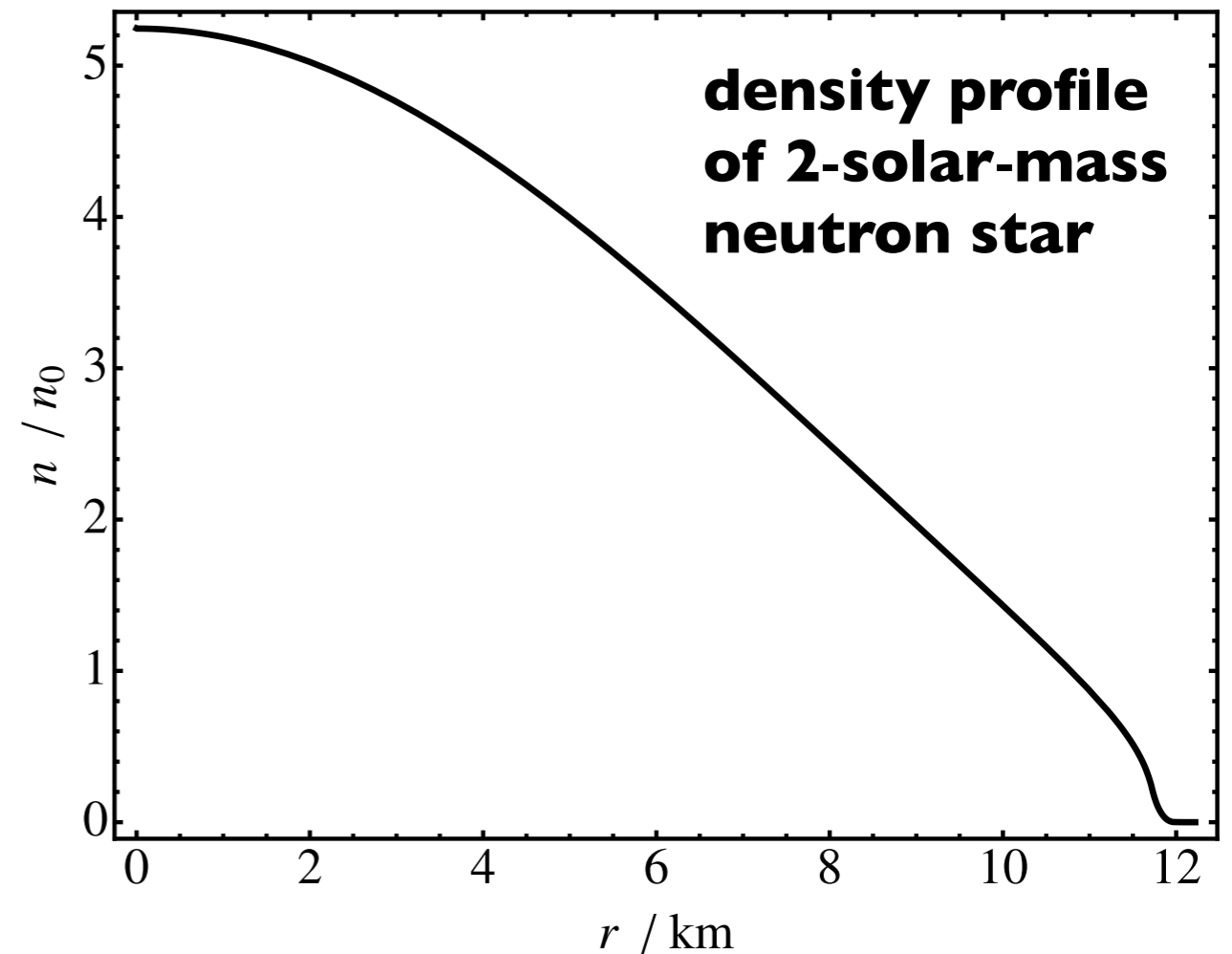
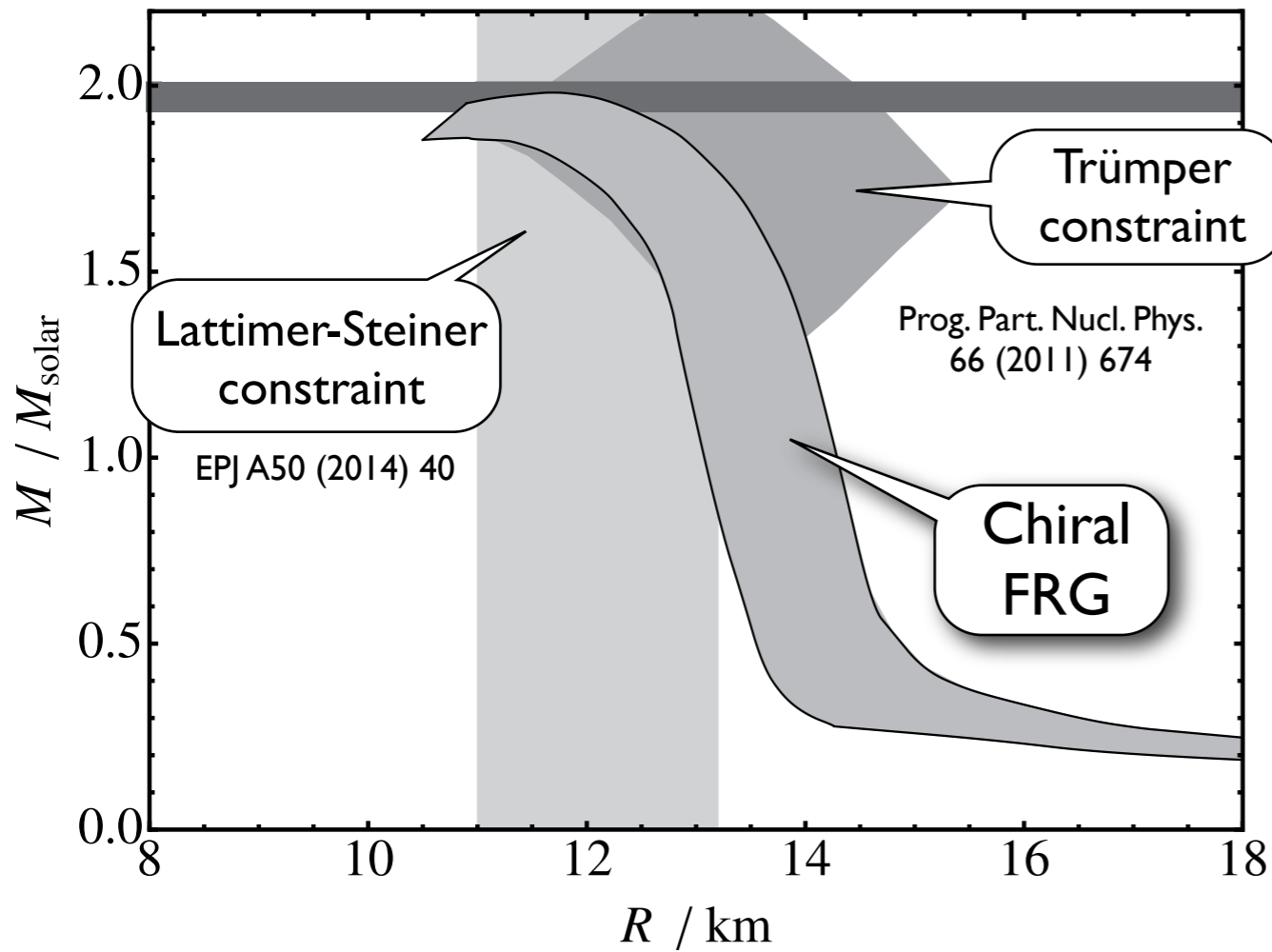
PSR J0348+0432

$$M = 2.01 \pm 0.04 M_{\odot}$$

# NEUTRON STAR MATTER

from  
**Chiral Nucleon-Meson Model**  
 and **FRG**

M. Drews, W.W. (2014)



- Neutron matter plus proton admixture (beta equilibrium)
- Symmetry energies 30 - 37 MeV
- **No** ultrahigh densities in the core of the neutron star

# CONCLUSIONS

- **FRG** provides non-perturbative approach to  
**Nuclear Chiral Thermodynamics**  
from symmetric to asymmetric nuclear matter and neutron matter
  - ▶ fluctuations beyond mean field include important two-pion exchange mechanisms and low-energy nucleonic particle-hole excitations
  - ▶ 1st order phase transition: Fermi liquid  $\leftrightarrow$  interacting Fermi gas
- **No** indication of **first-order chiral phase transition**
  - ▶ fluctuations work against early restoration of chiral symmetry
- New constraints from **neutron stars** for the **equation-of-state** of **dense & cold baryonic matter**:
  - ▶ Mass - radius relation: **stiff equation of state** required !  
**No** ultrahigh densities ( $\rho_{\max} \sim 5 \rho_0$ )
  - ▶ **Conventional** (nucleon-meson, “non-exotic”) **EoS** fulfills constraints

