

Spontaneous symmetry breaking in fermion systems with functional RG: Purely fermionic approach

Andreas Eberlein and Walter Metzner

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Why purely fermionic RG?

Introduction → Walter Metzner's talk

- Partially bosonized RG:
 - Relatively simple description of collective fluctuations
 - Ward identities can be implemented easily
 - More biased for simple truncations → Dynamical (re-) bosonization
- Purely fermionic RG:
 - Unbiased even for simple truncations
 - Accurate results for superfluid gap in ground state of attractive Hubbard model
 - Fermionic two-particle vertex at strong coupling contains contributions that cannot be described as boson exchange

Outline

- 1 Fermionic fRG for symmetry breaking
 - Katanin scheme
 - Example: Reduced BCS model
- 2 Application: Attractive Hubbard model
- 3 More recent developments
 - Other symmetry broken states
 - fRG+MF in 1PI formalism
- 4 Summary

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Fermionic flow equations: Katanin scheme

- Scheme proposed by Katanin in [PRB 70, 115109 \(2004\)](#)
- Aim: Derive truncation that captures ladder approximation *exactly*

$$\frac{d}{d\Lambda} \text{[shaded square]} = \text{[shaded square with bubble]} + \text{[shaded hexagon with bubble]}$$

$$\frac{d}{d\Lambda} \text{[shaded hexagon]} = \text{[shaded hexagon with three shaded squares]} + \text{[shaded square with bubble]} + \text{[shaded hexagon with bubble]}$$

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$$\frac{d}{d\Lambda} \text{[shaded hexagon]} = \text{[shaded hexagon with bubble]} + \text{[shaded square with bubble]} + \text{[shaded hexagon with bubble]} + \text{[shaded octagon with bubble]}$$

Neglect terms of $\mathcal{O}(\Gamma^4)$ or higher

Fermionic flow equations: Katanin scheme

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$$\frac{d}{d\Lambda} \text{[shaded square]} = \text{[shaded square with circle]} + \text{[shaded hexagon with circle]}$$

$$\frac{d}{d\Lambda} \text{[shaded hexagon]} = \text{[shaded hexagon with circle]} + O(\Gamma^4) = \frac{d}{d\Lambda} \left(\text{[shaded hexagon with circle]} \right) + O(\Gamma^4)$$

Right hand side can be rewritten as total derivative

Fermionic flow equations: Katanin scheme

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$$\frac{d}{d\Lambda} \text{[shaded square]} = \text{[shaded square with circle]} + \text{[shaded hexagon with circle]}$$

$$\text{[shaded hexagon]} = \text{[shaded hexagon with three shaded triangles]} + O(\Gamma^4)$$

Simple approximation for $\Gamma_{(3)} \rightarrow$ Insert into equation for Γ

Fermionic flow equations: Katanin scheme

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Two-loop flow equation for Γ :

$$\frac{d}{d\Lambda} \Gamma = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3}$$

Next step: Neglect terms with overlapping loops

Fermionic flow equations: Katanin scheme

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Two-loop flow equation for Γ :

$$\frac{d}{d\Lambda} \text{[shaded square]} = \text{[shaded bubble with top/bottom blocks]} + \text{[shaded bubble with top and bottom-side blocks]} + \text{[shaded bubble with top-side and bottom blocks]} \quad \text{(last term crossed out)}$$

Next step: Neglect terms with overlapping loops

Fermionic flow equations: Katanin scheme

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Two-loop flow equation for Γ :

$$\frac{d}{d\Lambda} \text{[shaded square]} = \text{[bubble with two shaded rectangles]} + \text{[shaded square on top, shaded trapezoid on bottom]} + \text{[shaded trapezoid on top, shaded square on bottom]} + \dots$$

Second term: $\dot{\Sigma}$ insertions for one-loop contribution
 Exploit $\dot{G}^\Lambda = S^\Lambda + G^\Lambda \dot{\Sigma}^\Lambda G^\Lambda$

Fermionic flow equations: Katanin scheme

- Scheme proposed by Katanin in [PRB 70, 115109 \(2004\)](#)
- Aim: Derive truncation that captures ladder approximation *exactly*

Two-loop flow equation for Γ :

$$\frac{d}{d\Lambda} \text{[shaded square]} = \text{[bubble with shaded blocks]} + \text{[shaded square on top arc, shaded circle on bottom arc]} + \text{[shaded square on bottom arc, shaded circle on top arc]} \approx \text{[shaded square on bottom arc, shaded circle on top arc]}$$

Second term: $\dot{\Sigma}$ insertions for one-loop contribution
 Exploit $\dot{G}^\Lambda = S^\Lambda + G^\Lambda \dot{\Sigma}^\Lambda G^\Lambda$

Fermionic flow equations: Katanin scheme

Result: Simple one-loop scheme (Katanin, PRB 70, 115109 (2004))

$$\begin{aligned}
 \frac{d}{d\Lambda} \text{ (circle) } &= \text{ (square with loop) } \\
 S^\Lambda = \frac{d}{d\Lambda} G^\Lambda |_{\Sigma^\Lambda = \text{const.}} & \\
 \frac{d}{d\Lambda} \text{ (square) } &= \text{ (two vertical rectangles) } - \text{ (two vertical rectangles with loop) } - \frac{1}{2} \text{ (two squares with loop) } \\
 & \quad \Pi^{\text{PH,d}} \quad \Pi^{\text{PH,cr}} \quad \Pi^{\text{PP}}
 \end{aligned}$$

- Standard truncation for instability analyses (See Metzner *et al.*, RMP 84, 299 (2012))
- Solves mean-field models exactly (Salmhofer *et al.*, PTP 112, 943 (2004))
- Yields reasonable results for SC gap in Hubbard model
(Gersch *et al.*, NJP 10, 045003 (2008); AE and Metzner, PRB 87, 174523 (2013), PRB 89, 035126 (2014))

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Reduced BCS model

- Microscopic action:

$$S[\bar{\psi}, \psi] = \sum_{k,\sigma} \bar{\psi}_{k\sigma} (-G_{0,k})^{-1} \psi_{k\sigma} - \Delta_0 \sum_k (\bar{\psi}_{k\uparrow} \bar{\psi}_{-k\downarrow} + \psi_{-k\downarrow} \psi_{k\uparrow}) \\ + g \sum_{k,k'} \bar{\psi}_{k\uparrow} \bar{\psi}_{-k\downarrow} \psi_{-k'\downarrow} \psi_{k'\uparrow}$$

- Model exactly solvable in thermodynamic limit
Solution: MF + RPA
- Exactly solvable with fRG using Katanin scheme

(Salmhofer *et al.*, PTP **112**, 943 (2004))

How to obtain exact solution with fRG?

Reduced BCS model

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How to obtain exact solution with fRG?

Reduced BCS model

- Ansatz for flowing action:

$$\begin{aligned}
 S^\Lambda[\bar{\psi}, \psi] = & \sum_{k,\sigma} \bar{\psi}_{k\sigma} (-G_{0,k}^\Lambda)^{-1} \psi_{k\sigma} - \Delta_\Lambda \sum_k (\bar{\psi}_{k\uparrow} \bar{\psi}_{-k\downarrow} + \psi_{-k\downarrow} \psi_{k\uparrow}) \\
 & + v_\Lambda \sum_{k,k'} \bar{\psi}_{k\uparrow} \bar{\psi}_{-k\downarrow} \psi_{-k'\downarrow} \psi_{k'\uparrow} \\
 & - \frac{w_\Lambda}{2} \sum_{k,k'} (\bar{\psi}_{k\uparrow} \bar{\psi}_{-k\downarrow} \bar{\psi}_{-k'\downarrow} \bar{\psi}_{k'\uparrow} + \text{conj.})
 \end{aligned}$$

- w_Λ important for symmetry breaking and fulfillment of Ward identity

- Flow of vertex: $\frac{d}{d\Lambda} \Gamma^\Lambda = \Pi_{\text{PH},d}^\Lambda + \Pi_{\text{PH},\sigma}^\Lambda + \Pi_{\text{PP}}^\Lambda$

- Solution: $\Gamma^\Lambda = \Gamma^{\Lambda_0} + \Gamma^{\Lambda_0} \mathcal{L}^\Lambda \Gamma^\Lambda$

Reduced BCS model

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- w_Λ important for symmetry breaking and fulfillment of Ward identity

- Flow of vertex: $\frac{d}{d\Lambda} \Gamma^\Lambda = \Pi_{\text{PH,d}}^\Lambda + \mathcal{O}((\beta V)^{-1})$

- Solution: $\Gamma^\Lambda = \Gamma^{\Lambda_0} + \Gamma^{\Lambda_0} \mathcal{L}^\Lambda \Gamma^\Lambda$

Reduced BCS model

- Ansatz for flowing action:

$$\begin{aligned}
 S^\Lambda[\bar{\psi}, \psi] = & \sum_{k,\sigma} \bar{\psi}_{k\sigma} (-G_{0,k}^\Lambda)^{-1} \psi_{k\sigma} - \Delta_\Lambda \sum_k (\bar{\psi}_{k\uparrow} \bar{\psi}_{-k\downarrow} + \psi_{-k\downarrow} \psi_{k\uparrow}) \\
 & + v_\Lambda \sum_{k,k'} \bar{\psi}_{k\uparrow} \bar{\psi}_{-k\downarrow} \psi_{-k'\downarrow} \psi_{k'\uparrow} \\
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Reduced BCS model

Solution for vertex: "Proof"

The diagram illustrates the proof for the vertex in the reduced BCS model. It shows the following steps:

$$\text{Shaded Square} = \text{Square with } \lambda_0 \text{ and dot} + \text{Shaded Rectangle}$$

$$\frac{d}{d\lambda} \text{Shaded Square} = \text{Shaded Rectangle with dot} + \text{Shaded Rectangle with dot}$$

$$= \text{Shaded Rectangle with dot} + \text{Shaded Rectangle with dot} = \text{Shaded Rectangle with dot}$$

The final result is a shaded rectangle with a dot, which is the derivative of the square vertex with respect to the coupling λ .

Reduced BCS model

Explicit solution for vertex

- Flow equations for coupling constants:

$$\dot{v}_\Lambda = (v_\Lambda^2 + w_\Lambda^2)\dot{B}^\Lambda + 2v_\Lambda w_\Lambda \dot{A}^\Lambda$$

$$\dot{w}_\Lambda = (v_\Lambda^2 + w_\Lambda^2)\dot{A}^\Lambda + 2v_\Lambda w_\Lambda \dot{B}^\Lambda$$

where $B^\Lambda = -\sum_p G^\Lambda(p)G^\Lambda(-p)$ and $A^\Lambda = \sum_p (F^\Lambda(p))^2$

- Introduction of amplitude and phase mode of superconducting gap:

$$a_\Lambda = v_\Lambda + w_\Lambda \qquad \varphi^\Lambda = v_\Lambda - w_\Lambda$$

$$\dot{a}_\Lambda = a_\Lambda^2 \frac{d}{d\Lambda} (B^\Lambda + A^\Lambda) \qquad \dot{\varphi}_\Lambda = \varphi_\Lambda^2 \frac{d}{d\Lambda} (B^\Lambda - A^\Lambda)$$

Simple solution:

$$a_\Lambda = \frac{g}{1 - g(B^\Lambda + A^\Lambda)}$$

$$\varphi_\Lambda = \frac{g}{1 - g(B^\Lambda - A^\Lambda)}$$

Reduced BCS model

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Reduced BCS model

Solution for self-energy

- Flow of self-energy: $\dot{\Delta}_\Lambda = -a_\Lambda \sum_k S_F^\Lambda(k)$

- Solution: $\Delta_\Lambda = -g \sum_k F(k)$

BCS gap equation! Can be seen by using solution for Γ^Λ and $\dot{G}^\Lambda = S^\Lambda + G^\Lambda \dot{\Sigma}^\Lambda G^\Lambda$

- Ward identity for global $U(1)$ charge symmetry:

$$\Delta_\Lambda = \Delta_0 + \Delta_0 \varphi_\Lambda (B^\Lambda - A^\Lambda) \quad \leftrightarrow \quad \varphi_\Lambda = g \frac{\Delta_\Lambda}{\Delta_0} \sim \Delta_0^{-1}$$

Reduced BCS model

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Reduced BCS model

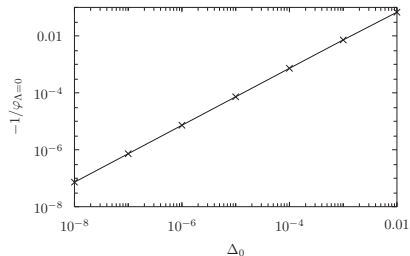
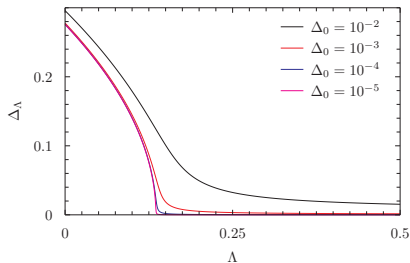
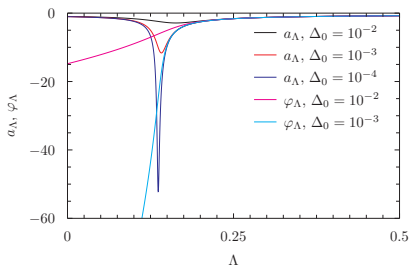
Solution for self-energy: "Proof"

$$\begin{aligned}
 \frac{d}{d\Lambda} \text{---}\bullet\text{---} &= \text{---}\text{---}\text{---}^{S^2} = \text{---}\text{---}\text{---} + \text{---}\text{---}\text{---} \\
 &= \text{---}\text{---}\text{---} + \text{---}\text{---}\text{---} = \frac{d}{d\Lambda} \text{---}\text{---}\text{---}
 \end{aligned}$$

Reduced BCS model

Numerical results

Constant density of states,
bandwidth $D = 2$,
momentum cutoff,
 $g = -0.5$



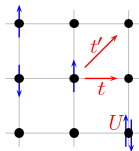
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Attractive Hubbard model

- Simple model for attractively interacting lattice fermions
- Can be simulated with cold atoms in optical lattices
- Hamiltonian:

$$\mathcal{H} = \sum_{i,j,\sigma} t_{ij} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

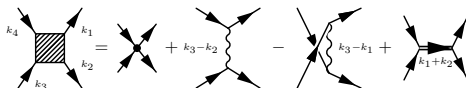


- $U < 0$, $d = 2$, $T = 0$
- Ground state: *s*-wave singlet superfluid;
Studied with fRG by Gersch *et al.*, NJP, **10**, 045003 (2008); AE and Metzner, PRB **87**, 174523 (2013); Strack *et al.*, PRB **78**, 014522 (2008); Obert *et al.*, PRB **88**, 144508 (2013)

Results presented here from AE and Metzner, PRB **87**, 174523 (2013)

Approximation for effective interactions and self-energies

- Vertex: Decomposition in interaction channels



Description of effective interactions: Boson exchange

$$V_{kk'}(q) = \sum_{\alpha, \beta} V_{\alpha\beta}(q) h_{\alpha}(q, k) h_{\beta}(q, k') \stackrel{U < 0}{\approx} V(q) g(k_0) g(k'_0)$$

(where V represents $A, \Phi, C, M, \nu, X, \tilde{X}$)

Extension of ideas by [Karrasch et al. JPCM 20, 345205 \(2008\)](#) and [Husemann and Salmhofer PRB 79, 195125 \(2009\)](#)

for description of symmetry breaking in Cooper channel

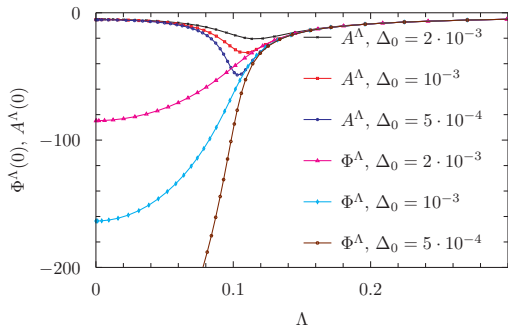
- Regulator and self-energy:

$$\text{Regulator: } ik_0 + R(k) = i \operatorname{sgn}(k_0) \sqrt{\Lambda^2 + k_0^2}$$

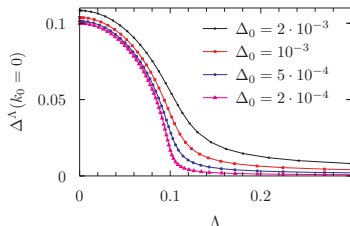
$$\text{Self-energy: } \Sigma(k) \approx \Sigma(k_0), \quad \Delta(k) \approx \Delta(k_0)$$

Amplitude and phase mode of superfluid gap

Scale dependence for $U = -2$, $t' = -0.1$, $n = 1/2$



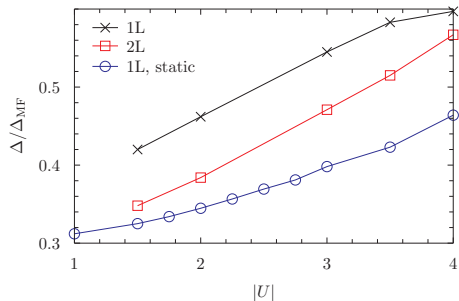
Flow of gap:



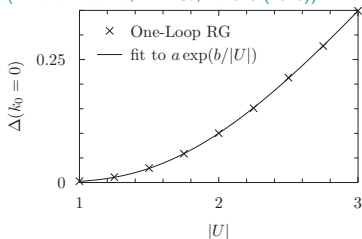
- Flow of couplings in Cooper channel similar to reduced BCS model

Reduction of superfluid gap by fluctuations

$$n = 1/2, t' = -0.1$$



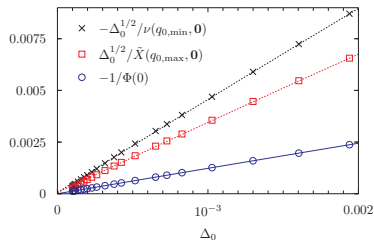
(AE and Metzner, PRB 87, 174523 (2013))



→ Good agreement with literature!

→ Well inside BCS regime

Δ/Δ_{MF}	Approximation, parameter
0.3	Perturbation theory, $n = 1/2$ or $n = 0.8$, $U \rightarrow 0$
0.5	One-loop fermionic RG, $n = 1/2$, $ U = 1 - 3.5$ (Gersch <i>et. al.</i> (2008))
0.39	DMFT, $n = 1/2$, $ U = 1.4$ (Bauer <i>et. al.</i> (2009))
0.76	One-loop mixed RG, $n = 1/2$, $ U = 4$ (Obert <i>et. al.</i> (2013))
0.5	QMC, $n = 0.2$, $ U = 4$ (Singer <i>et. al.</i> (1996))

Singular behaviour of vertex for $\Delta_0 \rightarrow 0$ ($U = -3, n = 1/2, t' = -0.1$)

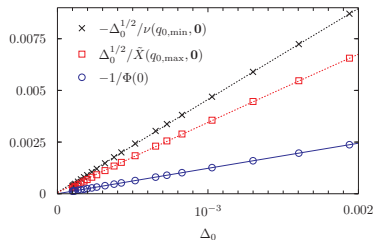
One-loop level:

$$\Phi(0, \mathbf{0}) \sim \Delta_0^{-1}$$

$$\nu(q_0 \sim \Delta_0^{1/2}, \mathbf{0}) \sim \Delta_0^{-1/2}$$

$$\tilde{X}(q_0 \sim \Delta_0^{1/2}, \mathbf{0}) \sim \Delta_0^{-1/2}$$

(AE and Metzner, PRB **87**, 174523 (2013))

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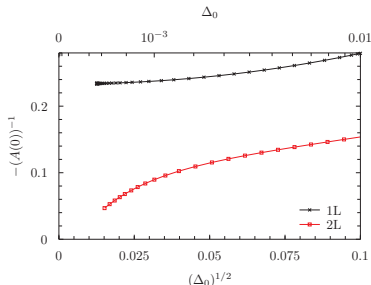
(AE and Metzner, PRB 87, 174523 (2013))

In addition at two-loop level:

(AE, PRB 90, 115125 (2014))

$$A(0) \sim \Delta_0^{-1/2}$$

$$\partial_\Lambda A^\Lambda(0) \sim \int \frac{d^3 p}{(2\pi)^3} \Phi^\Lambda(p) \partial_{\Lambda, V} \Phi^\Lambda(p)$$

Expected, see Strack *et al.*, PRB 78, 014522 (2008)No other singularities for $\Delta_0 \rightarrow 0$!

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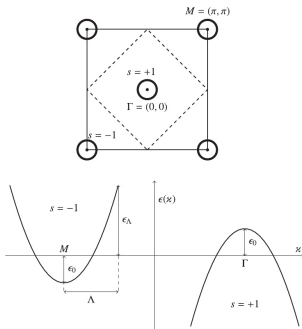
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Commensurate antiferromagnetism

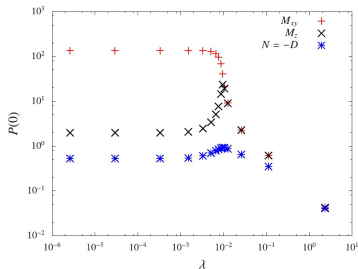
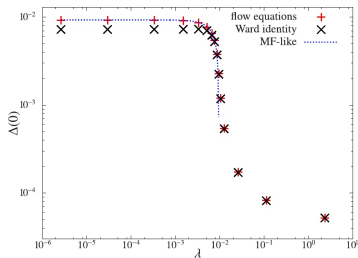
- Simple two-patch model with perfect nesting



- Translational and spin-rotational symmetries broken

(Maier and Honerkamp, PRB 86, 134404 (2012),

Maier et al., PRB 90, 035140 (2014))



d-wave superconductivity

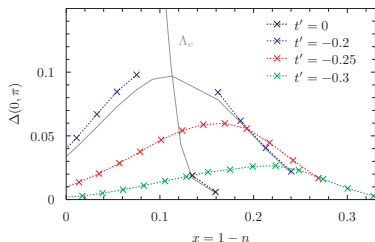
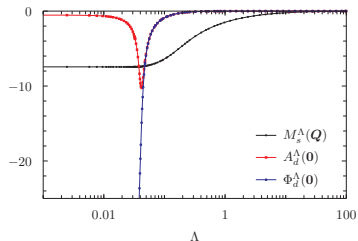
(AE and Metzner, PRB **89**, 035126 (2014))

Hubbard model:

$$\mathcal{H} = \sum_{i,j,\sigma} t_{ij} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

- Simple model for electronic structure of CuO₂ planes
- $U > 0$, $d = 2$, $T = 0$
- Ground state: AF, dSC, ...

$$\text{dSC: } \Delta(\mathbf{k}) \propto \cos k_x - \cos k_y$$



Extends instability analyses by Zanchi and Schulz (2000), Halboth and Metzner (2000), Honerkamp *et al.* (2001), Husemann and Salmhofer (2009), Husemann *et al.* (2012), Giering and Salmhofer (2012), ...

See also mixed fermion-boson RG study by Friederich *et al.*, PRB **83**, 155125 (2011)

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fRG+MF in 1PI formalism

Simple description of **competing** order? Combine fRG and MF theory!

fRG+MF in Wick-ordered formalism: [Reiss et al., PRB 75, 075110 \(2007\)](#)

Idea in 1PI formalism: ([Wang, AE, Metzner, PRB 89, 121116 \(2014\)](#))

- Solve full fRG flow including fluctuations for $\Lambda > \Lambda_{MF}$

$$\partial_\Lambda \Gamma^\Lambda = \Pi_{PH,d}^\Lambda - \Pi_{PH,cr}^\Lambda - \frac{1}{2} \Pi_{PP}^\Lambda \quad \partial_\Lambda \Sigma^\Lambda = \Gamma^\Lambda S^\Lambda$$

- Continue fRG flow for $\Lambda < \Lambda_{MF}$ and in symmetry broken phase below Λ_c for reduced model

$$\partial_\Lambda \Gamma^\Lambda = \Pi_{PH,d}^\Lambda = \Gamma^\Lambda \dot{L}^\Lambda \Gamma^\Lambda \quad \partial_\Lambda \Sigma^\Lambda = \Gamma^\Lambda S^\Lambda$$

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$$\Gamma^\Lambda = \tilde{\Gamma}^{\Lambda_{MF}} + \tilde{\Gamma}^{\Lambda_{MF}} L^\Lambda \Gamma^\Lambda \quad \Sigma^\Lambda = \Sigma^{\Lambda_{MF}} + \tilde{\Gamma}^{\Lambda_{MF}} (G^\Lambda - G^{\Lambda_{MF}})$$

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Recipe: (Wang, AE, Metzner, PRB **89**, 121116 (2014))

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$$\Sigma^{\Lambda=0} = \Sigma^{\Lambda_{\text{MF}}} + \tilde{\Gamma}^{\Lambda_{\text{MF}}} (G^{\Lambda=0} - G^{\Lambda_{\text{MF}}})$$

Remarks:

- No anomalous vertices and propagators appear in fRG flow
- Full fRG flow below Λ_c mean-field-like at least for SC at $T = 0$

(AE and Metzner, PRB **87**, 174523 (2013), PRB **89**, 035126 (2014); Obert *et al.*, PRB **88**, 144508 (2013))

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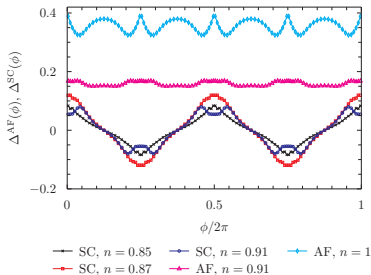
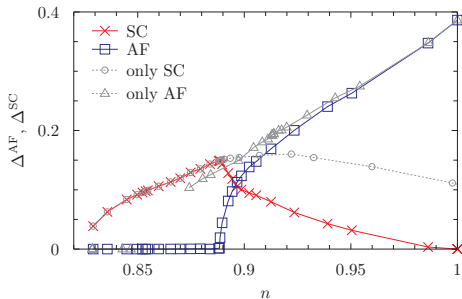
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Competition and coexistence of AF and dSC

Hubbard model, $U = 3$, $t' = -0.15$

So far: Only dSC and commensurate AF (Wang, AE, and WM, PRB 89, 121116(R) (2014))



- AF and dSC compete and coexist in sizeable doping range, differently from [Friederich *et al.*, PRB 83, 155125 \(2011\)](#)
- Quantum critical point between dSC and dSC+AF
- Note: Incommensurate AF important for these parameters

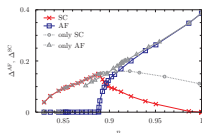
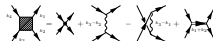
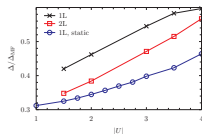
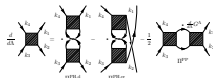
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- Yields reasonable results for gaps in various models
- Decomposition of vertex in interaction channels allows to capture its singular dependences on momentum and frequency
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→ fRG+MF in 1PI formalism as simple alternative



The End

Thank you for your attention!

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