

Anomalous scaling at non-thermal fixed points of Gross-Pitaevskii and KPZ turbulence

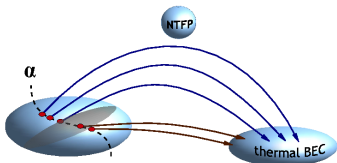
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Non-thermal fixed points

Non-thermal fixed point are far from equilibrium quasi stationary states of matter.

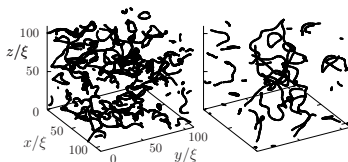
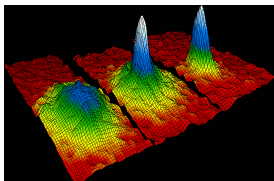


$$\text{Scale invariance } \epsilon(k) \sim k^{-d+\eta}$$

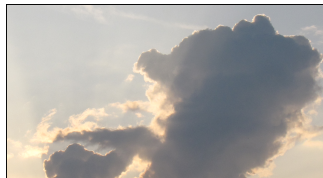
Depending on the initial conditions, the system may take an **algebraically long** time on the way to thermalisation.

Plan

Ultra-cold Bose gases,
Gross-Pitaevskii Equation (GPE)



Interface dynamics,
Kardar-Parisi-Zhang (KPZ)



Driven-Dissipative GPE

Classical field equation for the average Bose wave-function $\phi(\mathbf{x}, t)$:

$$i\partial_t\phi(\mathbf{x}, t) = \left[-\left(\frac{1}{2m}\right) \nabla^2 + \mu + g|\phi(\mathbf{x}, t)|^2 \right] \phi(\mathbf{x}, t)$$

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Classical field equation for the average Bose wave-function $\phi(\mathbf{x}, t)$:

$$i\partial_t\phi(\mathbf{x}, t) = \left[- \left(\frac{1}{2m} - i\nu \right) \nabla^2 + \mu + g|\phi(\mathbf{x}, t)|^2 \right] \phi(\mathbf{x}, t) + \zeta(\mathbf{x}, t)$$

With complex parameters

$$\mu = \mu_1 + i\mu_2$$

Single particle pump

$$g = g_1 - ig_2$$

2 particle losses

and stochastic driving

$$\langle \zeta(\mathbf{x}, t) \rangle = 0 \quad \langle \zeta(\mathbf{x}, t) \zeta(\mathbf{x}', t') \rangle = \gamma \delta(t - t') \delta(\mathbf{x} - \mathbf{x}')$$

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We focus on the kinetic energy density,

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d	1	2	3
η_{num}	1	small	small

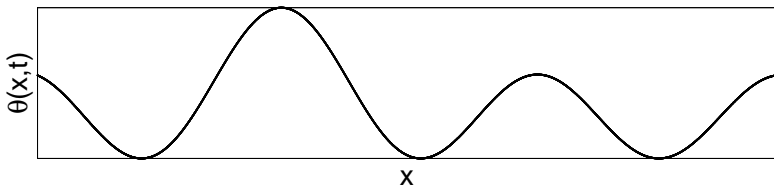
Kardar–Parisi–Zhang equation

A model for interface growth,

$$\partial_t \theta(\mathbf{x}, t) = \nu \nabla^2 \theta(\mathbf{x}, t) + \frac{\lambda}{2} [\nabla \theta(\mathbf{x}, t)]^2 + \eta(\mathbf{x}, t)$$

with diffusion, perpendicular expansion and stochastic driving,

$$\langle \eta(\mathbf{x}, t) \rangle = 0 \quad \langle \eta(\mathbf{x}, t) \eta(\mathbf{x}', t') \rangle = D \delta(t - t') \delta(\mathbf{x} - \mathbf{x}')$$



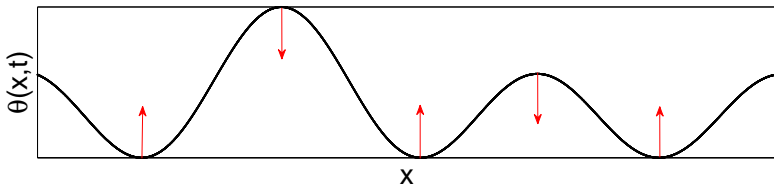
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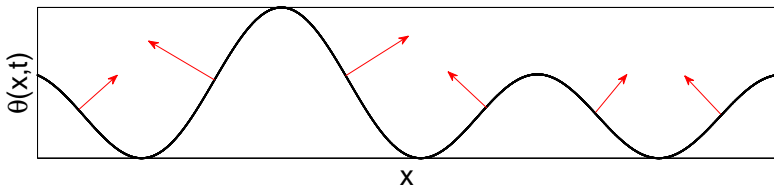
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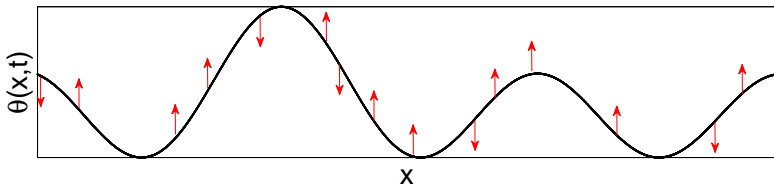
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Scaling in interface growth

The stationary state has scaling correlation functions,

$$\langle \theta(t + \tau, \mathbf{x} + \mathbf{r}) \theta(t, \mathbf{x}) \rangle_c = r^{2\chi} g(\tau/r^z)$$

with exponents given by:

d	1	2	3	4
χ	1/2	0.379	0.300	?
$z = 2 - \chi$	3/2	1.6210	1.700	?

From DDGPE to KPZ

Density and phase decomposition $\phi(\mathbf{x}, t) = \sqrt{n(\mathbf{x}, t)} e^{-i\theta(\mathbf{x}, t)}$

$$\partial_t \theta - \frac{1}{2m} (\nabla \theta)^2 - \nu \nabla^2 \theta = U,$$

$$\partial_t n - \frac{1}{m} \nabla \cdot (n \nabla \theta) = S,$$

with sources of phase and density fluctuations,

$$U = U[\theta, n] + \frac{\text{Re}(\zeta e^{i\theta})}{\sqrt{n}}$$

$$S = S[\theta, n] + 2\sqrt{n} \text{Im}(\zeta e^{i\theta}).$$

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Comparing scaling exponents

$$\begin{aligned}\epsilon_{\text{kin}} &= \frac{1}{2m} \langle |\nabla \phi|^2 \rangle & \epsilon_{\text{kin}}(k) &\sim k^{-d+\eta} \\ &= \frac{n}{2m} \int_{\mathbf{k}} k^2 \langle |\theta(\mathbf{k}, t)|^2 \rangle & \epsilon_{\text{kin}}(k) &\sim k^{z-d-\chi}\end{aligned}$$

$$\rightarrow \boxed{\eta = z - \chi}$$

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GPE simulations

d	1	2	3
η_{num}	1	small	small

KPZ literature

d	1	2	3
η	1	1.242	1.400

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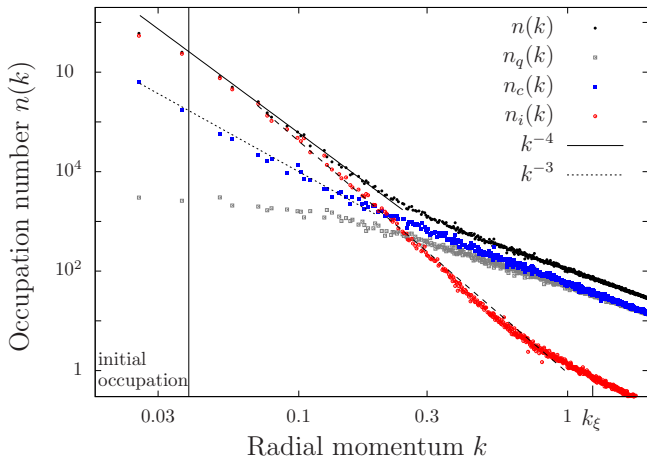
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$$d = 2$$

$$\epsilon(k) = k^2 n(k)$$

$$\nu = \mu_2 = 0$$

$$g_2 = \zeta = 0$$



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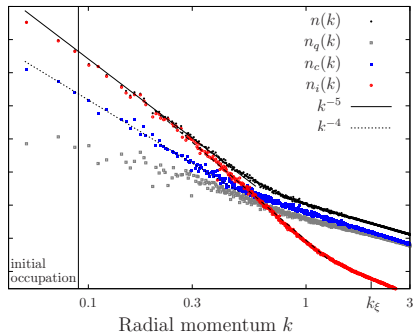
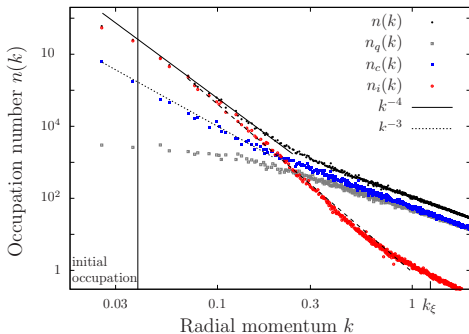
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$d = 2$

$d = 3$



Compressible excitations

$$\begin{aligned} \epsilon_c &= \frac{1}{2m} \langle |\nabla \phi|^2 \rangle & \epsilon_c(k) &\sim k^{-d+\eta+1} \\ &= \frac{n}{2m} \int_{\mathbf{k}} k^2 \langle |\theta(\mathbf{k}, t)|^2 \rangle & \epsilon_c(k) &\sim k^{z-d-\chi} \end{aligned} \quad \rightarrow \quad \boxed{\eta = z - \chi - 1}$$

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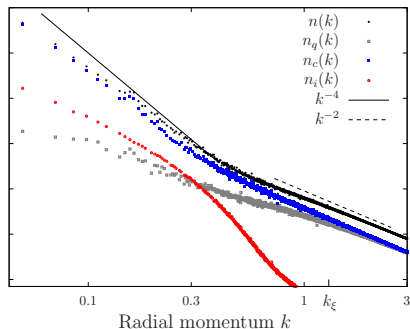
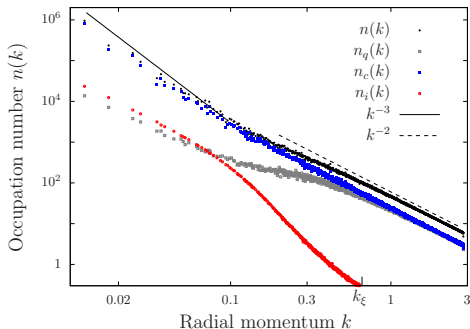
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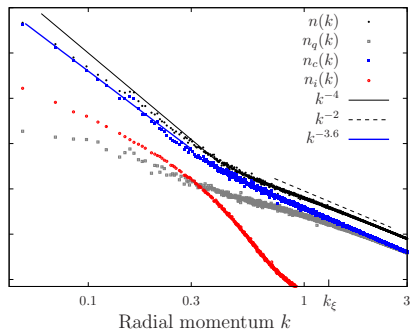
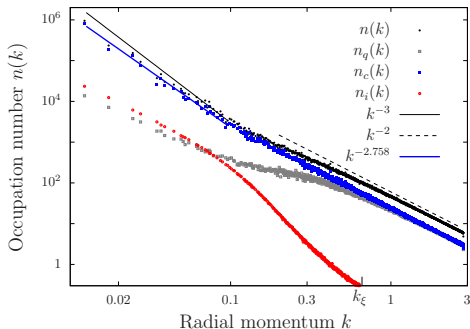
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 $d = 2$
 $d = 3$


Conclusions

- Interface dynamics described by KPZ equation does **not** capture **vortex** dynamics.
- It does captures **the rest**.
- We have made an estimation of **anomalous scaling exponents** of the ultra-cold Bose gas at a non-thermal fixed point.

Gasenzer, SM, Pawłowski, arXiv:1405.7652 [cond-mat.quant-gas]