

Phases of Unitary Imbalanced Fermi Gases

Dietrich Roscher

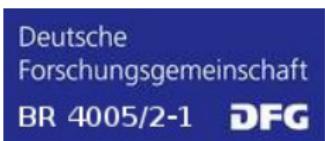
[I. Boettcher, J. Braun, T.K. Herbst, J.M. Pawłowski, D. Roscher, C. Wetterich, arXiv:1409.5070]

[D. Roscher, J. Braun, J.E. Drut, arXiv:1410.xxxx in preparation]

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Institut für Kernphysik / Technische Universität Darmstadt

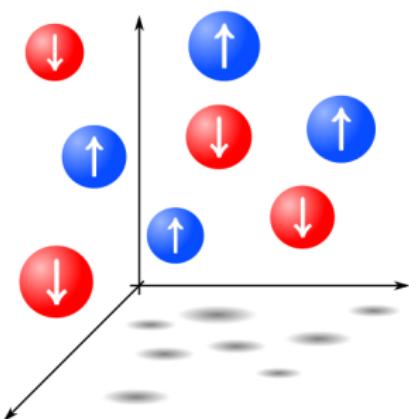
September 24, 2014



Ultracold Fermi Gases:

$$S = \int_0^{\frac{1}{T}} \int d^3x \sum_{\sigma=\uparrow,\downarrow} \psi_{\sigma}^{*} \left(\partial_{\tau} - \frac{\nabla^2}{2m} - \mu \right) \psi_{\sigma}$$

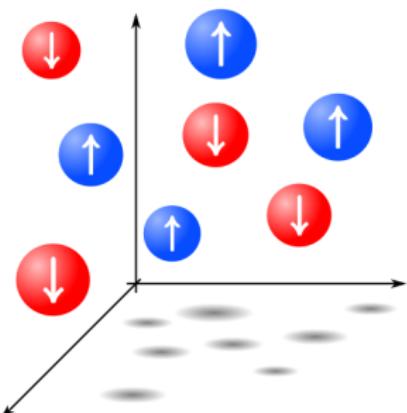
- Spatial dimensions: 3
- Two components: $\sigma = \uparrow, \downarrow$



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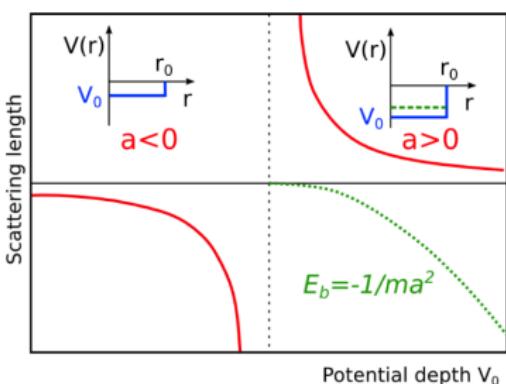
- Spatial dimensions: 3
- Two components: $\sigma = \uparrow, \downarrow$
- System temperature $T \sim T_F = \frac{1}{2m} (6\pi^2 n)^{\frac{1}{3}} \sim \mathcal{O}(\mu K)$



Unitary Fermi Gases:

$$S = \int_0^{\frac{1}{\bar{T}}} \int d^3x \left[\sum_{\sigma=\uparrow,\downarrow} \psi_{\sigma}^{*} \left(\partial_{\tau} - \frac{\nabla^2}{2m} - \mu \right) \psi_{\sigma} + \frac{\bar{g}}{2} \psi_{\uparrow}^{*} \psi_{\uparrow}^{*} \psi_{\downarrow} \psi_{\downarrow} \right]$$

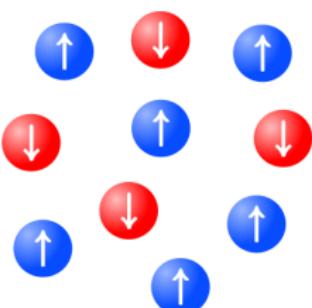
- Contact interaction: effective range $r_e = 0$
- Scattering length a tunable due to Feshbach resonance in experiment
- Here: *unitary limit* $\frac{1}{a} \rightarrow 0$ (strong coupling)



Unitary Imbalanced Fermi Gases:

$$S = \int_0^{\frac{1}{T}} \int d^3x \left[\sum_{\sigma=\uparrow,\downarrow} \psi_{\sigma}^{*} \left(\partial_{\tau} - \frac{\nabla^2}{2m} - \mu_{\sigma} \right) \psi_{\sigma} + \frac{\bar{g}}{2} \psi_{\uparrow}^{*} \psi_{\downarrow}^{*} \psi_{\downarrow} \psi_{\uparrow} \right]$$

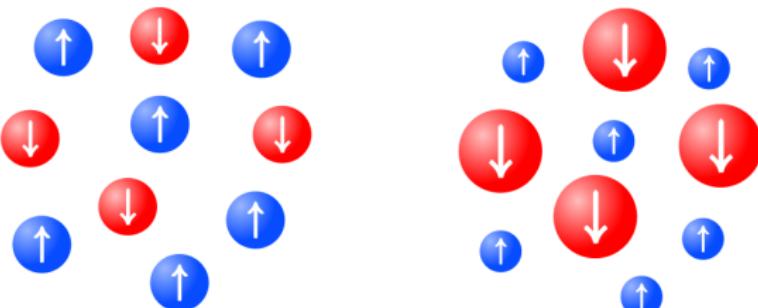
- Population/Spin imbalance: $h = \frac{\mu_{\uparrow} - \mu_{\downarrow}}{2}$



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- Population/Spin imbalance: $h = \frac{\mu_\uparrow - \mu_\downarrow}{2}$
- Mass imbalance $\bar{m} = \frac{m_\downarrow - m_\uparrow}{m_\uparrow + m_\downarrow}$
- Set reduced mass $\frac{4m_\uparrow m_\downarrow}{m_\uparrow + m_\downarrow} = 1$ and average chemical potential $\frac{\mu_\uparrow + \mu_\downarrow}{2} = 1$ for convenience



Unitary Imbalanced Fermi Gases

To be explored:

- Spontaneous breaking of $U(1)$ gauge & galilean invariance
- Various quantum and thermodynamic properties

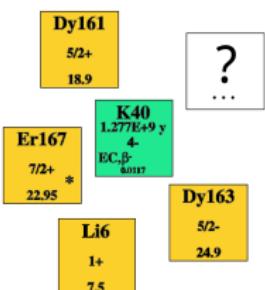
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Physical relevance:

- Upcoming cold atoms experiments



[Isotopes explored by *Grimm* group, Innsbruck]

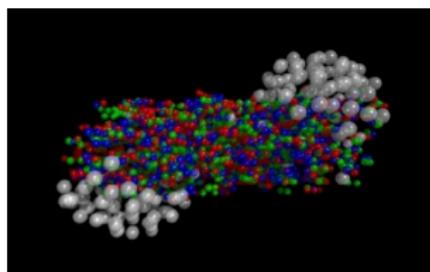
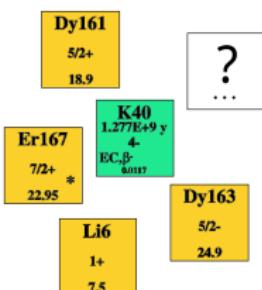
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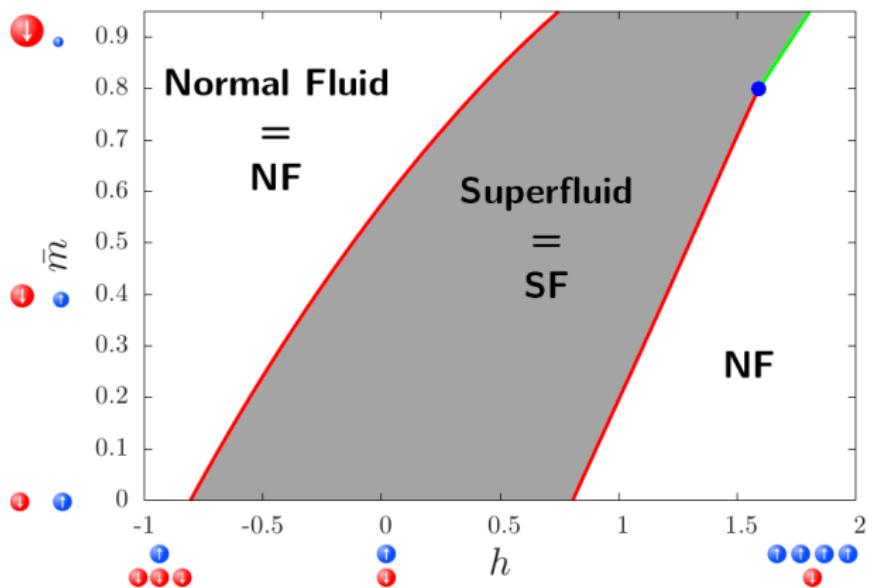
- Upcoming cold atoms experiments
- General insight: symmetry breaking patterns of fermionic field theories [e.g. V. Schoen, M. Thies '00]
- „Simulation“ of more complicated systems (heavy ion collisions, color superconductivity) [e.g. A. Adams, L.D. Carr, T. Schaefer, P. Steinberg, J. Thomas '12]



[Isotopes explored by Grimm group, Innsbruck]

First Impressions: Mean Field at $T = 0$

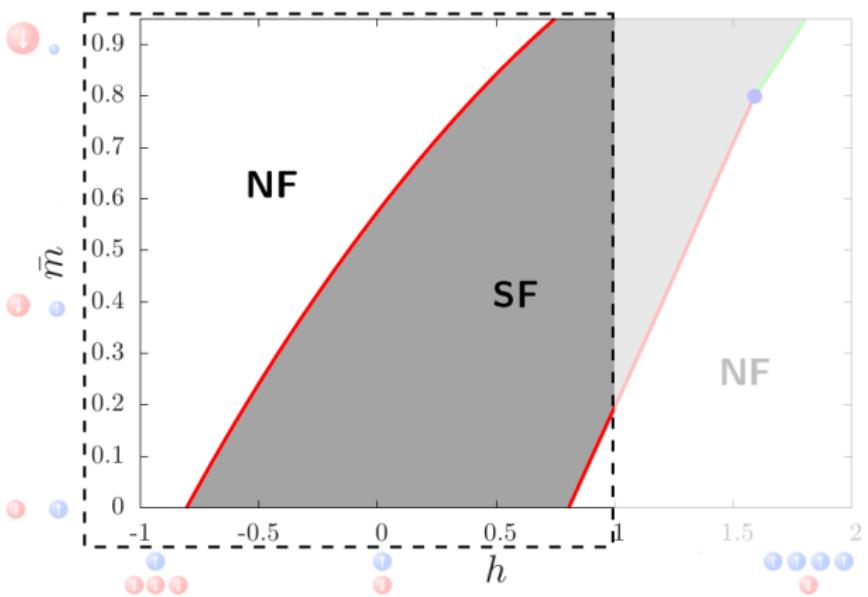
$$U_{\text{MF}}(\Delta_0) \sim \int dq q^2 \left\{ \frac{|\Delta_0|^2}{2q^2} - T \sum_{\sigma=\pm 1} \ln \left[1 + e^{-\frac{1}{T} (q^2 \bar{m} - h + \sigma \sqrt{(q^2 - 1)^2 + |\Delta_0|^2})} \right] \right\}$$



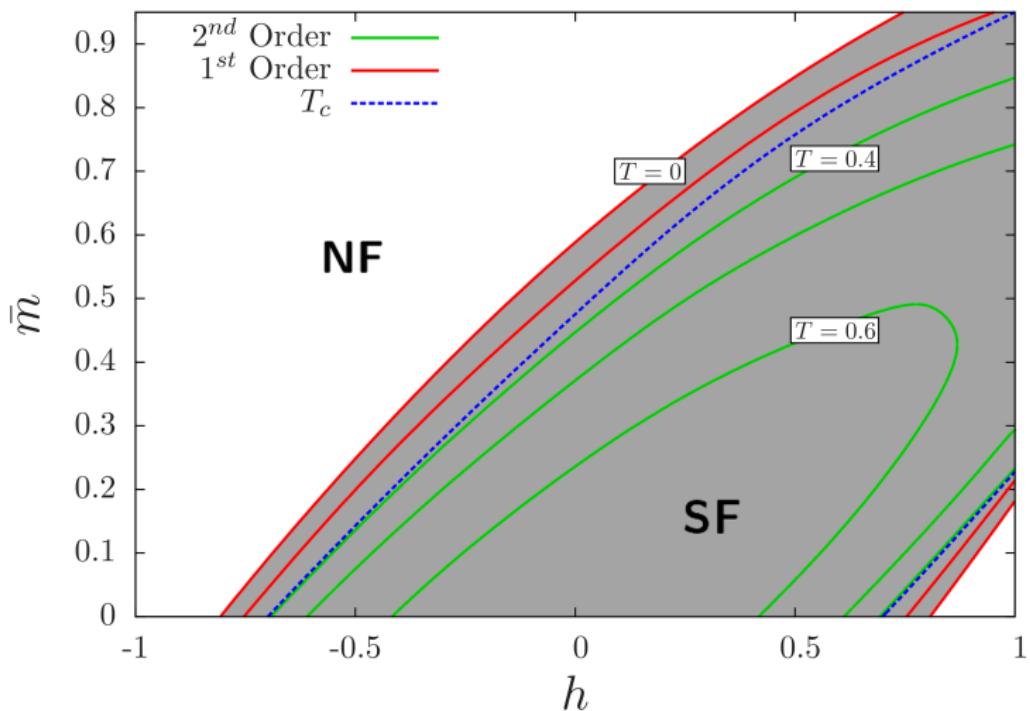
[M.M. Parish, F.M. Marchetti, A. Lamacraft, D.B. Simons '07]

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Mean Field: Finite Temperature



Examining Δ_0 More Closely

$$U_{\text{MF}}(\Delta_0) \sim \int dq q^2 \left\{ \frac{|\Delta_0|^2}{2q^2} - T \sum_{\sigma=\pm 1} \ln \left[1 + e^{-\frac{1}{T} (q^2 \bar{m} - h + \sigma \sqrt{(q^2 - 1)^2 + |\Delta_0|^2})} \right] \right\}$$

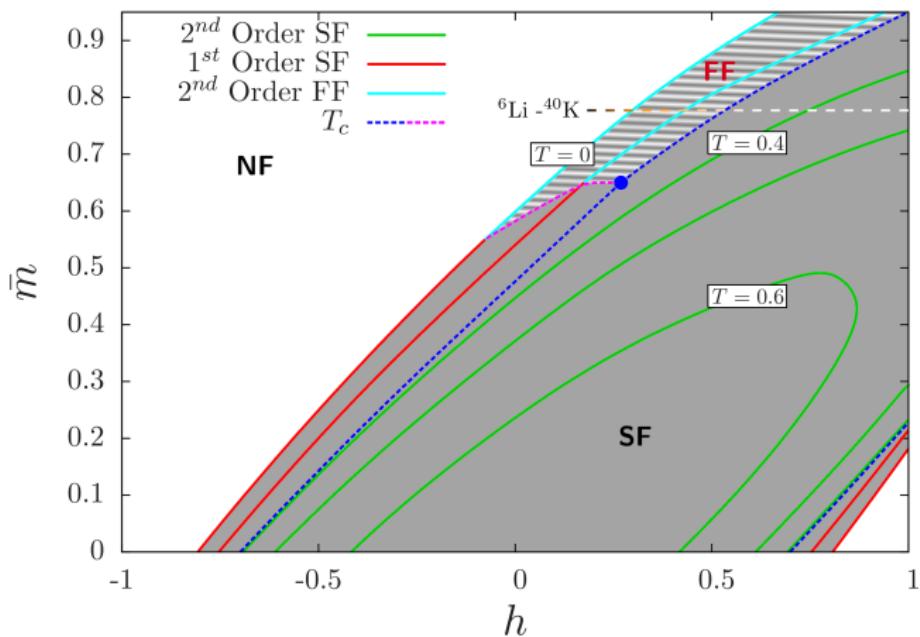
- Implicit assumption so far: $\Delta_0 \neq \Delta_0(\vec{x})$
- Analytical calculations for general $\Delta_0(\vec{x})$ non-trivial even at mean field [see e.g. M. Thies et al '06 ; G. Basar, G.V. Dunne '08]

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- FF-Ansatz: $\Delta_0(\vec{x}) \equiv \Delta_0 e^{i2\vec{Q} \cdot \vec{x}}$ [P. Fulde, R.A. Ferrell '64; A.I. Larkin, Y.N. Ovchinnikov '64]

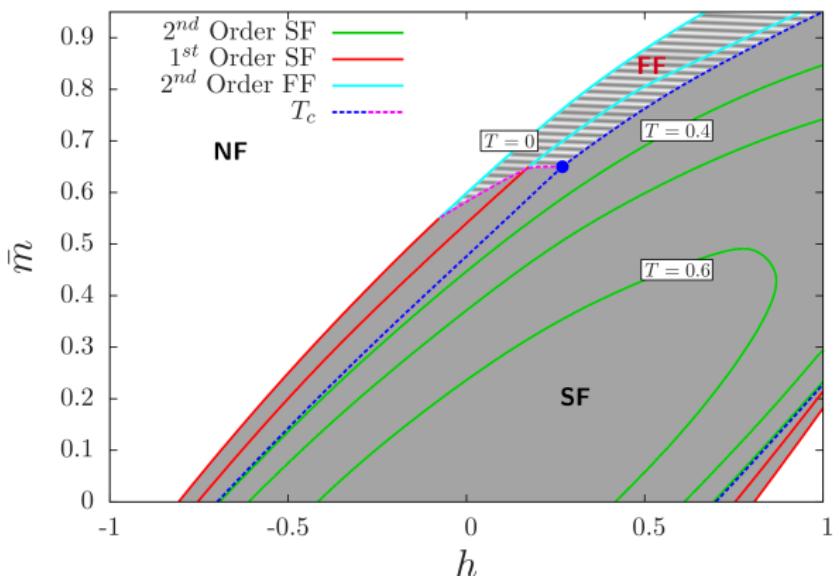
Mean Field Phase Diagram Including FF Region



[For ${}^6\text{Li}-{}^{40}\text{K}$: J.E. Baarsma, H.T.C. Stoof '13]

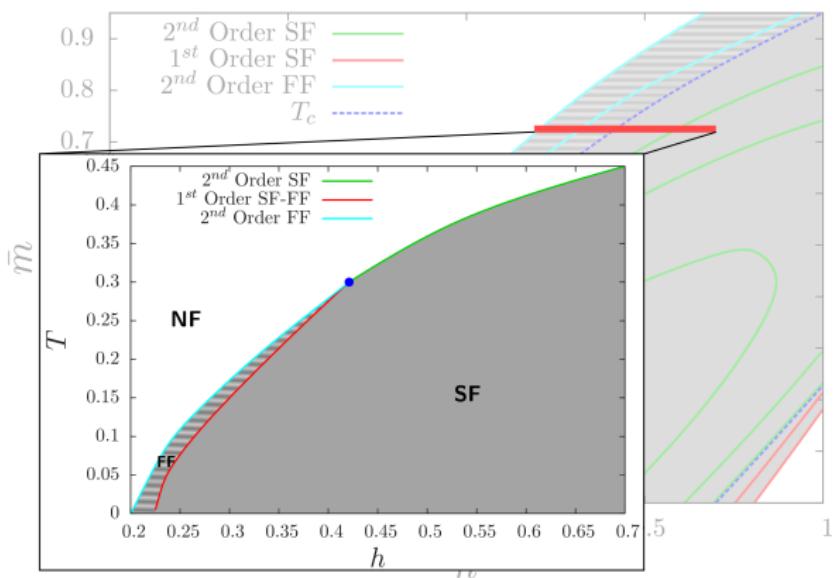
Spontaneous Breaking of Galilean Invariance!

...and beyond Mean Field?



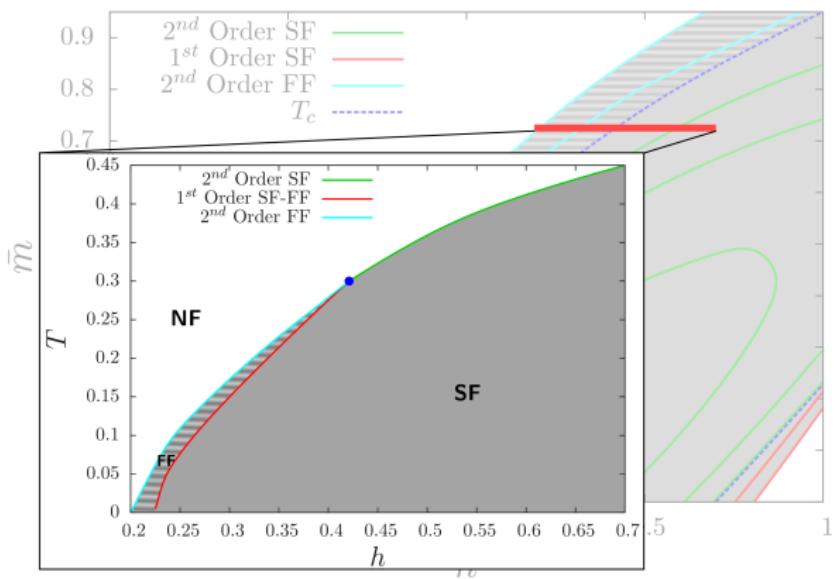
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⇒ FRG Analysis

Based on framework by: [S. Diehl et al '07, '10; S. Floerchinger et al '08; ...]

FRG Implementation

$$\partial_t \Gamma_k = \frac{1}{2} S \text{Tr} \left[\left(\Gamma_k^{(2)} + R_k \right)^{-1} \cdot (\partial_t R_k) \right] \quad [\text{Wetterich '93}]$$

Ansatz for Γ_k :

$$\begin{aligned} \Gamma_k = \int_0^{\frac{1}{T}} \int d^3x \left[\sum_{\sigma=\uparrow\downarrow} \psi_\sigma^* (\partial_\tau - \Delta(1 + \sigma \bar{m}) - \mu(1 + \sigma h)) \psi_\sigma + \varphi^* \left(\frac{Z_\varphi}{A_\varphi} \partial_\tau - \frac{1 - \bar{m}^2}{2} \Delta \right) \varphi \right. \\ \left. + U_k(\rho = \varphi^* \varphi) - \frac{h_\varphi}{2} (\varphi^* \psi_\uparrow \psi_\downarrow - \varphi \psi_\uparrow^* \psi_\downarrow^*) \right] \end{aligned}$$

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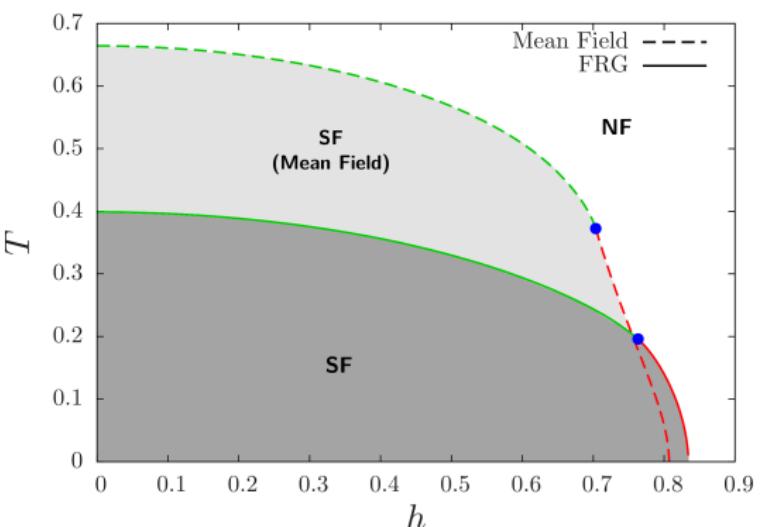
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- Flow of the effective potential:

$$k \partial_k U_k = \eta_{A_\varphi} \rho U'_k + [k \partial_k U_k]^\psi + [k \partial_k U_k]^\varphi$$

- Order parameter of $U(1)$ symmetry breaking: $\Delta_{0,k} = h_\varphi \sqrt{\rho_{0,k}}$, $U_k(\rho_{0,k})$ minimum
- Numerical implementation: discretization of $U_k(\rho)$ on a grid

First Results & Benchmarking: $\bar{m} = 0$ case

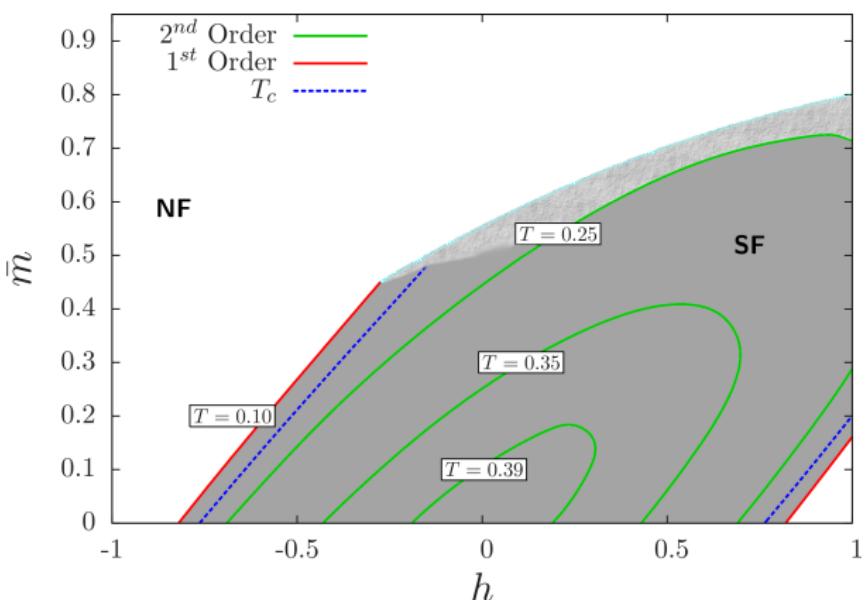


[I. Boettcher, J. Braun, T.K. Herbst, J.M. Pawłowski, D. Roscher, C. Wetterich '14]

See also: talk by Tina Herbst

- For small h : $T_{c,\text{FRG}} < T_{c,\text{MF}}$
- Endpoint of the transition line at $T = 0$: $h_{c,\text{FRG}} > h_{c,\text{MF}}$
- Quantitative agreement for two different regularization schemes

FRG Phase Diagram with Mass Imbalance



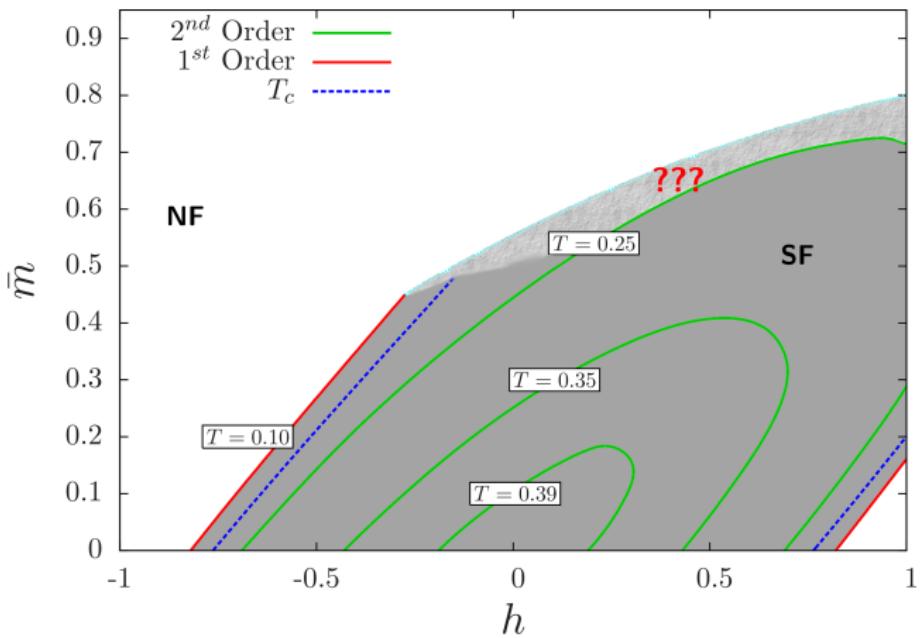
[D. Roscher, J. Braun, J.E. Drut *in prep.*]

Similar as for the $\bar{m} = 0$ case:

- For small h and \bar{m} : $T_{c,\text{FRG}} < T_{c,\text{MF}}$
- Superfluid region slightly broadened for low T

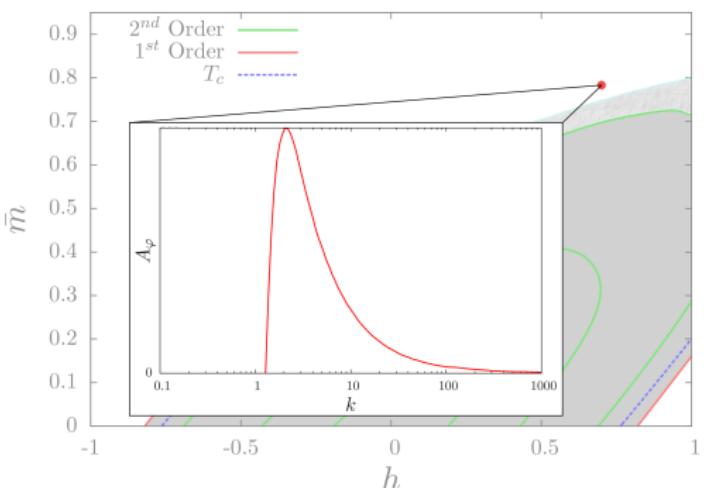
Full Phase Diagram with Mass Imbalance

Convergence Problems?



Behaviour of the Boson Propagator

$$P_\varphi^{-1}(q_0, \vec{q}) = iZ_\varphi q_0 + \frac{1-\bar{m}^2}{2} A_\varphi \vec{p}^2$$



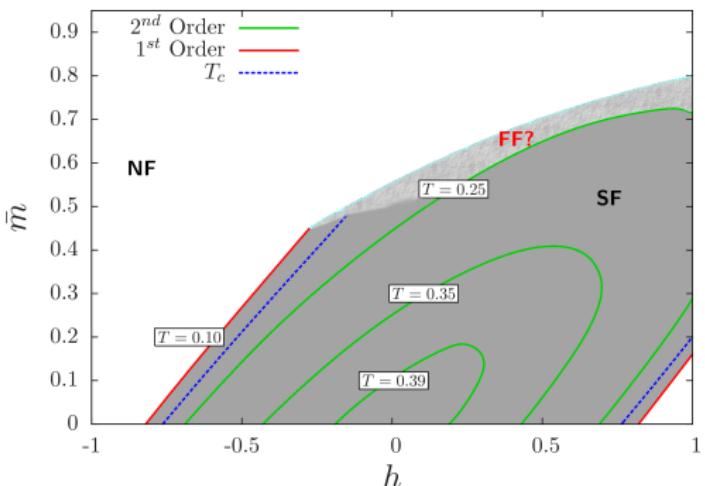
[D. Roscher, J. Braun, J.E. Drut *in prep.*]

- Sign change of A_φ invalidates lowest order deriv. expansion for P_φ^{-1}
- Interpretation: formation of bosons with momentum $|\vec{Q}| > 0$ favored

[H.C. Krahl, S. Friederich, C. Wetterich '09]

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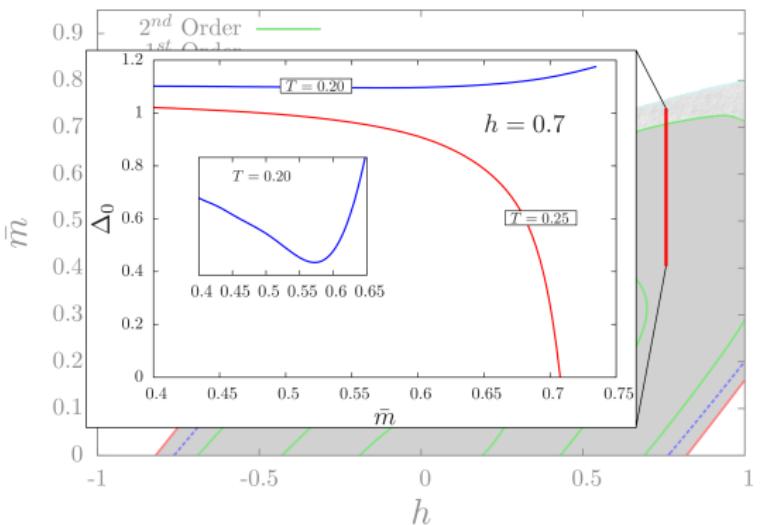
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 - Interpretation: formation of bosons with momentum $|\vec{Q}| > 0$ favored
- [H.C. Krahl, S. Friederich, C. Wetterich '09]
- Rough coincidence with FF region at mean field

Further Precursors of Inhomogeneous Pairing

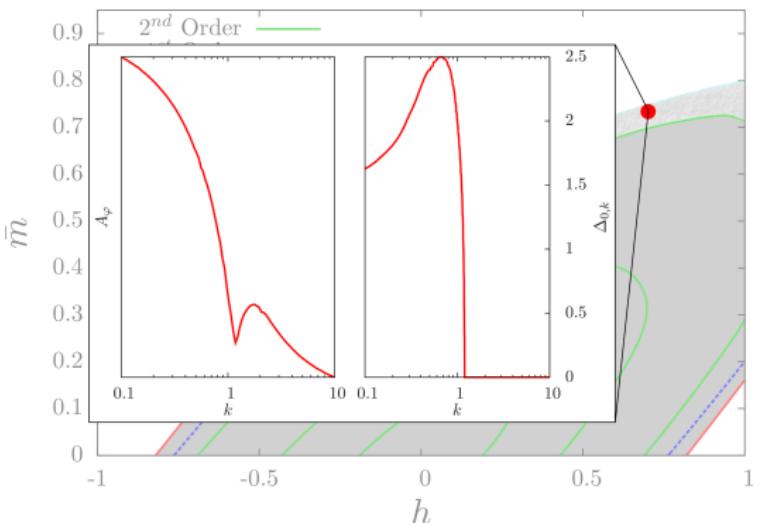


[D. Roscher, J. Braun, J.E. Drut *in prep.*]

Inside the problematic region:

- Unphysical growing of Δ_0 towards transition to normal phase

Further Precursors of Inhomogeneous Pairing



[D. Roscher, J. Braun, J.E. Drut *in prep.*]

Inside the problematic region:

- Unphysical growing of Δ_0 towards transition to normal phase
- Early occurrence of $\Delta_{0,k} > 0$ inhibits sign change of A_φ
- Conclusion: $\Delta_0 \neq \Delta_0(\vec{x})$ probably not the true ground state

Summary

- Application of FRG to compute the phase diagram of a unitary Fermi gas with spin and mass imbalance beyond mean field
- Observations: T_c mostly lowered, superfluid phase slightly extended
- Precursors of inhomogeneous phases found, qualitative agreement with mean field predictions

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Outlook

- More elaborate derivative expansion/explicit inhomogenous condensate → w.i.p. with J. Braun, S. Rechenberger
- Full BEC-BCS crossover: finite a , dimensions $d \neq 3$
→ planned with I. Boettcher, T. Herbst
- Complementary MC simulations → w.i.p. with J. Braun, J. Drut

[D. Roscher, J. Braun, J.-W. Chen, J. Drut '13; J. Braun, J. Drut, D. Roscher '14]