

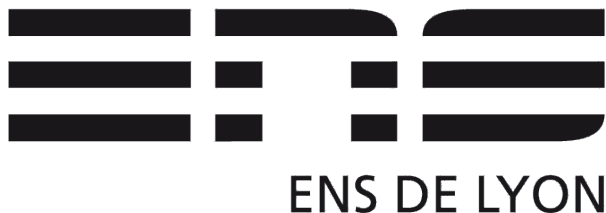
Fate of the Higgs mode in the vicinity of a quantum critical point

Adam Rançon

ENS-Lyon

Nicolas Dupuis

LPTMC (Paris)



ERG 2014



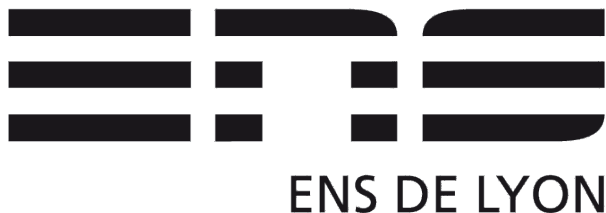
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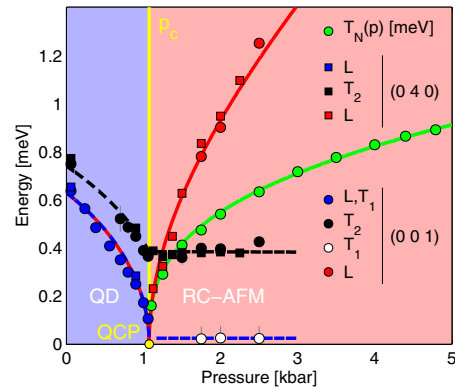


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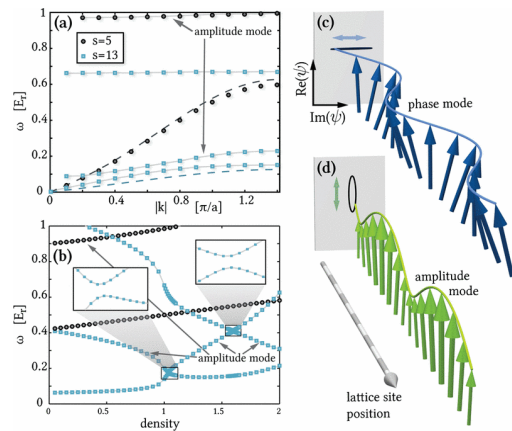


Experimental observations

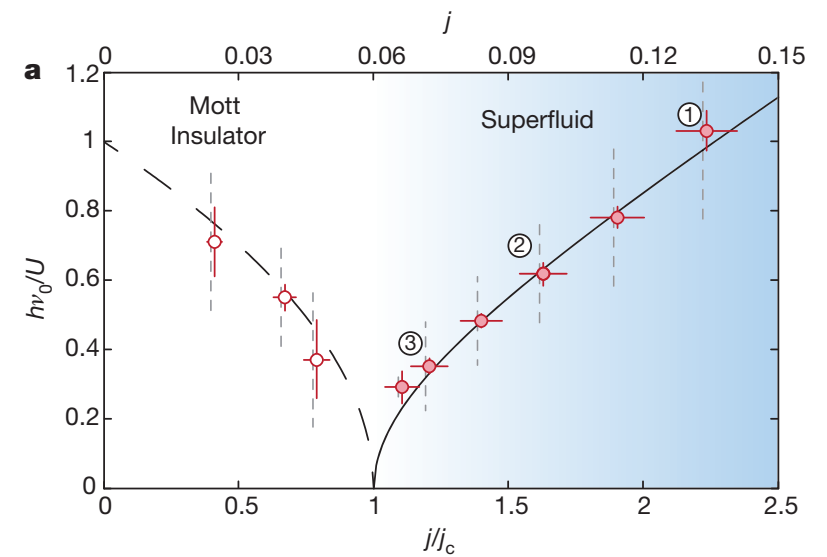
Ruegg et al. PRL 2008
Quantum spins - 3D



Blissbort et al. PRL 2008
Cold atoms - 3D

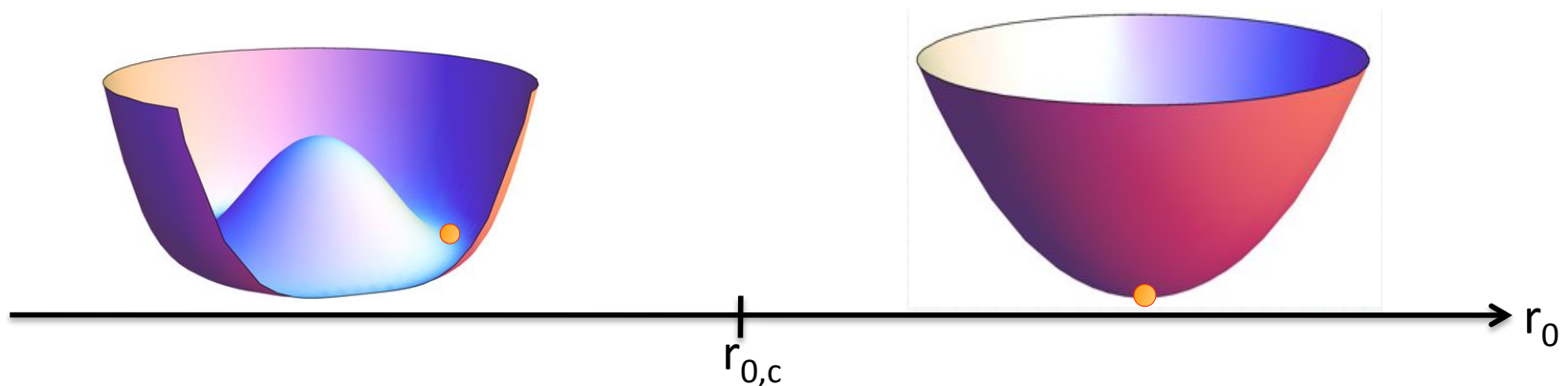


Enders et al. Nature 2012
Cold atoms - 2D



Quantum O(N) model

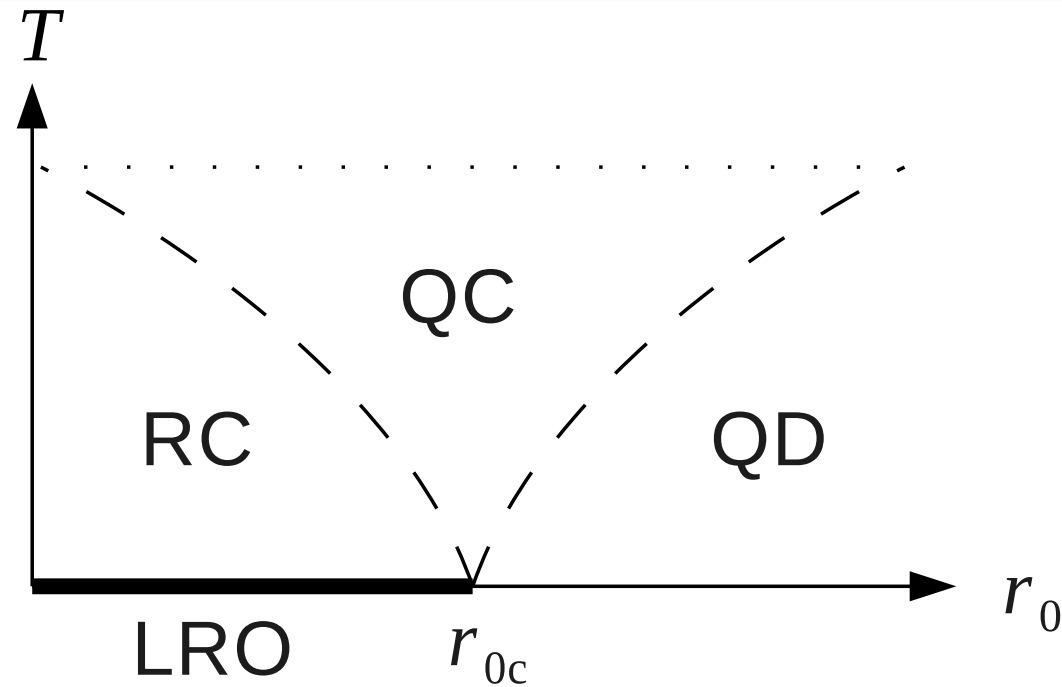
$$S = \int_0^\beta d\tau \int d^d x \left\{ \frac{1}{2} \psi_i (-c^{-2} \partial_\tau^2 - \nabla^2 + r_0) \psi_i + \frac{u_0}{4!N} (\psi_i^2)^2 \right\}$$



- Generalization of classical O(N) model
- at $T=0$, Lorentz symmetry
- Describes critical regime of a number of systems :
 - bosons in optical lattices
 - antiferromagnets
 - Josephson junction arrays
 - granular superconductors, ...

Chakravarty et al. 1989
Chubukov et al. 1994
Sachdev's book, etc.

Typical phase diagram in 2D



Quantum critical point described by the classical $O(N)$ model in 2+1 D

Non trivial critical exponent : η, ν

One (quantum) exponent : dynamical exponent $z = 1$ (Lorentz symmetry)

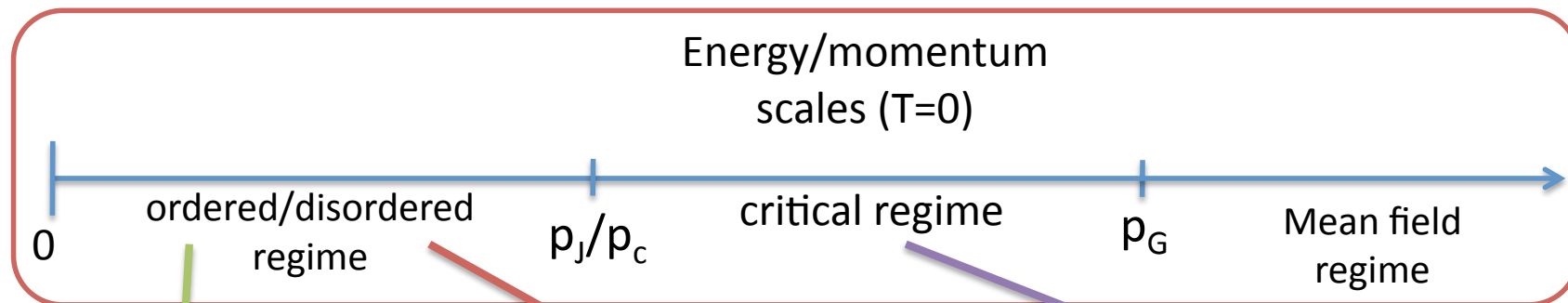
Vanishing energy scale close to the critical point (gap in the disordered phase)

$$\Delta \propto |r_0 - r_{0c}|^{z\nu}$$

Excitation spectrum in ordered phase ?

-Critical behavior and phase diagram well-known

-Excitation spectrum ? Much less clear in the ordered phase



$$G_t(p) = \frac{1}{p^2}$$

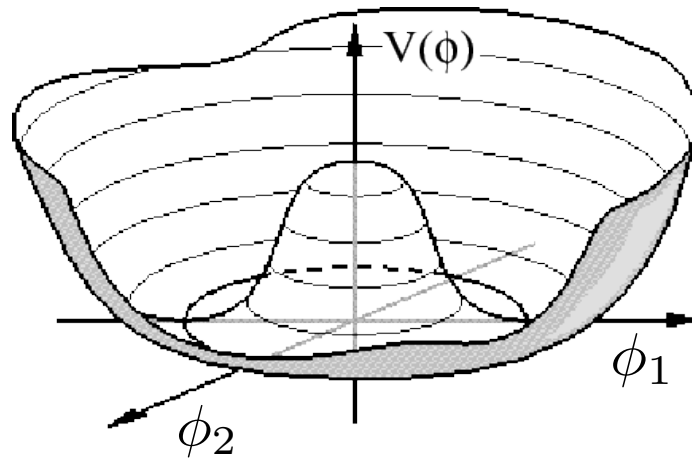
$$G_l(p) = ?$$

$$G(p) = \frac{1}{p^2 + \Delta^2}$$

$$G(p) = \frac{1}{p^{2-\eta}}$$

$$p^2 = \omega_n^2 + \mathbf{p}^2$$

Excitation spectrum : Mean Field



Broken symmetry phase $\langle \phi_i \rangle = \delta_{i1} \phi_0$

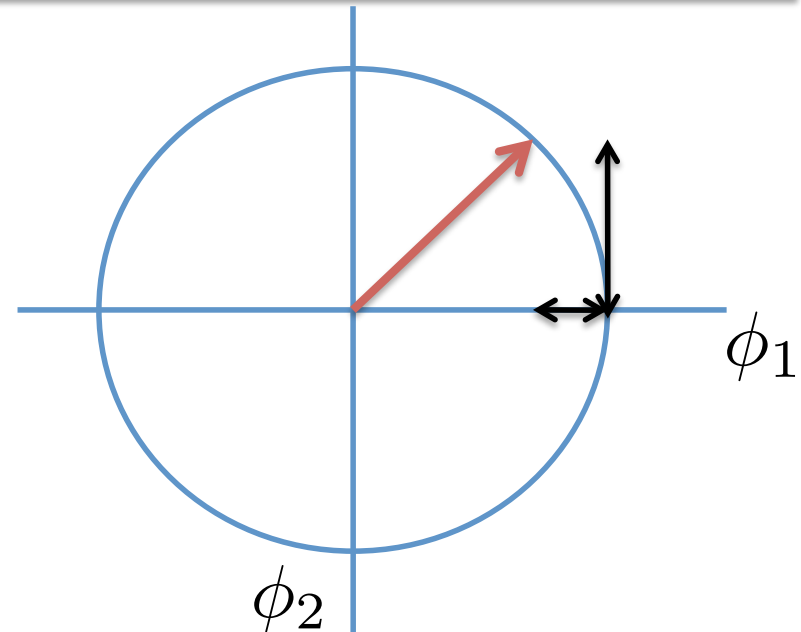
Mean-field picture : - N-1 transverse (massless) Goldstone modes $G_t(p) = \frac{1}{p^2}$
- 1 gapped longitudinal mode $G_l(p) = \frac{1}{p^2 + 2\Delta^2}$

Amplitude fluctuations = longitudinal fluctuations (in MF !)

MF spectrum has one gapped mode arbitrarily close to the QCP ($\Delta \rightarrow 0$)
at energy $\omega = \sqrt{2}\Delta$

Spectrum : Goldstone regime

At low energy (dominated by the Goldstones), amplitude $\langle |\phi| \rangle$ constant.



Strong coupling between longitudinal and transverse fluctuations

$$G_l(p) \propto \frac{1}{p} \text{ at low energy}$$

Longitudinal (amplitude ?) fluctuations seem not to be a well defined mode close to criticality:

the $1/p$ divergence might hide the resonance as Δ goes to zero.

Amplitude mode : scalar fluctuations

Podolsky et al. 2011 : it depends on the correlation function !

$$\phi = \sqrt{\rho} \mathbf{n} \quad G_\rho = \langle \rho \rho \rangle$$

magnitude



Not the same correlation function !

Example for antiferromagnets:

- longitudinal = neutron scattering
- scalar (amplitude) = Raman scattering

At low energy, coupling for decay of an amplitude fluctuation into two Goldstone is small (derivative coupling – proportional to ω^2).

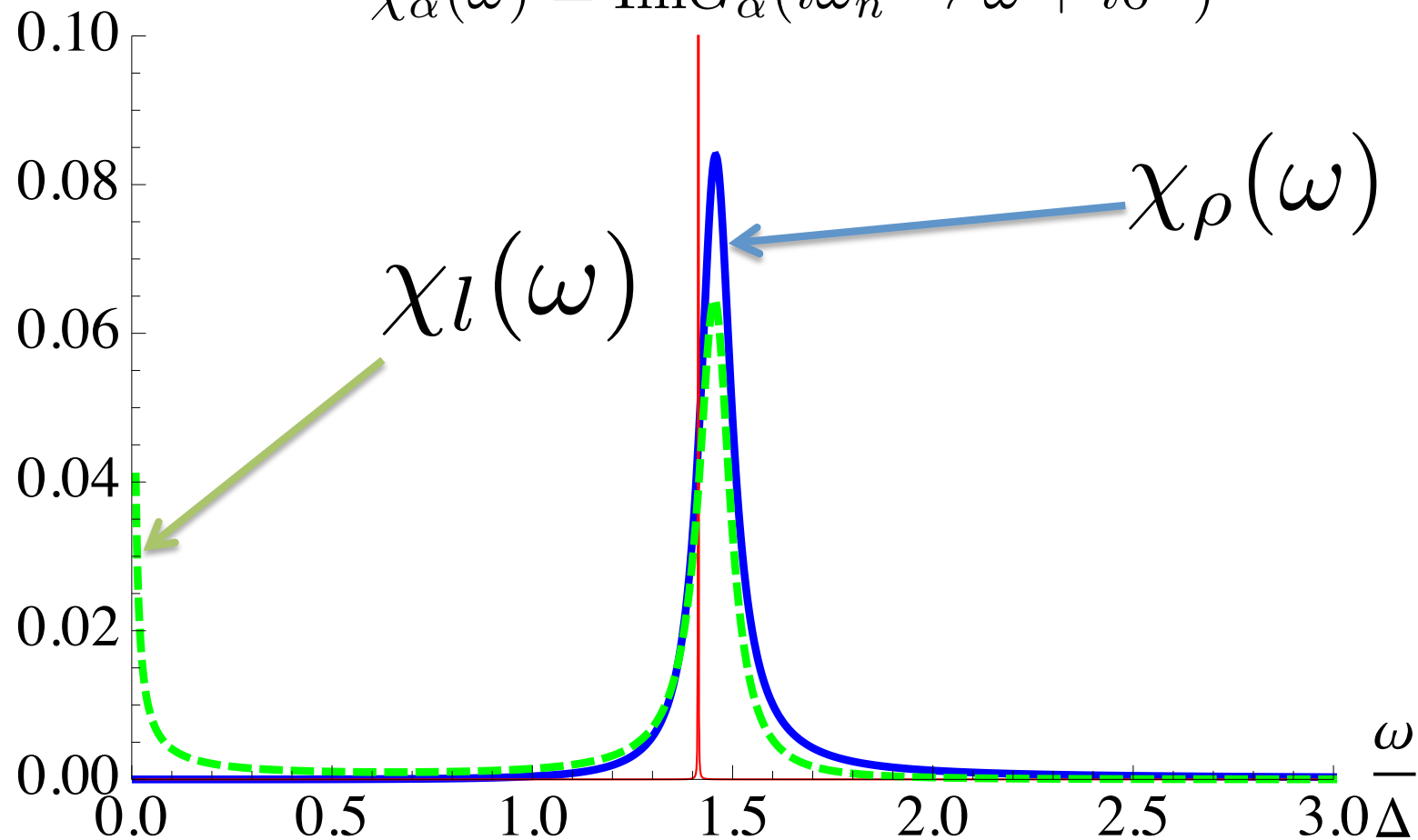
We expect the spectral function to be of order ω^{d+1} at small ω (no divergence), so the resonance might still be visible close to the QCP.

But is the amplitude mode still well defined close to the quantum critical point ?

Amplitude mode : Large N

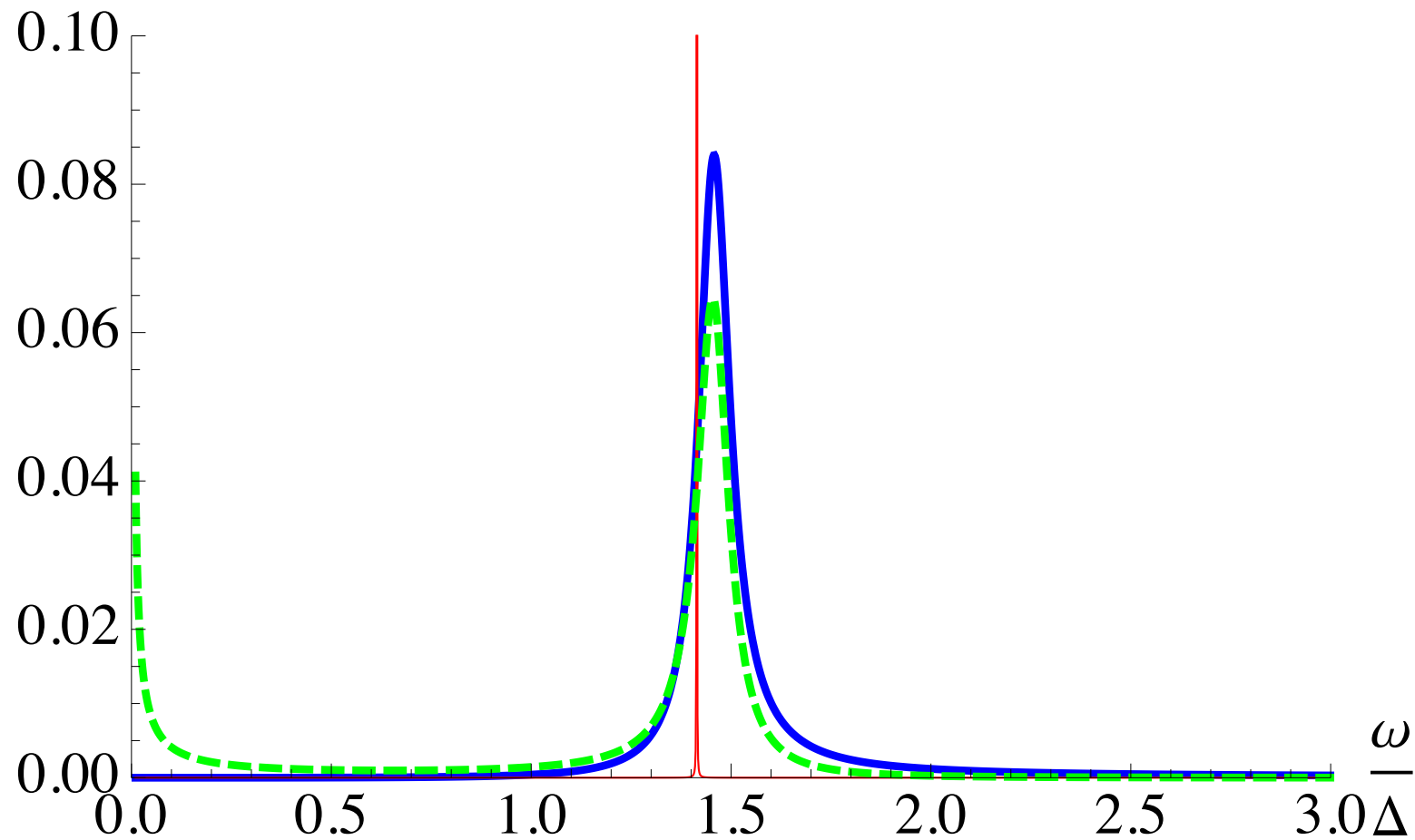
Spectral functions : far from criticality

$$\chi_\alpha(\omega) = \text{Im}G_\alpha(i\omega_n \rightarrow \omega + i0^+)$$



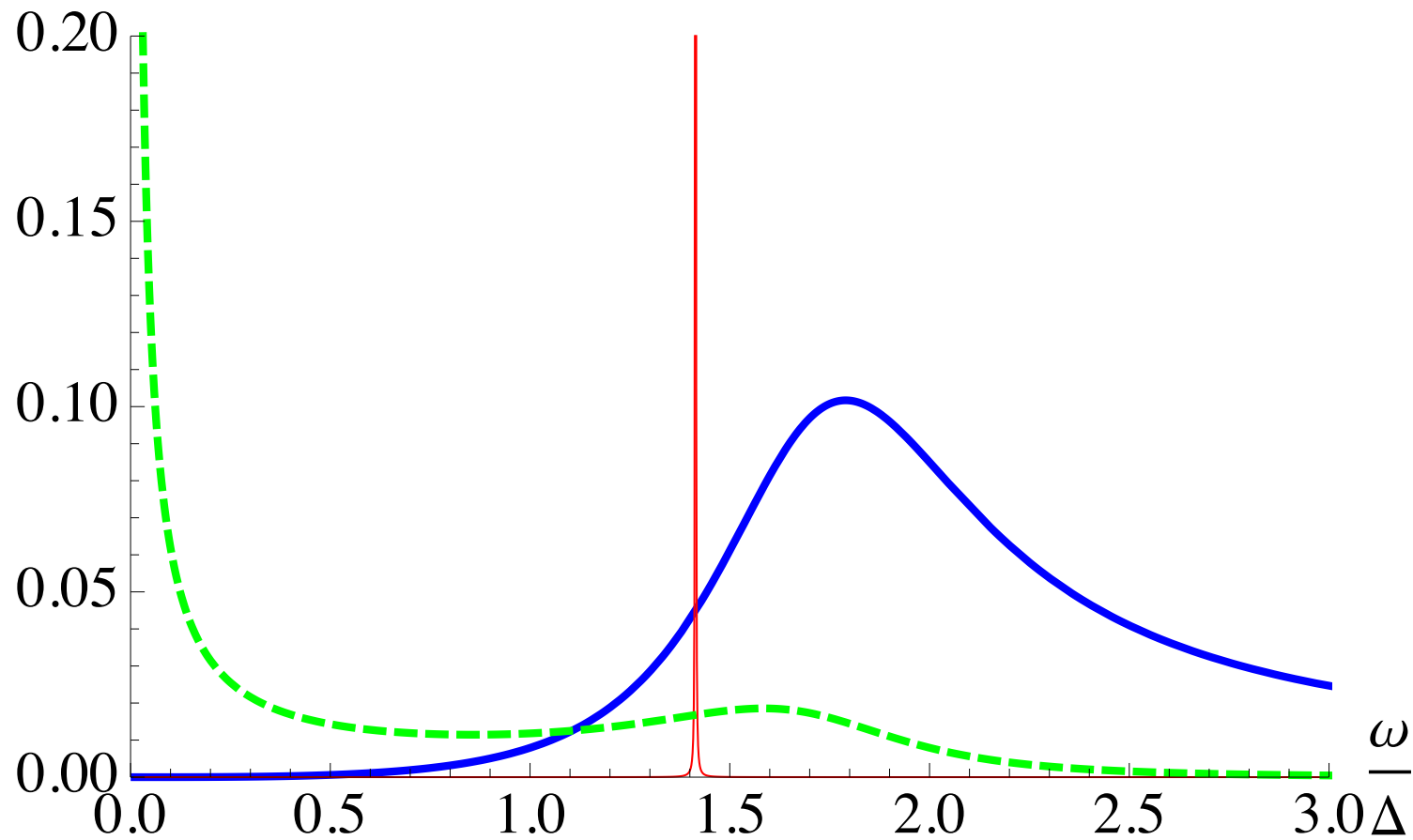
Amplitude mode : Large N

Spectral functions : far from criticality



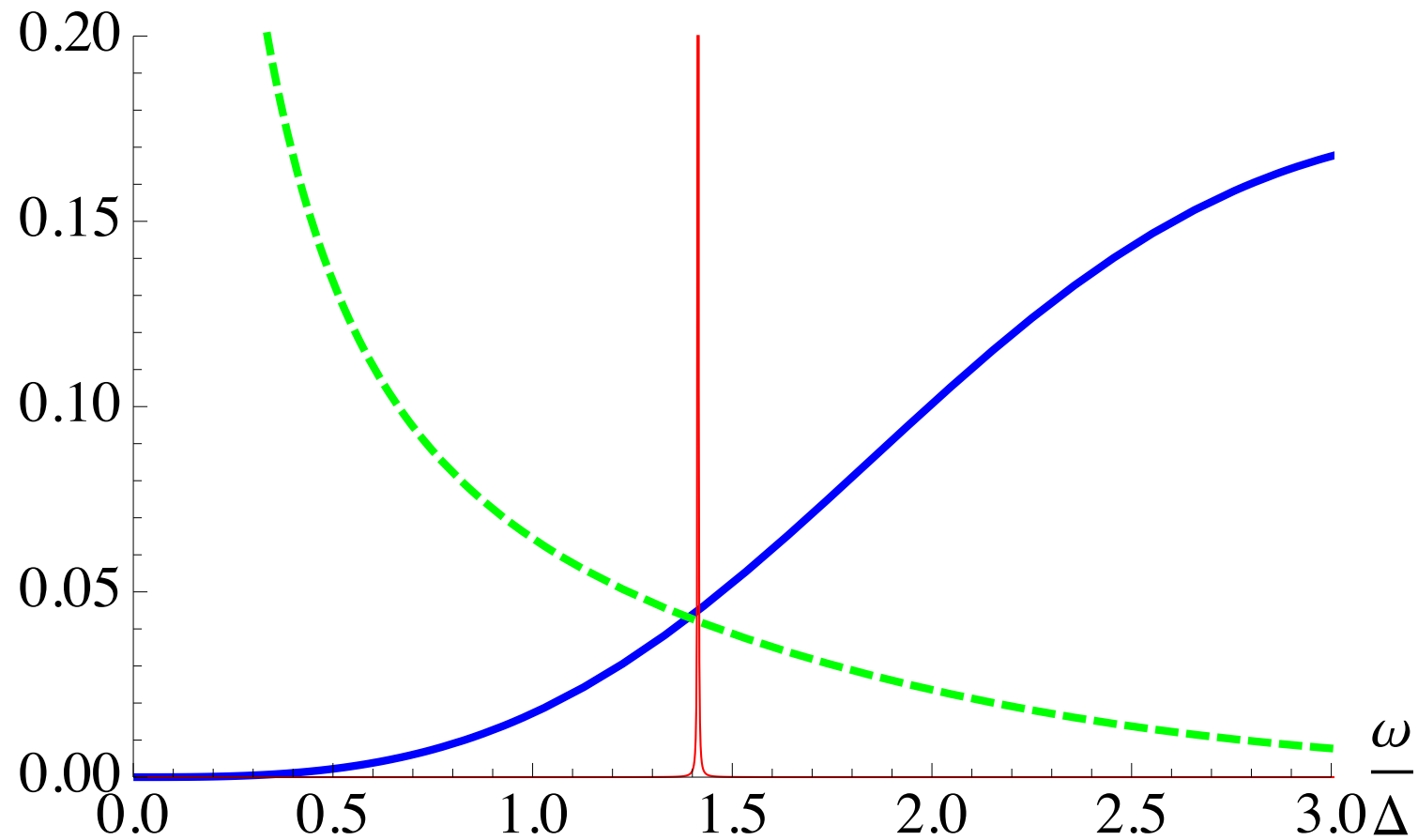
Amplitude mode : Large N

Spectral functions : closer to criticality



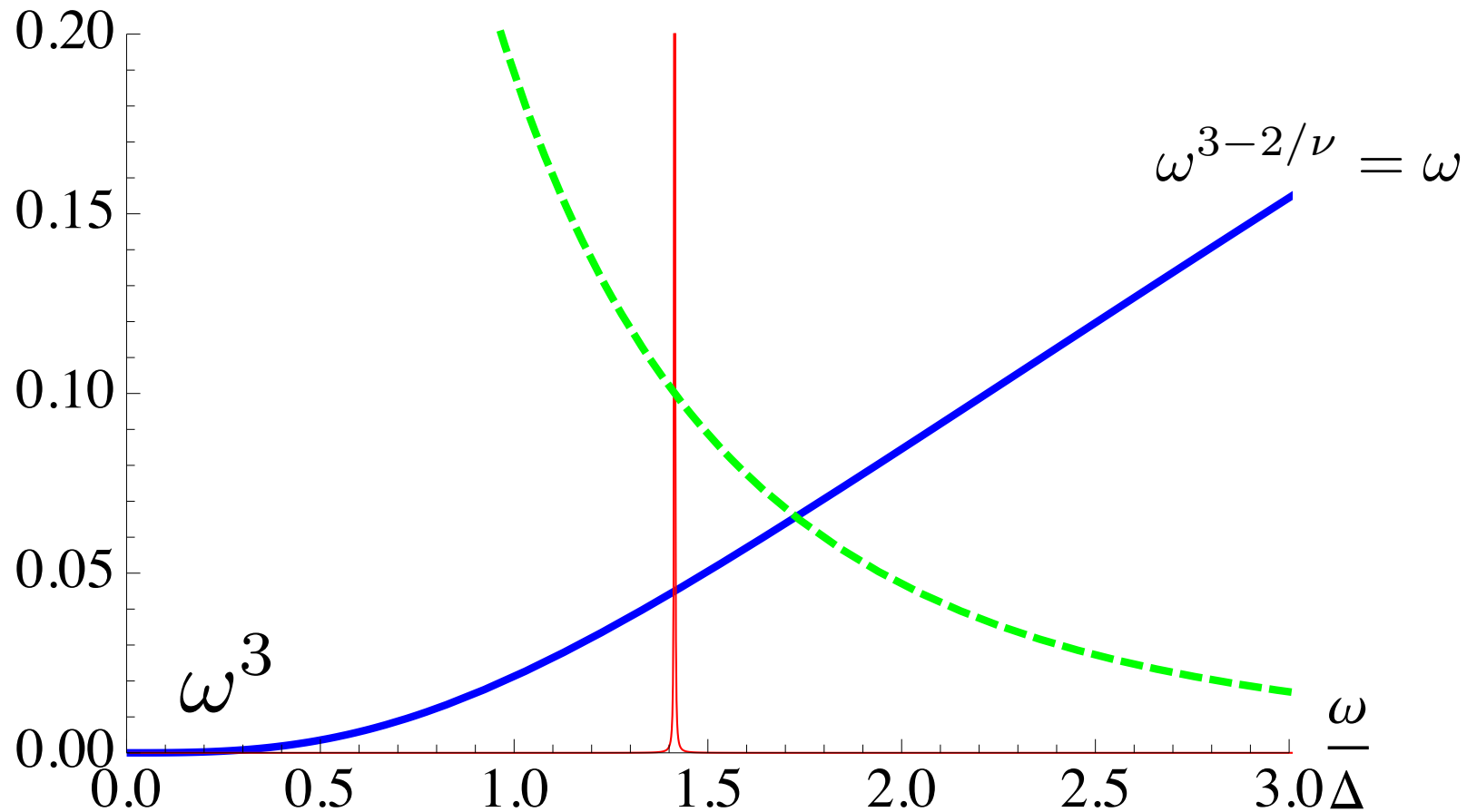
Amplitude mode : Large N

Spectral functions : closer to criticality



Amplitude mode : Large N

Spectral functions : critical regime – no amplitude mode



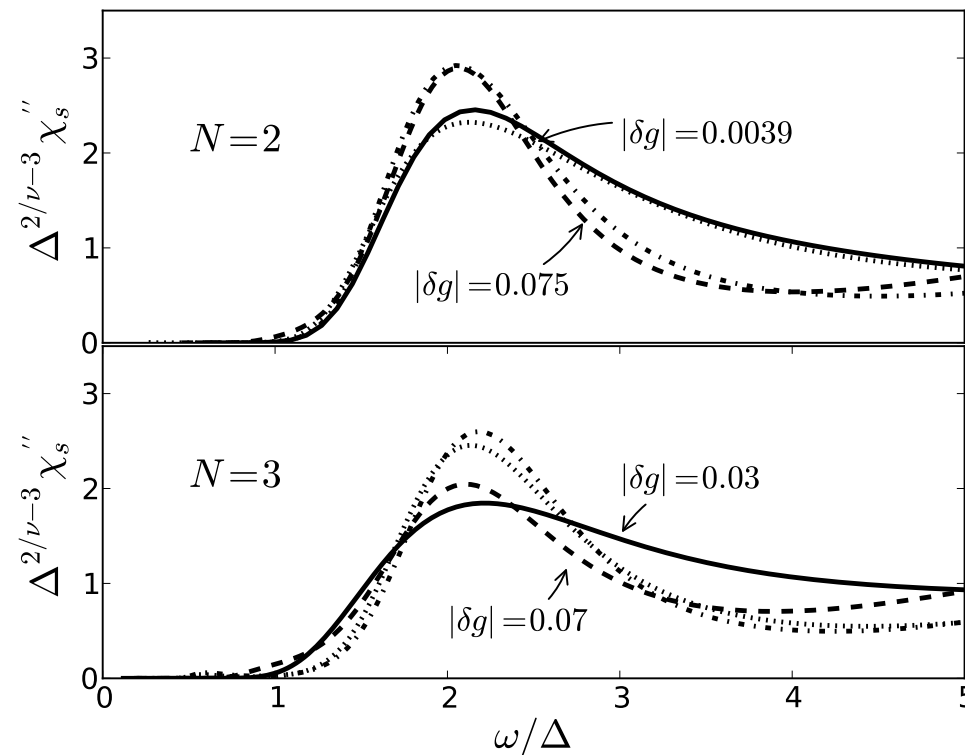
For $1/N$ corrections, see (if you dare) Podolsky & Sachdev, PRB2012

Amplitude mode : Small N Monte Carlo

Recent Monte Carlo simulations : Polet et al. PRL 2012, Gazit et al. PRL 2013, Chen et al. PRL 2013 for $N=2$ and 3

resonance between
 2.1Δ and 3.3Δ

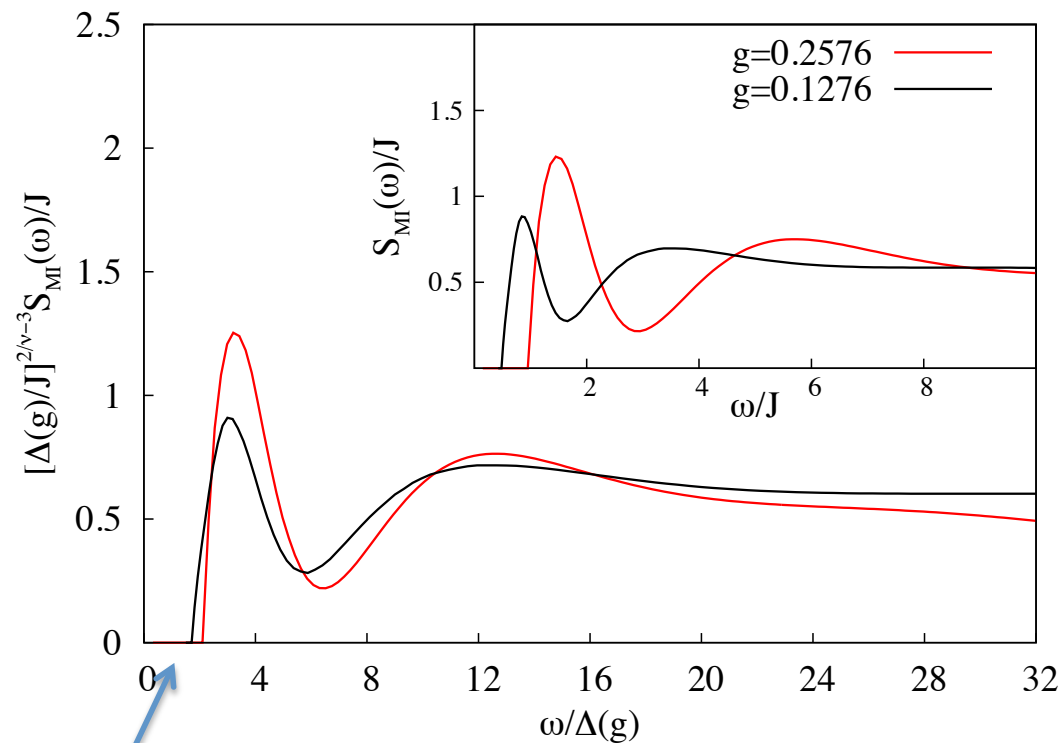
$\approx 2.2\Delta$



Amplitude mode defined close to criticality,
but how does it disappear as N increases ?

Amplitude mode : disordered phase ?

Chen et al. : amplitude mode in the symmetric (disordered) phase !?

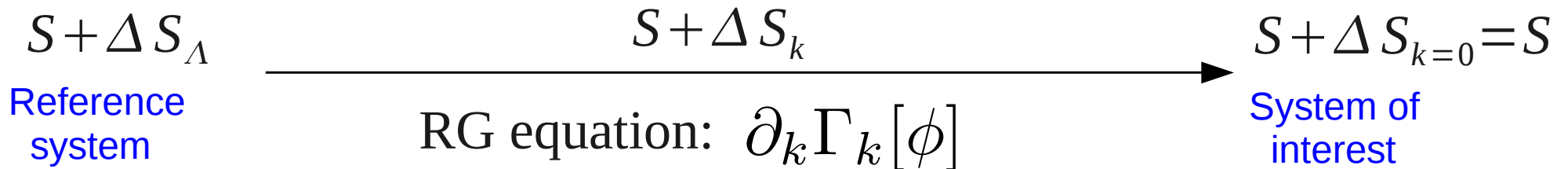


Gapped system : threshold at $\omega=2\Delta$ ($\rho \rightarrow 2$ massive bosons)

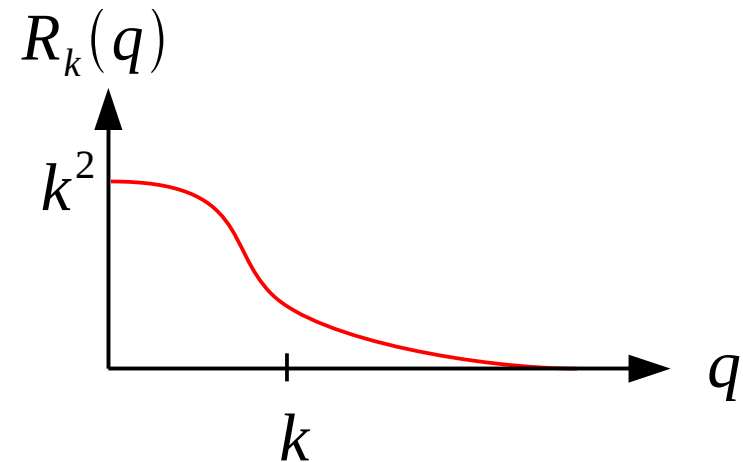
Gave an RG argument, but is it true ?

Non-Perturbative Renormalization Group

Family of actions indexed by momentum scale k



$$\Delta S_k[\psi] = \sum_i \int_q R_k(q) \psi_i(q) \psi_i(-q)$$



Exact Flow equation (Wetterich '93) : $\partial_k \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \left\{ \partial_k R_k \left(\Gamma_k^{(2)} + R_k \right)^{-1} \right\}$

Average Effective Action ("Gibbs" free energy)

Bare action :

$$S = \int_0^\beta d\tau \int d^d x \left\{ \frac{1}{2} \psi_i (-c^{-2} \partial_\tau^2 - \nabla^2 + r_0) \psi_i^2 + \frac{u_0}{4!N} (\psi_i^2)^2 \right\}$$

"Functional free energy" gives access to all the physics :

$$W_k[J, h] = \ln \int D\psi_i e^{-S - \Delta S_k + J_i \psi_i + h\rho}$$

modified Legendre transform

$$\phi_i(\tau) = \langle \psi_i(\tau) \rangle$$

$$\Gamma_k[\phi, h]$$

Source term for the amplitude

$$G_\rho(p) = -\Gamma^{(0,2)}(p) + \Gamma^{(1,1)}(p) G_l(p) \Gamma^{(1,1)}(p)$$

Ansatz

$$\Gamma[\phi, h] = \Gamma[\phi, 0] + \int_{x,y} \frac{1}{2}(\phi^2(x) - \phi_0^2)H_1(x-y)h(y) + \int_{x,y} h(x)H_2(x-y)h(y)$$

Standard derivative and field expansion

Parameterization of longitudinal propagator

$$\Gamma_l^{(2,0)}(p) = p^2 + \Delta_A(p) + \phi_0^2 \Delta_B(p)$$

Computed in a “poor man’s” BMW approximation (Benitez et al. PRB 2008)

Flows

h-dependent part of the flow

$$\partial_k \text{ (wavy line) } \bullet \text{ (wavy line)} = \text{ (wavy line) } \bullet \text{ (loop with } x \text{)} \bullet \text{ (wavy line)}$$

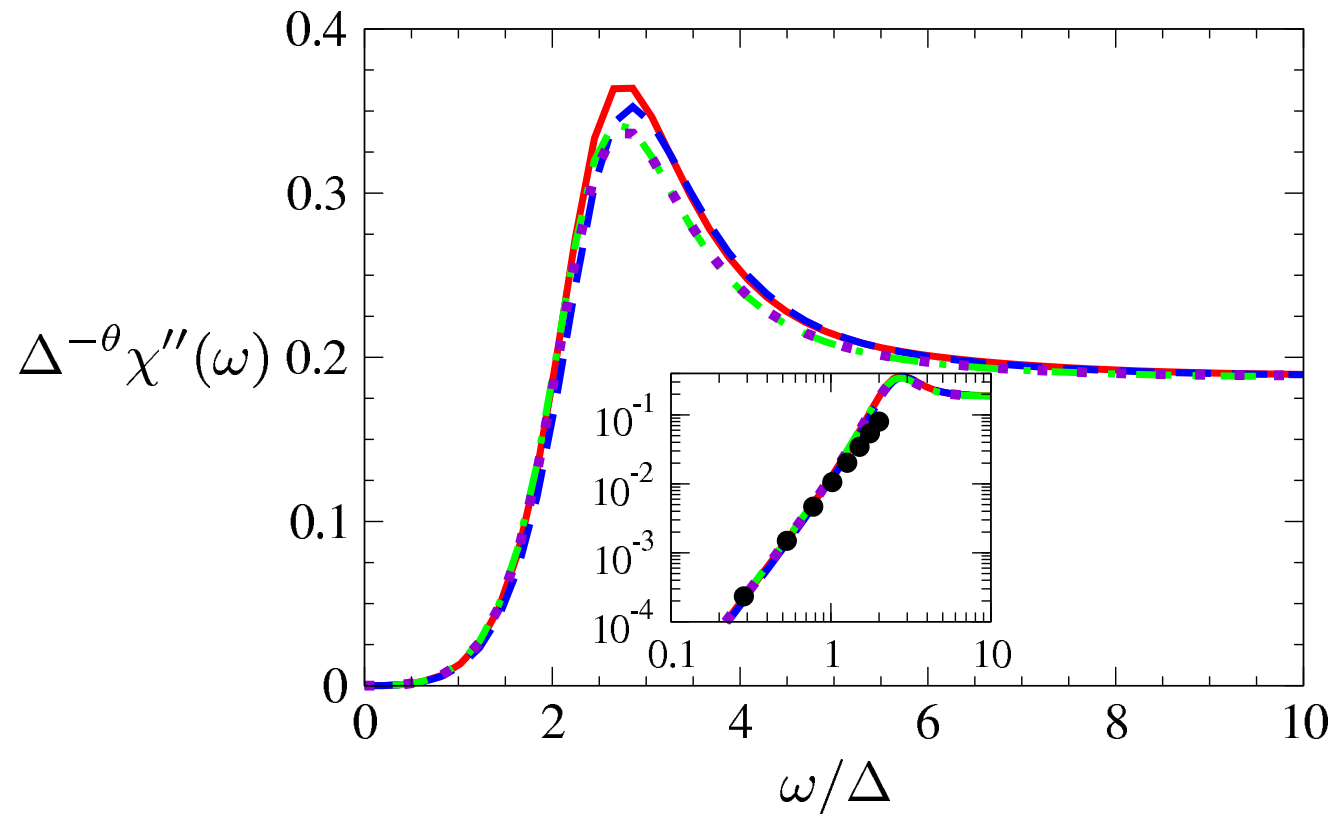
$$\partial_k \text{ (wavy line) } \bullet \text{ (arrow)} = \text{ (wavy line) } \bullet \text{ (loop with } x \text{)} \bullet \text{ (arrow)}$$

Propagators approximated by a derivative expansion but keep the Matsubara frequency dependence of the vertices

At the end of the day, numerical analytic continuation (Pade approximant)

Amplitude mode : N=2

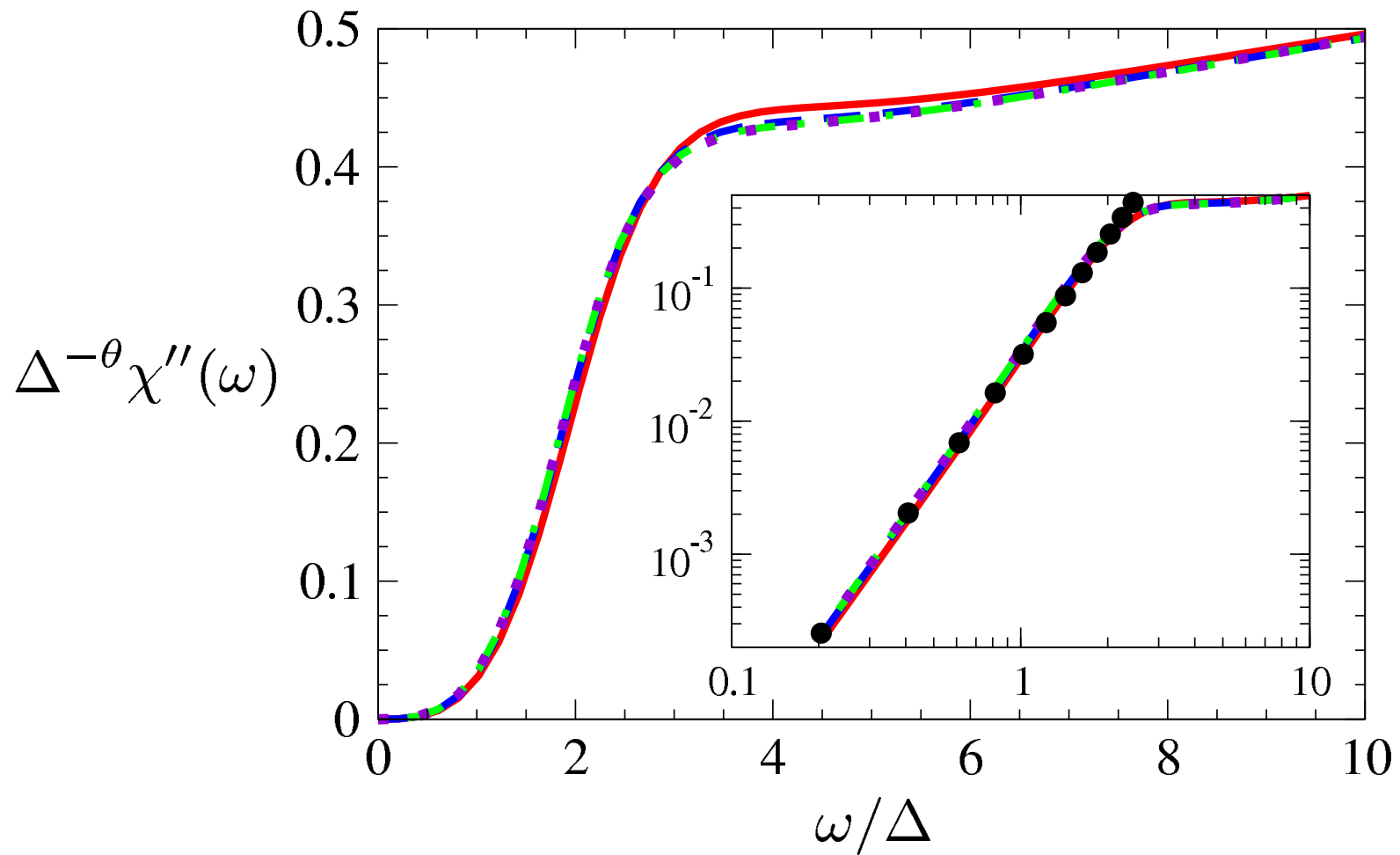
N=2, $m \approx 2.4 \Delta$



Wrong scaling exponent

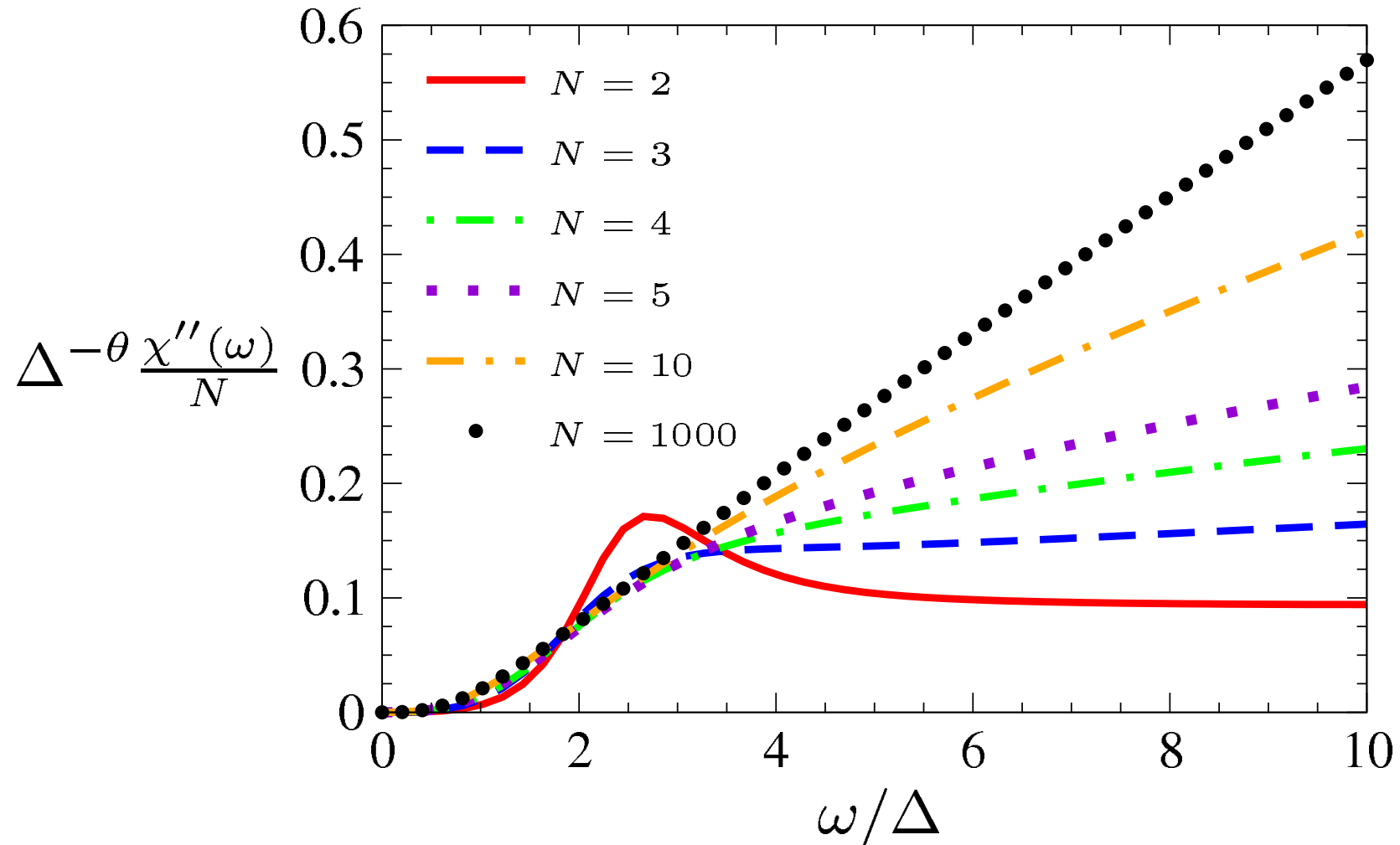
$$\theta \neq 3 - 2/\nu$$

Amplitude mode : N=3



Amplitude mode : all N

No resonance for $N > 2$



T=0 disordered phase - 1

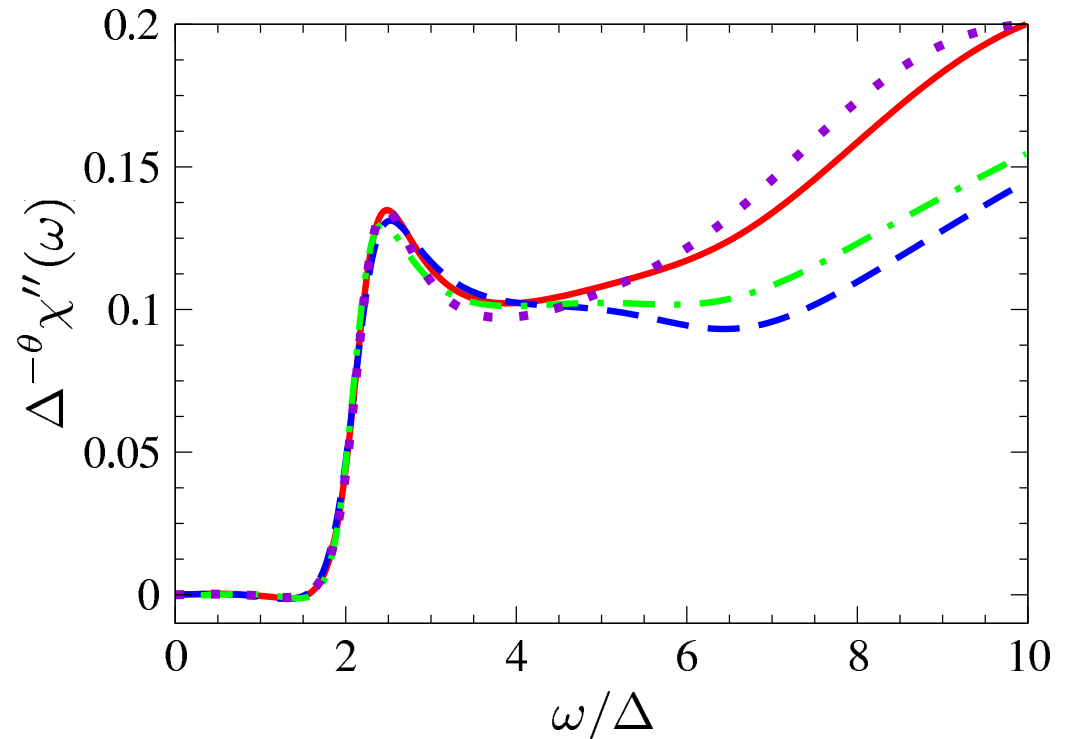
- “Resonance” close to gap at 2Δ :
related to amplitude mode ?
- If so : resonance must appear in the
critical regime (otherwise : not related)

-But in critical regime : scaling

$$\chi''_{\text{crit}}(\omega) \propto \omega^{3-2/\nu}$$

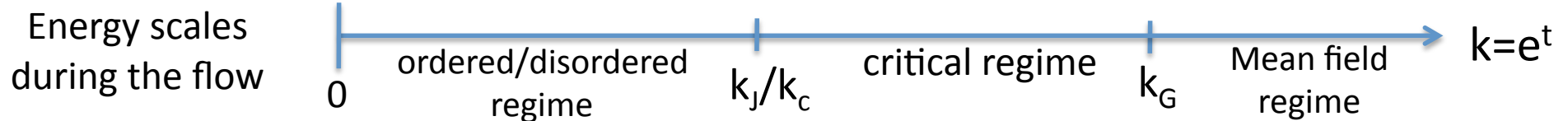
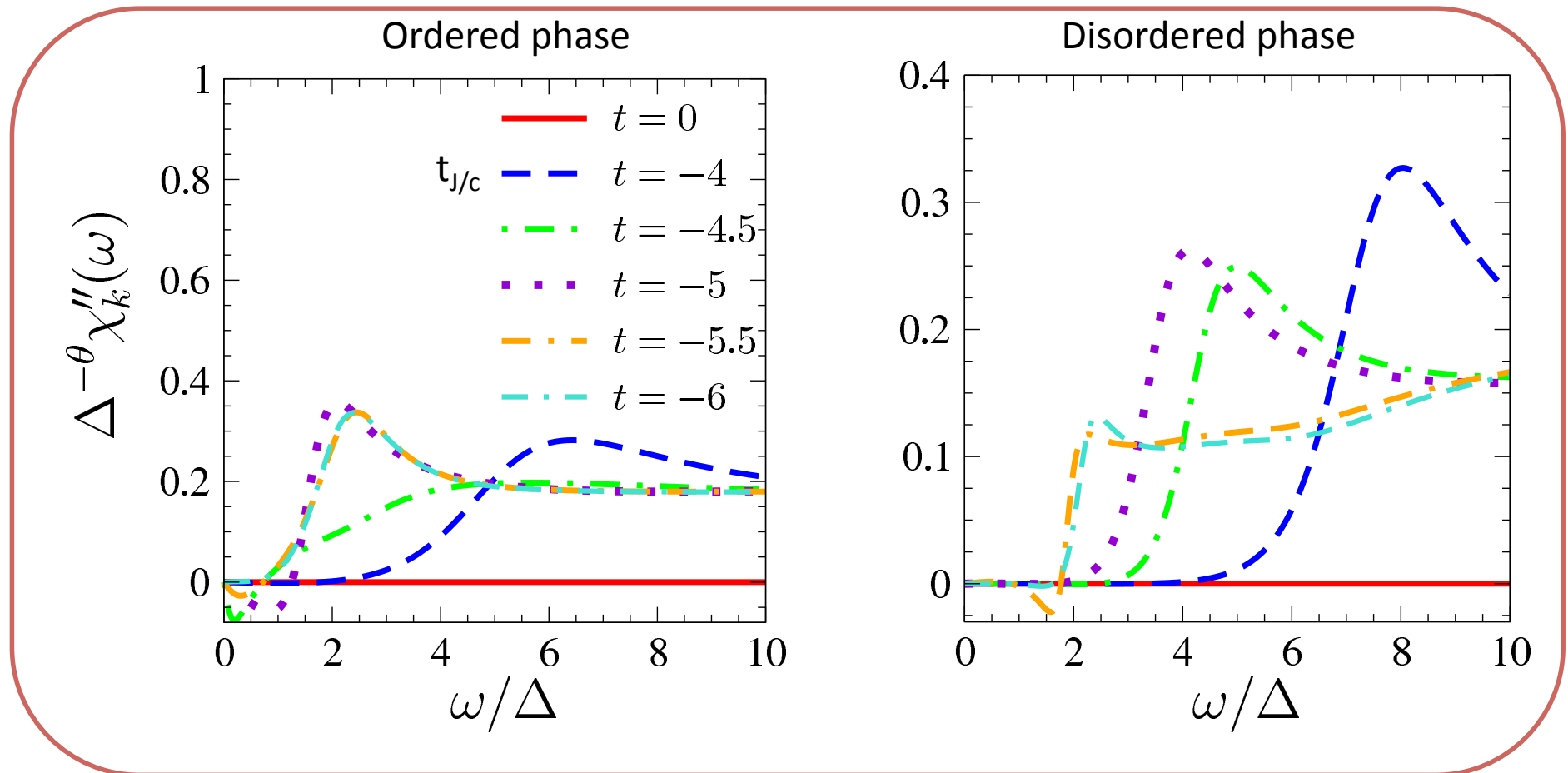
-Implies no resonance in critical regime

- Thus resonance in ordered/ disordered regime are not related – not the same physics



Quantum disordered phase - 2

No resonance in the disordered phase : no amplitude mode precursor in the renormalization flow when the peak appears in the disordered phase.



Amplitude mode : Kosterlitz-Thouless phase

$N=2, T>0$: BKT phase
(quasi-long range order)

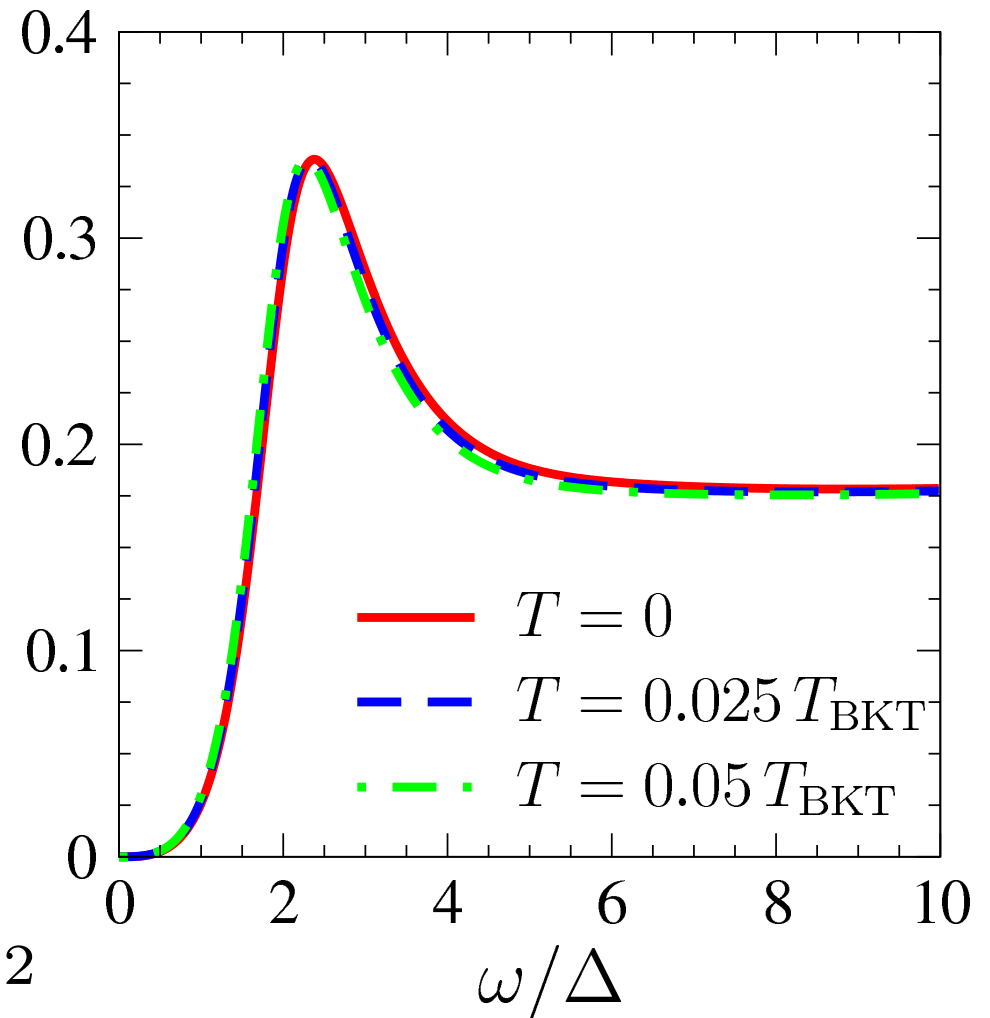
$$T_{BKT} \simeq 0.42\Delta$$

AR et al. PRE 2013

Numerical analytic continuation hard
for $\omega < 2\pi T$

Perturbative calculation :

$$\chi''(\omega) \propto \omega^3 \coth\left(\frac{\omega}{2T}\right) \rightarrow T\omega^2$$



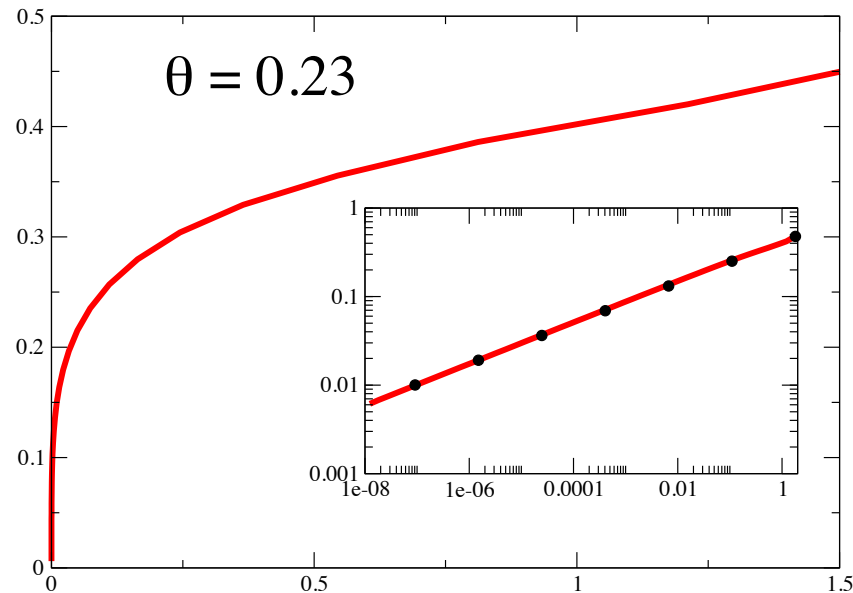
Conclusion and Perspectives

- NPRG study of amplitude mode close to a non-trivial critical point. Available approaches : Monte-Carlo, NPRG, others ?
- Study of the amplitude mode for all N. Nice agreement with Monte-Carlo for N=2. Case N=3 needs further studies :

F. Rose, F. Leonard, B. Delamotte and N. Dupuis :
Self-consistent BMW approach

N = 3

Critical regime :
correct exponents !



Perspectives bis

- No amplitude mode resonance in the disordered regime
- Resonance exists in the BKT phase (vanishing order parameter !). Up to what temperature ? ($T_{\text{BKT}} \approx m_{\text{H}}/5$)

-Numerical analytic continuation hard at large T : Strodtzoff et al. scheme ?

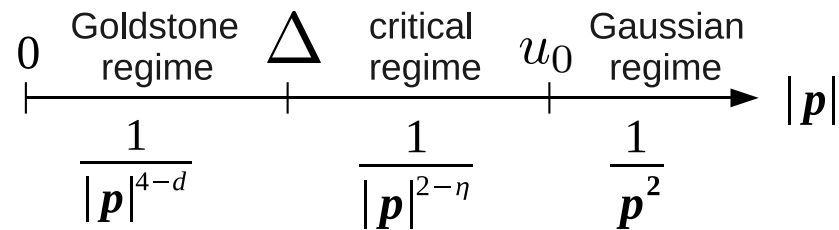
-Large N, resonance exists for $2.7 < d \leq 3$, and “high temperature” (fraction of T_c)

Typical energy scales

$$S = \int_0^\beta d\tau \int d^d x \left\{ \frac{1}{2} \psi_i (-c^{-2} \partial_\tau^2 - \nabla^2 + r_0) \psi_i^2 + \frac{u_0}{4!N} (\psi_i^2)^2 \right\}$$

Four energy scales : T , Δ , u_0 , $\omega_c \propto \sqrt{r_0 - r_{0,c}}$ (not all independent)

(a) Critical regime: $\omega_c \ll u_0$



(b) Non-critical regime: $\omega_c \gg u_0$

