

# Dynamical Locking of the Chiral and the Deconfinement Phase Transition in QCD at Finite Chemical Potential

**Paul Springer**

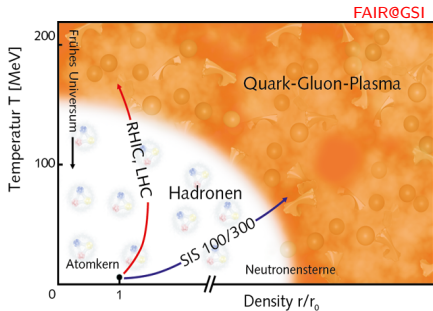
Jens Braun, Marc Leonhardt, Stefan Rechenberger

7th International ERG Conference  
Lefkada, Greece

September 22, 2014

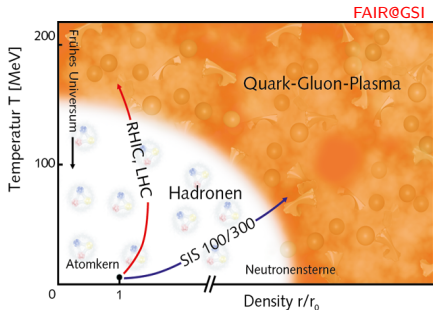


# Motivation



- Confinement  $\Leftrightarrow$  **gauge degrees of freedom**
- $\chi SB \Leftrightarrow$  **quark self-interactions**

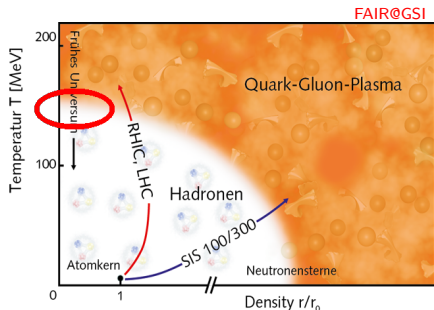
# Motivation



- Confinement  $\Leftrightarrow$  **gauge degrees of freedom**
- $\chi SB \Leftrightarrow$  **quark self-interactions**



However, quark self-interactions are generated by **gluodynamics**



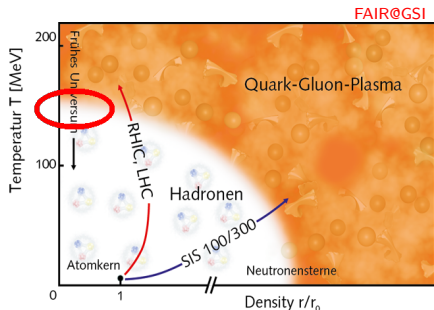
- Confinement  $\Leftrightarrow$  **gauge degrees of freedom**
- $\chi SB \Leftrightarrow$  **quark self-interactions**



However, quark self-interactions are generated by **gluodynamics**

Lattice QCD:

At  $\mu = 0$  pseudo-critical temperatures are very similar for both crossovers  
e.g. [Karsch et al., 2003], [Endrodi et al., 2006], [Aoki et al., 2009] etc.



- Confinement  $\Leftrightarrow$  **gauge degrees of freedom**
- $\chi SB \Leftrightarrow$  **quark self-interactions**



However, quark self-interactions are generated by **gluodynamics**

Lattice QCD:

At  $\mu = 0$  pseudo-critical temperatures are very similar for both crossovers  
e.g. [Karsch et al., 2003], [Endrodi et al., 2006], [Aoki et al., 2009] etc.

## Deeper relation between chiral and confining dynamics???

We investigate  $\lambda_\psi$ -deformed QCD (model) without gluons (basically PNJL model) with two massless flavors,  $N_c$  colors and finite chemical potential:

$$\mathcal{L} = \bar{\psi}(i\not{\partial} + \gamma_0 \bar{g} \langle A_0 \rangle + i\gamma_0 \mu)\psi + \frac{\bar{\lambda}_\psi}{2} [(\bar{\psi}\psi)^2 - (\bar{\psi}\vec{\tau}\gamma_5\psi)^2]$$

- **two** parameters:  $\lambda_\psi(\Lambda)$ ,  $\langle A_0 \rangle$

- large  $\lambda_\psi$  triggers  $\chi$ SB

- deconfinement order parameter:

$$\text{Tr}_F L[\langle A_0 \rangle] = \frac{1}{N_c} \text{Tr}_F [\mathcal{P}e^{i\beta \bar{g} \langle A_0 \rangle}] \stackrel{\text{PNJL}}{=} \frac{1}{N_c} \langle \text{Tr}_F [\mathcal{P}e^{i\bar{g} \int_0^\beta A_0}] \rangle$$

e.g. [Meisinger, Ogilvie, 1996]

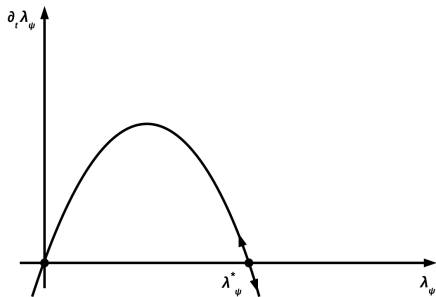
- Tool: **Wetterich flow equation** [C. Wetterich, 1993]

# $\lambda_\psi$ -deformed QCD: RG fixed-point analysis

$$\mathcal{L} = \bar{\psi}(i\cancel{D} + \bar{\lambda}_\psi \psi + \frac{\bar{\lambda}_\psi}{2} [(\bar{\psi}\psi)^2 - (\bar{\psi}\vec{\tau}\gamma_5\psi)^2])$$

RG-flow equation:  
 $T = 0, \mu = 0, \langle A_0 \rangle = 0$   
( $k$  is momentum scale)

$$k\partial_k \lambda_\psi = 2\lambda_\psi - C\lambda_\psi^2$$



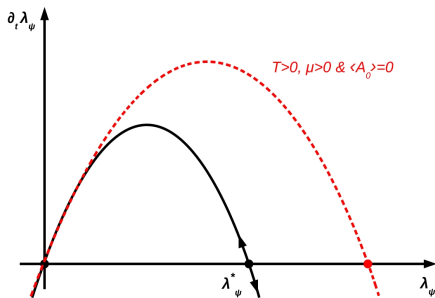
$$\lambda_\psi(\Lambda) > \lambda_\psi^* \Rightarrow \chi SB$$

# $\lambda_\psi$ -deformed QCD: RG fixed-point analysis

$$\mathcal{L} = \bar{\psi}(i\not{\partial} + \quad + i\gamma_0\mu)\psi + \frac{\bar{\lambda}_\psi}{2} [(\bar{\psi}\psi)^2 - (\bar{\psi}\vec{\tau}\gamma_5\psi)^2]$$

RG-flow equation:  
 $T \neq 0, \mu \neq 0, \langle A_0 \rangle = 0$   
 (k is momentum scale)

$$k\partial_k \lambda_\psi = 2\lambda_\psi - C\left(\frac{T}{k}, \frac{\mu}{k}\right)\lambda_\psi^2$$



$\lambda_\psi(\Lambda) > \lambda_\psi^*, T$  or (and)  $\mu$  increase  $\Rightarrow$  restoration of  $\chi$ -Symmetry

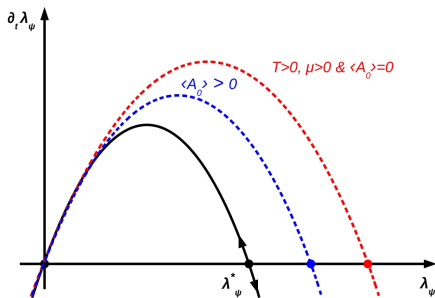


# $\lambda_\psi$ -deformed QCD: RG fixed-point analysis

$$\mathcal{L} = \bar{\psi}(i\not{\partial} + \gamma_0 \bar{g} \langle A_0 \rangle + i\gamma_0 \mu)\psi + \frac{\bar{\lambda}_\psi}{2} [(\bar{\psi}\psi)^2 - (\bar{\psi}\vec{\tau}\gamma_5\psi)^2]$$

RG-flow equation:  
 $T \neq 0, \mu \neq 0, \langle A_0 \rangle \neq 0$   
 (color-confined regime)

$$k\partial_k \lambda_\psi = 2\lambda_\psi - C\left(\frac{T}{k}, \frac{\mu}{k}, \langle A_0 \rangle\right) \lambda_\psi^2$$



Finite  $\langle A_0 \rangle \Rightarrow$  fixed point “moves” to the left

# $\lambda_\psi$ -deformed QCD: RG fixed-point analysis

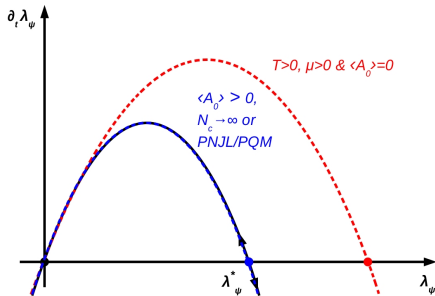
$$\mathcal{L} = \bar{\psi}(i\not{\partial} + \gamma_0 \bar{g} \langle A_0 \rangle + i\gamma_0 \mu)\psi + \frac{\bar{\lambda}_\psi}{2} [(\bar{\psi}\psi)^2 - (\bar{\psi}\vec{\tau}\gamma_5\psi)^2]$$

## Analytical result:

as long

- $N_c \rightarrow \infty$ , **or**
- $\text{Tr}_F L[\langle A_0 \rangle] \approx \langle \text{Tr}_F L[A_0] \rangle$   
(PNJL/PQM-models)

$$\Rightarrow \lambda^*(T, \mu, \langle A_0 \rangle) = \lambda^*(0, 0, 0)$$



$$T_\chi \geq T_d$$

- **PNJL/PQM-models  $\iff$  Large- $N_c$  in the coupling of the matter and gauge sector** (should not be confused with the standard large- $N_c$  approximation, such as neglecting pion fluctuations etc.)
- $T_\chi \geq T_d$  in the phase diagram of PNJL/PQM-models
  - $\Rightarrow$  existence of quarkyonic phase in PNJL/PQM-models under debate
  - $\Rightarrow$  **Constraint on parametrization of Polyakov potential**

Thank you for your attention!

# Numerical Results

- We use the data for  $\langle A_0 \rangle$  for pure  $SU(N_c)$  gauge theory, i. e., we drop the back coupling of fermions to the gauge sector:  $T_d$  is fixed!

[Braun, Gies, Pawłowski, 2010], [Braun, Eichhorn, Gies, Pawłowski, 2010]

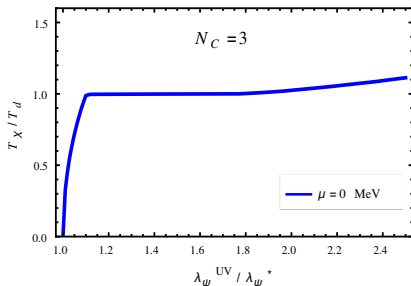
- Back coupling  $\rightarrow$  corrections, but the main results should be the same on the qualitative level

# Numerical Results

- We use the data for  $\langle A_0 \rangle$  for pure  $SU(N_c)$  gauge theory, i. e., we drop the back coupling of fermions to the gauge sector:  $T_d$  is fixed!

[Braun, Gies, Pawłowski, 2010], [Braun, Eichhorn, Gies, Pawłowski, 2010]

- Back coupling  $\rightarrow$  corrections, but the main results should be the same on the qualitative level

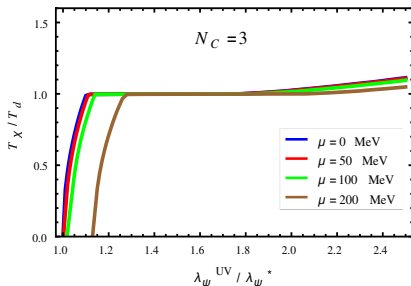


# Numerical Results

- We use the data for  $\langle A_0 \rangle$  for pure  $SU(N_c)$  gauge theory, i. e., we drop the back coupling of fermions to the gauge sector:  $T_d$  is fixed!

[Braun, Gies, Pawłowski, 2010], [Braun, Eichhorn, Gies, Pawłowski, 2010]

- Back coupling  $\rightarrow$  corrections, but the main results should be the same on the qualitative level

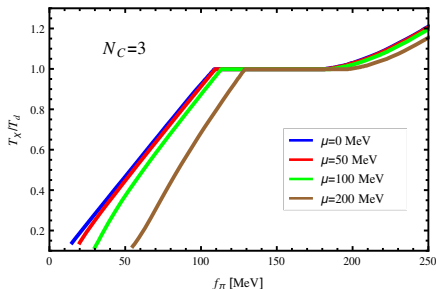
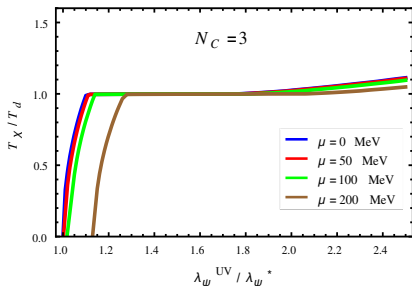


# Numerical Results

- We use the data for  $\langle A_0 \rangle$  for pure  $SU(N_c)$  gauge theory, i. e., we drop the back coupling of fermions to the gauge sector:  $T_d$  is fixed!

[Braun, Gies, Pawłowski, 2010], [Braun, Eichhorn, Gies, Pawłowski, 2010]

- Back coupling  $\rightarrow$  corrections, but the main results should be the same on the qualitative level



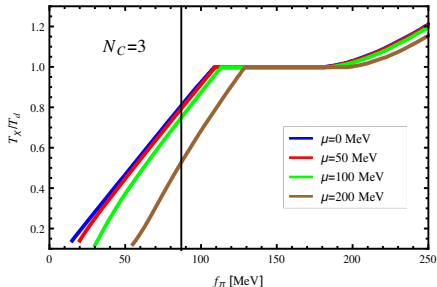
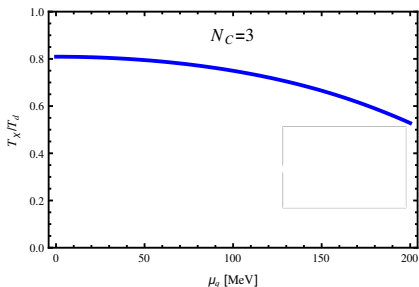


# Numerical Results

- We use the data for  $\langle A_0 \rangle$  for pure  $SU(N_c)$  gauge theory, i. e., we drop the back coupling of fermions to the gauge sector:  $T_d$  is fixed!

[Braun, Gies, Pawłowski, 2010], [Braun, Eichhorn, Gies, Pawłowski, 2010]

- Back coupling  $\rightarrow$  corrections, but the main results should be the same on the qualitative level



$$\kappa = -T_\chi \frac{dT_\chi(\mu^2)}{d(\mu^2)} \Big|_{\mu=0} = 0.385(5)$$