

# FRG approaches to graphene - an overview

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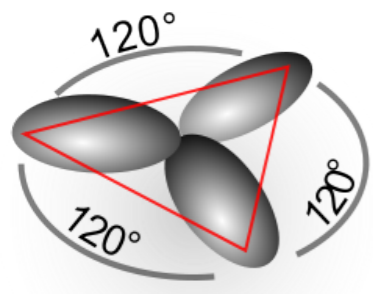
September 24, 2014, ERG 2014, Lefkada Island, Greece

# Outline

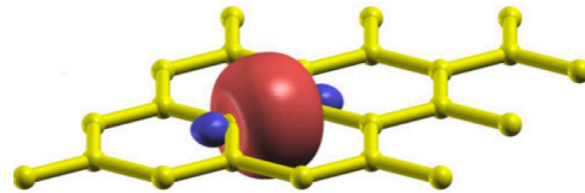
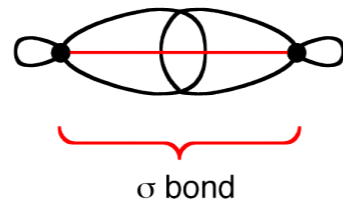
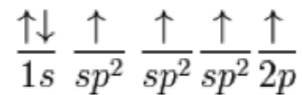
- **Part I: Graphene basics**
- **Part II: Instabilities and phase transitions from (F)RG**
  - ▶ undoped graphene, Dirac electrons
  - ▶ doped graphene @van-Hove singularity
- **Conclusions & Outlook**

# Real space structure

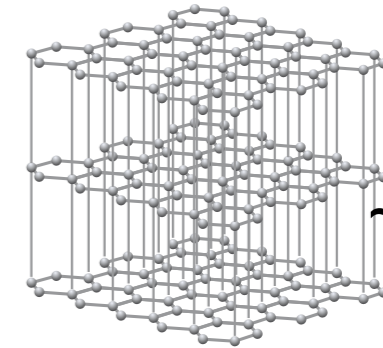
- Graphene: flat single layer of C-atoms in **2D hexagonal lattice**
- $sp^2$ -Hybrid-orbitals of C-atoms  $\rightarrow$   $\sigma$ -bonds away from  $\epsilon_F$



$sp^2$ , planar

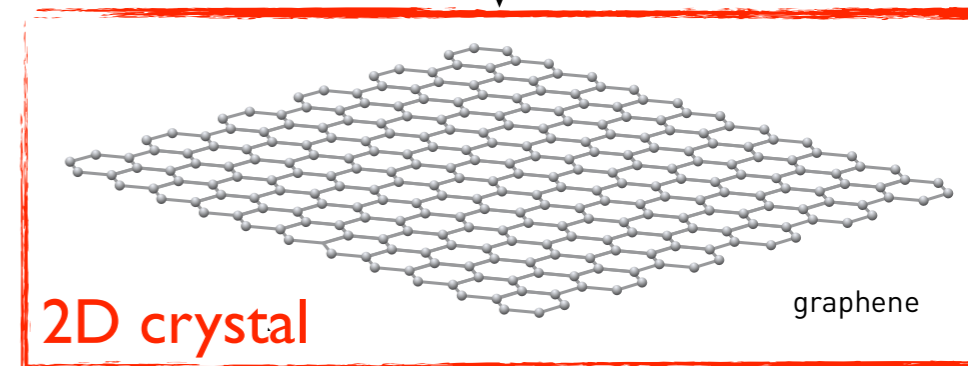
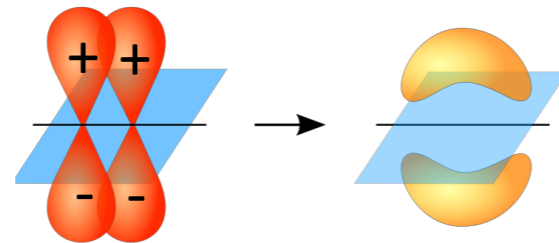
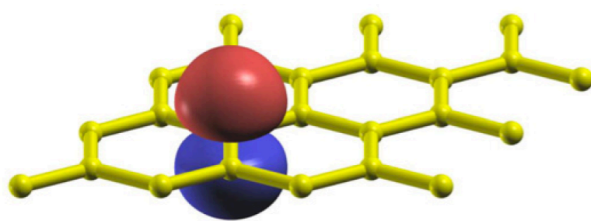


graphite



$\sim 0.335\text{nm}$

- Conduction & valence  $e^-$ :  $\pi$ -bonds of  $p_z$ -orbitals

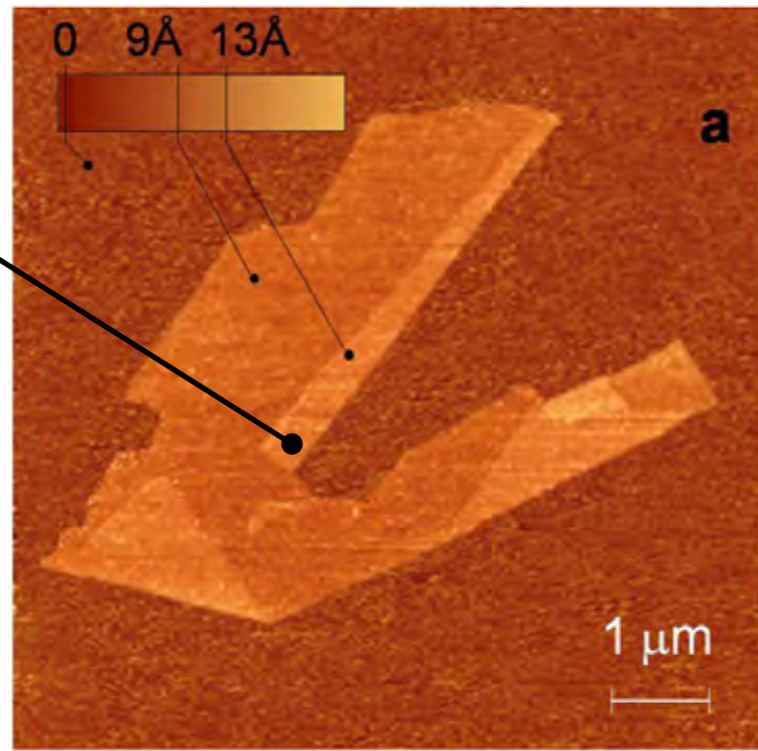


- Van der Waal's bonds between layers

# In the lab

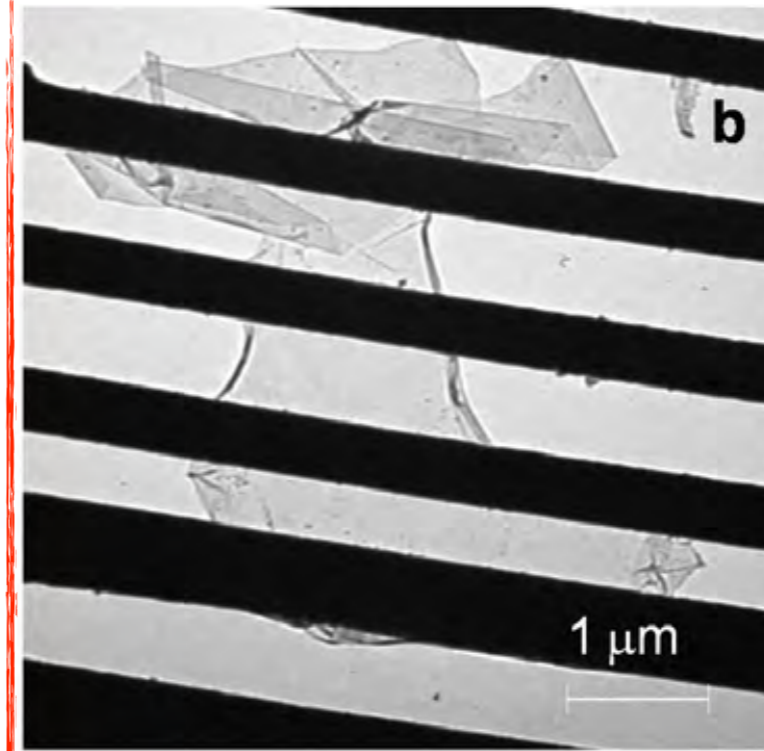
## Atomic-force microscopy

rel. height  $\sim 0.4\text{nm}$   
→ single layer

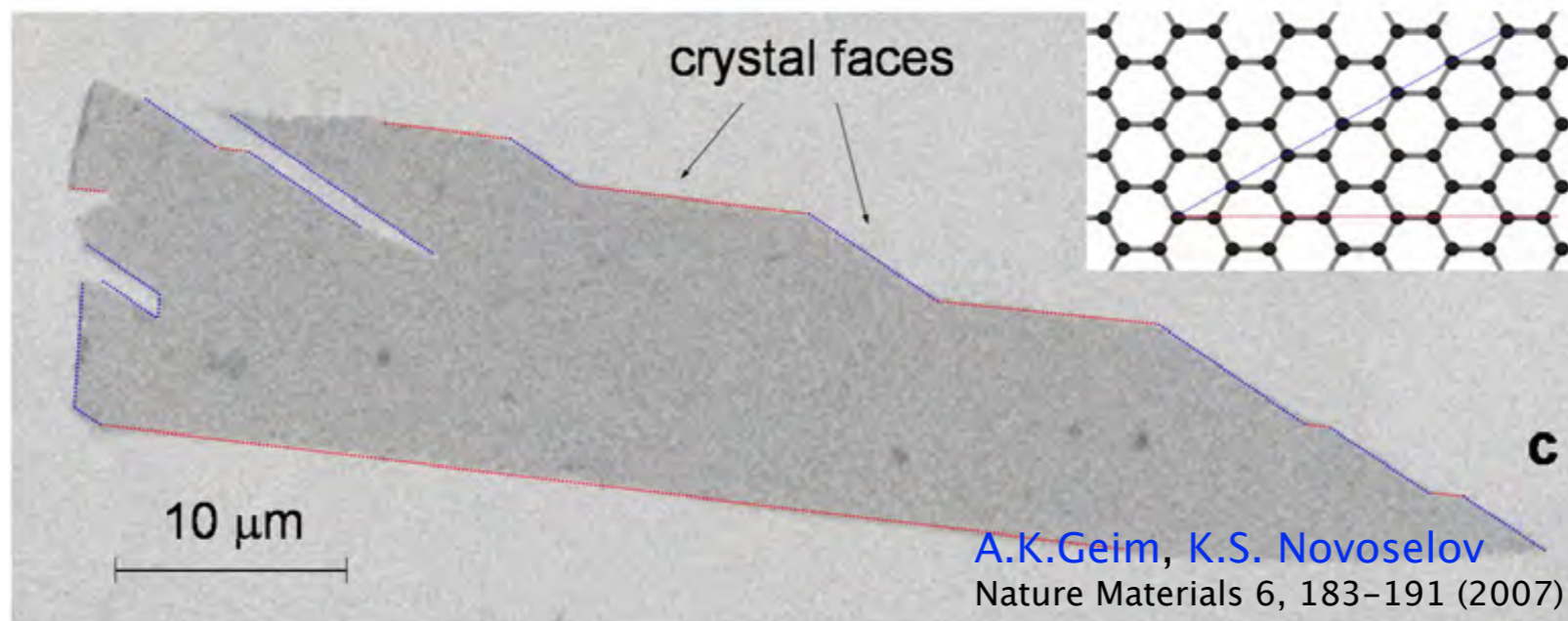


## Transmission-electron spectroscopy

suspended on scaffold



large flake



## Scanning-electron micrograph

# Fun facts

- **Density:**  $0.77\text{mg/m}^2$ 
  - ▶ even the smallest gas atom (He) cannot pass through it
- **Optical transparency:**
  - ▶ Absorbs only 2.3% of the visible light intensity
- **Strength:** breaking strength= $42\text{N/m}$ .
  - ▶ Thin film of steel ( $0.335\text{nm}$ ) has 2D breaking strength  $\sim 0.3\text{N/m}$
- **Electrical conductivity:** as well as copper
  - ▶ 2D sheet conductivity  $\sigma = e n \mu$ 
    - **high electron mobility**  $\mu$ , weakly depends on T even at  $\sim 300\text{K}$
    - mobility remains high even in doped devices  $\rightarrow$  in contrast to bulk semiconductors



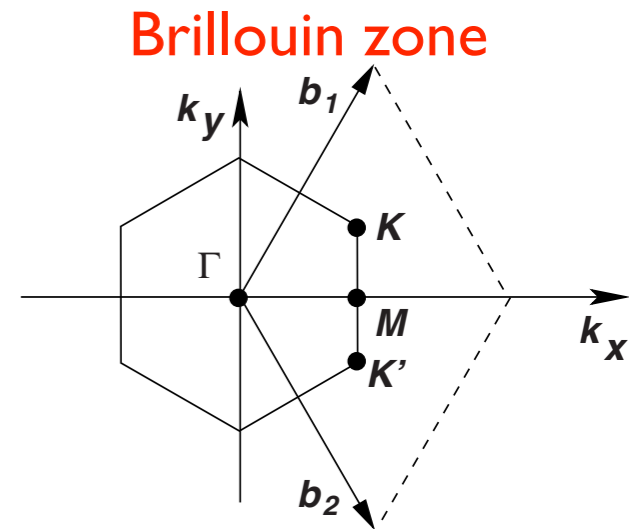
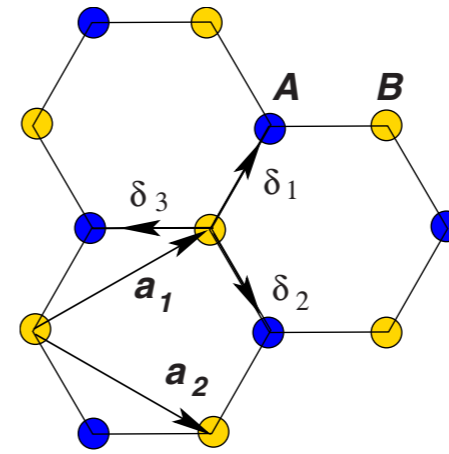
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# Tight-binding model

$$H_0 = -t \sum_{\langle i,j \rangle, \sigma} \left( a_{i,\sigma}^\dagger b_{j,\sigma} + \text{H.c.} \right) - t' \sum_{\langle\langle i,j \rangle\rangle, \sigma} \left( a_{i,\sigma}^\dagger a_{j,\sigma} + b_{i,\sigma}^\dagger b_{j,\sigma} + \text{H.c.} \right)$$

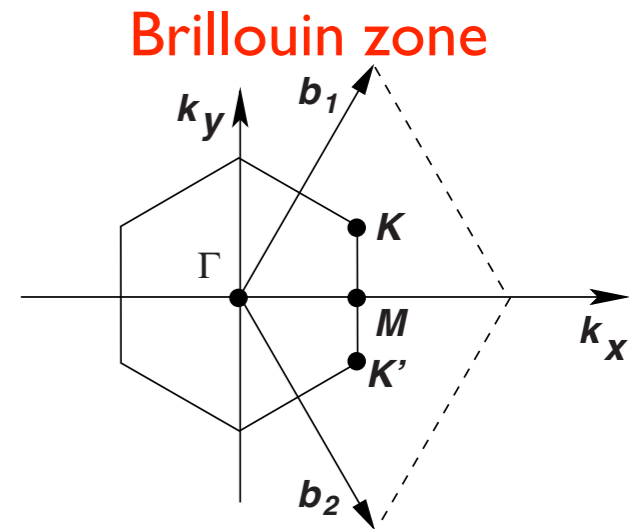
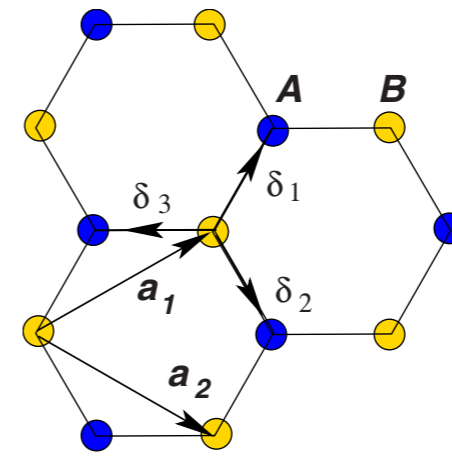


►  $\pi$  band structure of graphene:

$$E_{\pm}(\mathbf{k}) = \pm t \sqrt{3 + f(\mathbf{k})} - t' f(\mathbf{k}), \quad f(\mathbf{k}) = \sum_{i=1}^3 e^{i\delta_i \cdot \mathbf{k}}$$

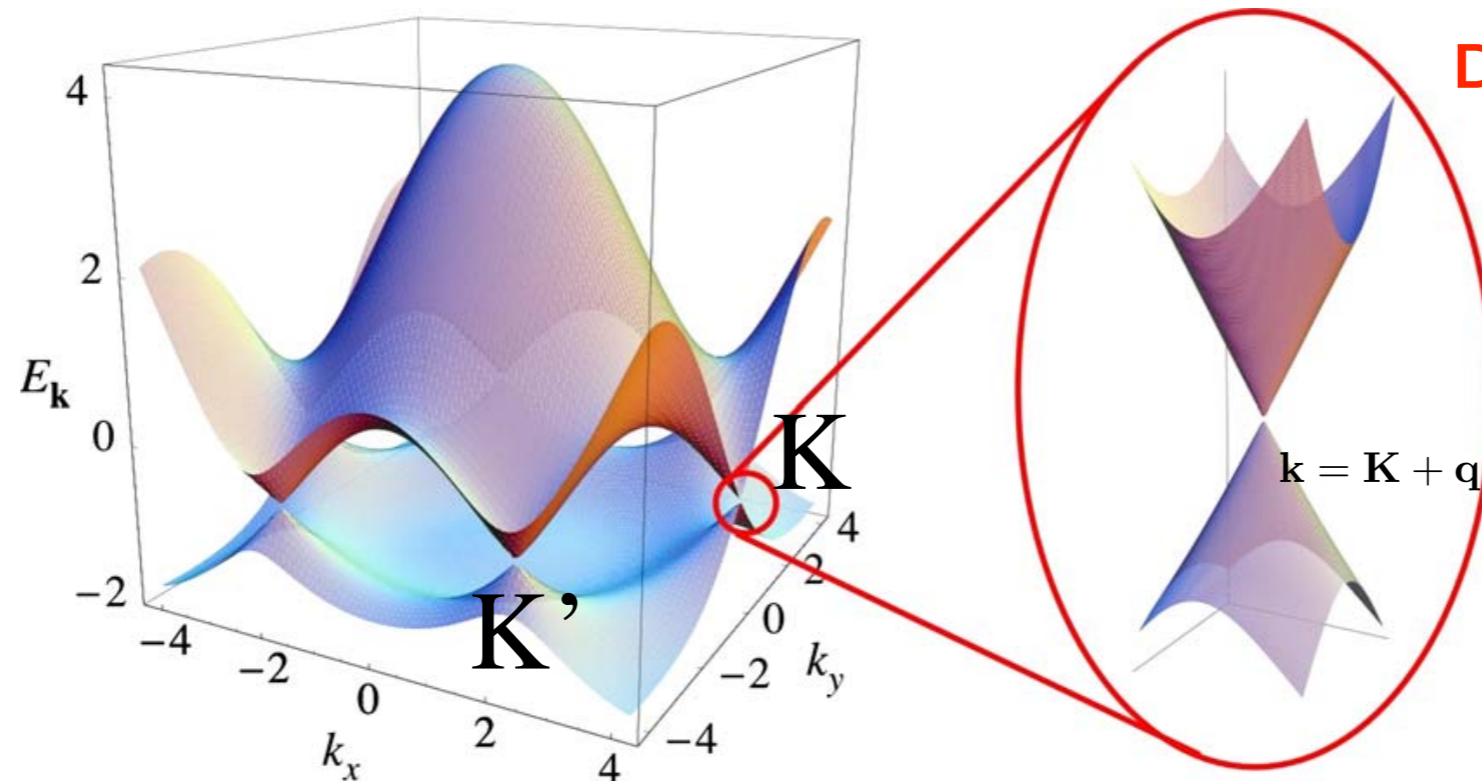
# Tight-binding model

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►  $\pi$  band structure of graphene:  $E_{\pm}(\mathbf{k}) = \pm t \sqrt{3 + f(\mathbf{k})} - t' f(\mathbf{k})$ ,  $f(\mathbf{k}) = \sum_{i=1}^3 e^{i\delta_i \cdot \mathbf{k}}$

► e.g. from *ab initio* calculations:  $t = 2.8 \text{ eV}, t' = -0.2t$



$$E_{\pm}(\mathbf{q}) \approx 3t' \pm v_F |\mathbf{q}|$$

► adjust Fermi level by chemical potential  $\rightarrow$  **doping**



# Dirac fermions in undoped graphene

- Expand operators in Hamiltonian around Fermi level  $\rightarrow$  K, K' points

- ▶ **2D massless Dirac equation (around K point):**

$$-i v_F \vec{\sigma} \cdot \nabla \psi(\vec{r}) = E \psi(\vec{r})$$

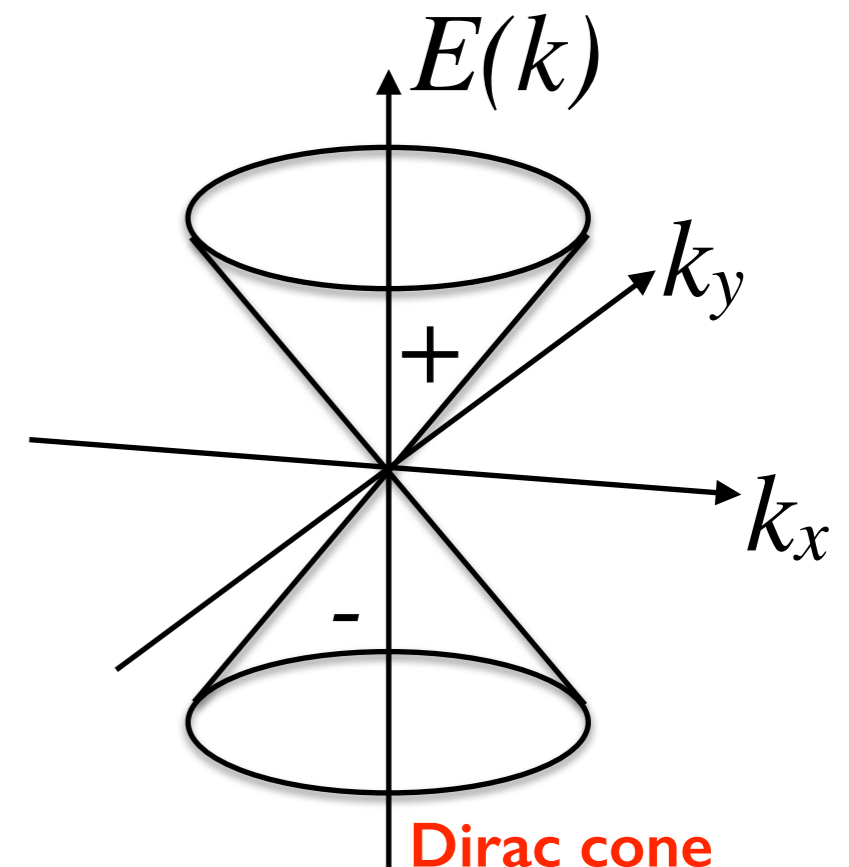
- ▶ near Dirac point electrons have well-defined **chirality**

include K' point

➔ **two copies of massless Dirac-like Hamiltonian**

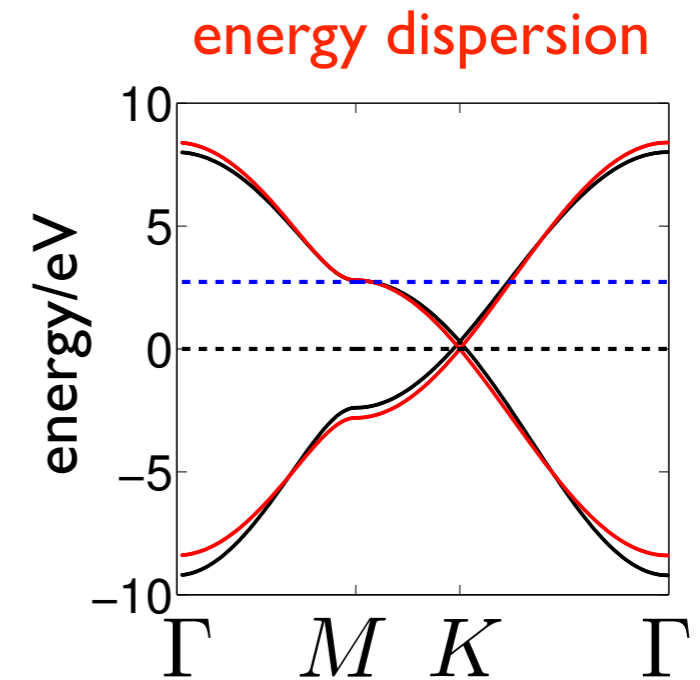
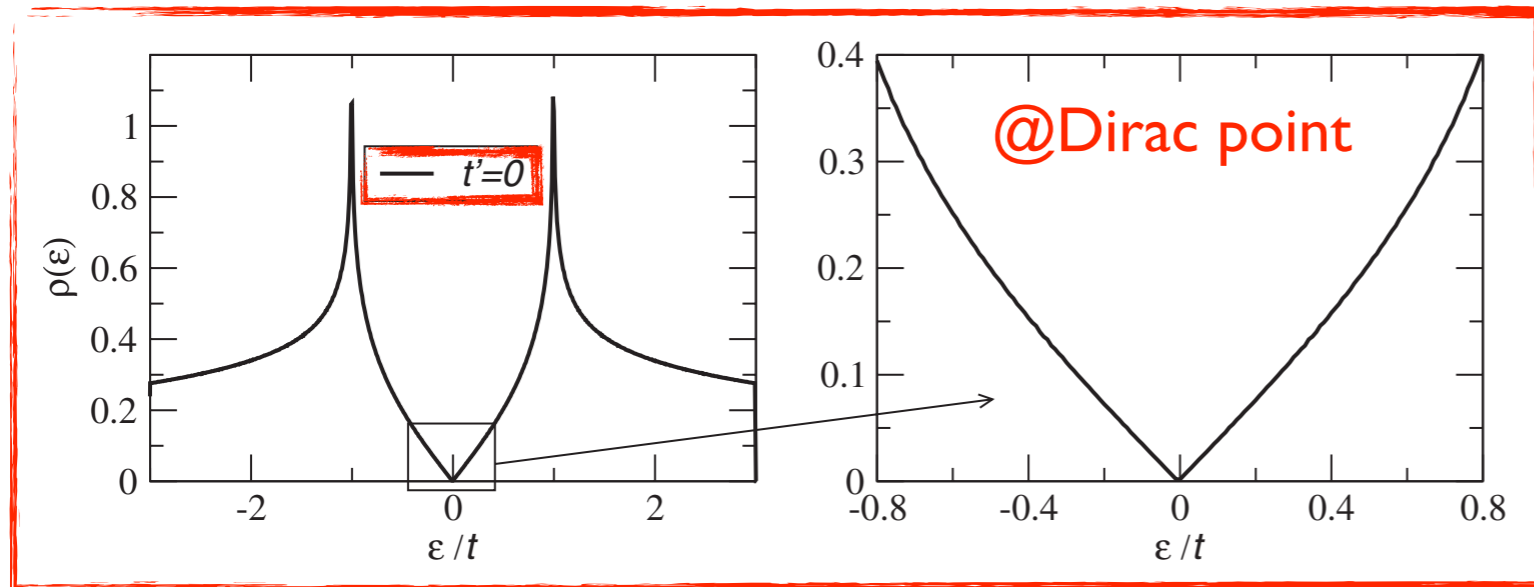
- Unprecedented phenomena in condensed matter:

- *half integer quantum Hall effect*
- *Klein paradox and suppression of backscattering*



# Density of states

- Density of states:  $\rho(\epsilon)$

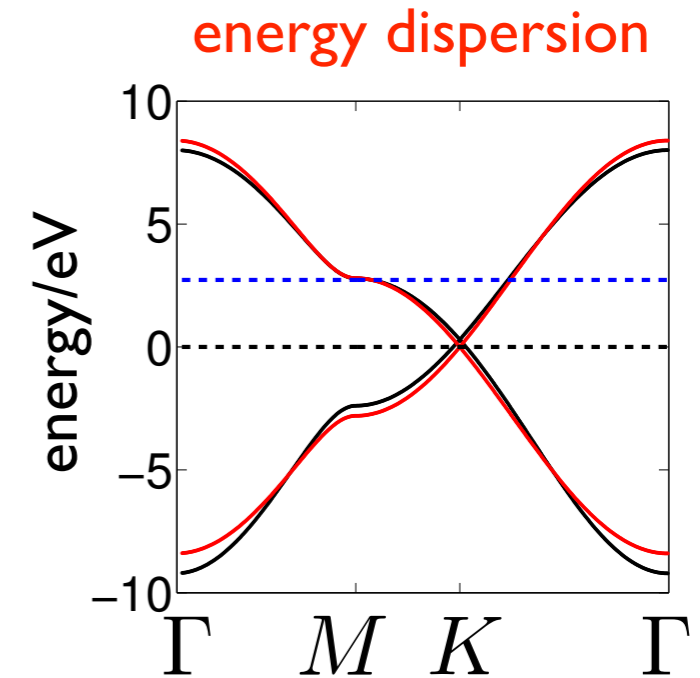
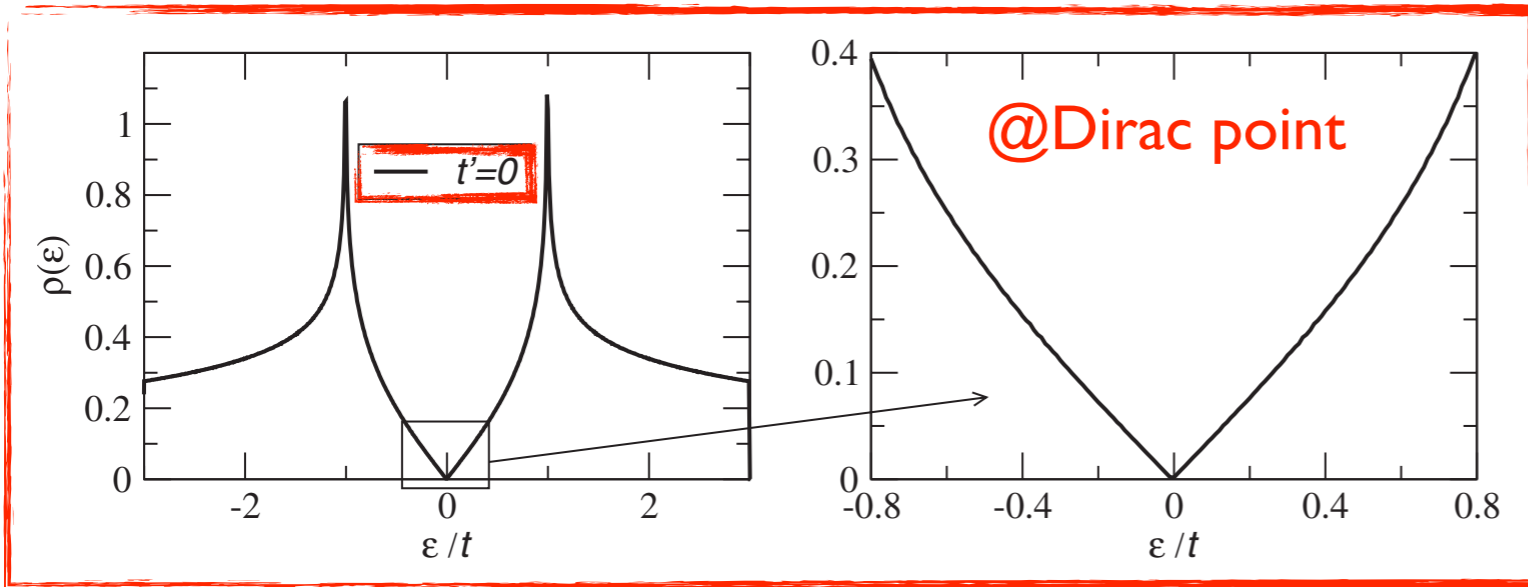


Kiesel et al, Phys. Rev. B 86, 020507 (2012)

- ▶ close to Dirac point:  $\rho(\epsilon) \propto |\epsilon|$
- ▶ van-Hove singularities @ finite doping  $\rightarrow$  logarithmic divergence of  $\rho(\epsilon)$

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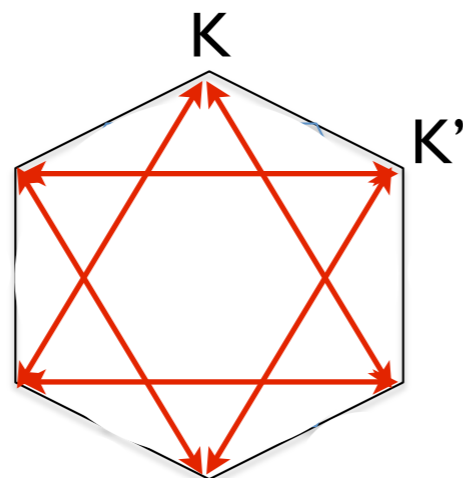


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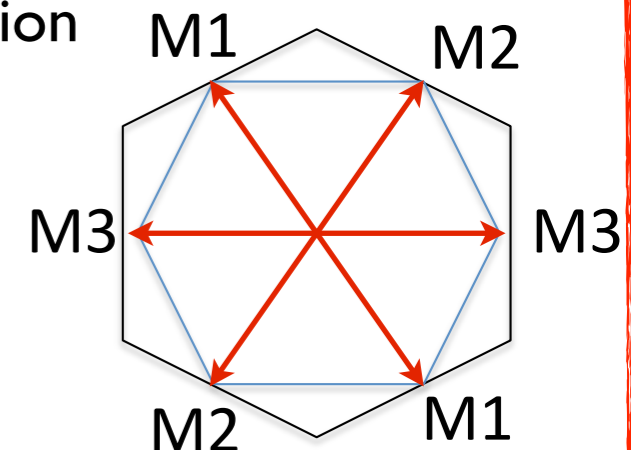
## • Graphene @ neutrality point:

- ▶ Dirac fermions
- ▶ nesting
- ▶ vanishing DOS



## • Graphene @ van-Hove singularity:

- ▶ saddle-point dispersion
- ▶ nesting
- ▶ diverging DOS



# Interactions & phase transitions

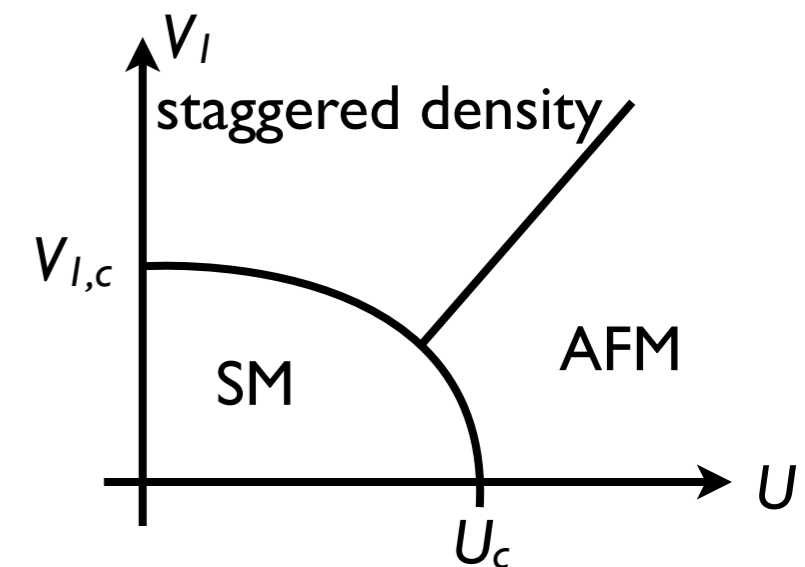
$$H_{\text{int}} = U \sum_i n_{i,\uparrow} n_{i,\downarrow} + V_1 \sum_{\langle i,j \rangle, \sigma, \sigma'} n_{i,\sigma} n_{j,\sigma'} + V_2 \sum_{\langle\langle i,j \rangle\rangle, \sigma, \sigma'} n_{i,\sigma} n_{j,\sigma'} + V_3 \sum_{\langle\langle\langle i,j \rangle\rangle\rangle, \sigma, \sigma'} n_{i,\sigma} n_{j,\sigma'}$$

	Graphene		Graphite	
	Bare	cRPA	Bare	cRPA
$U_{00}^{A \text{ or } B}$ (eV)	17.0	9.3	17.5, 17.7	8.0, 8.1
$U_{01}$ (eV)	8.5	5.5	8.6	3.9
$U_{02}^{A \text{ or } B}$ (eV)	5.4	4.1	5.4, 5.4	2.4, 2.4
$U_{03}$ (eV)	4.7	3.6	4.7	1.9

☞ Wehling et al, Phys. Rev. Lett. **106**, 236805 (2011)

- **Undoped graphene** → interactions can induce phase transitions:

- \* Dirac fermions have vanishing DOS at Fermi level
- \* Stoner-type criterion → **critical interaction strength** required
- \* Experimental data: Graphene below critical strength

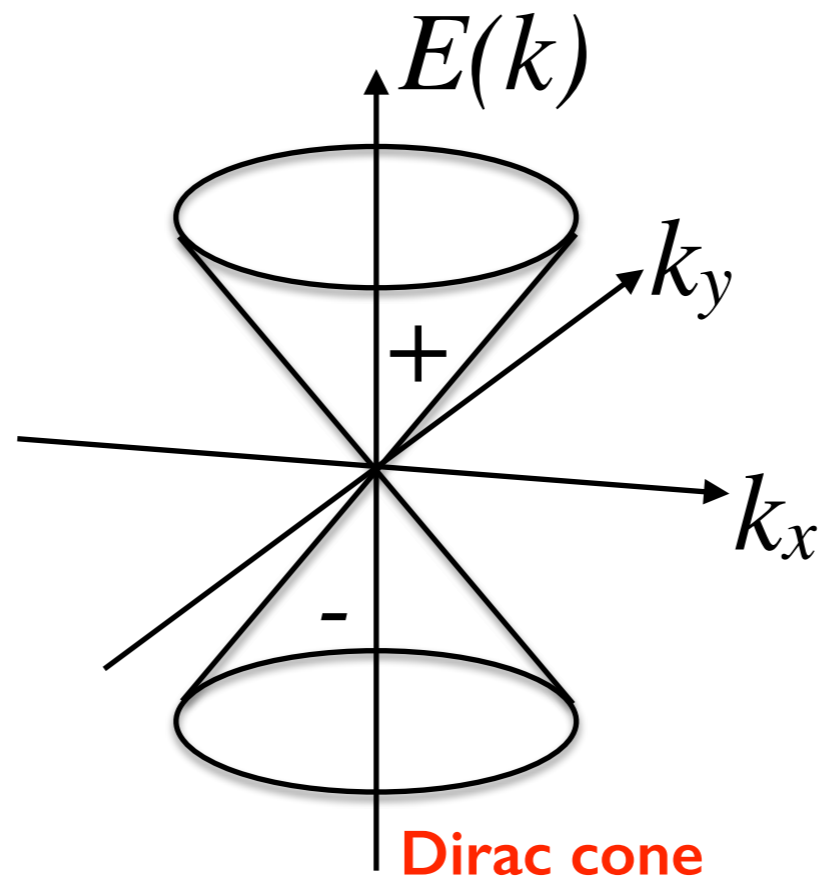


- **Graphene @ VHS** → interactions can induce phase transitions:

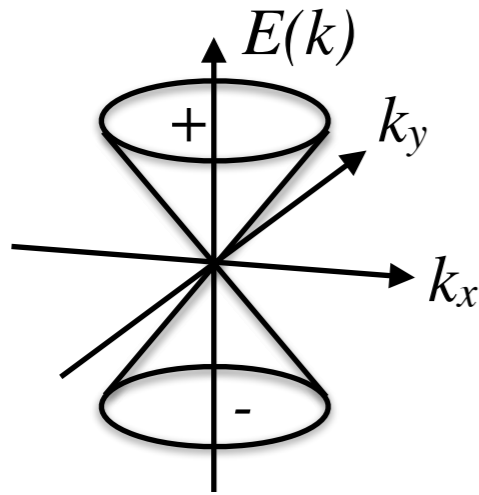
- \* nesting: ph-channel diverges @ low T
- \* also: pp-channel diverges @ low T

➡ **competing orders**

## (F)RG aspects of undoped graphene



# Symmetries & Fierz transformations



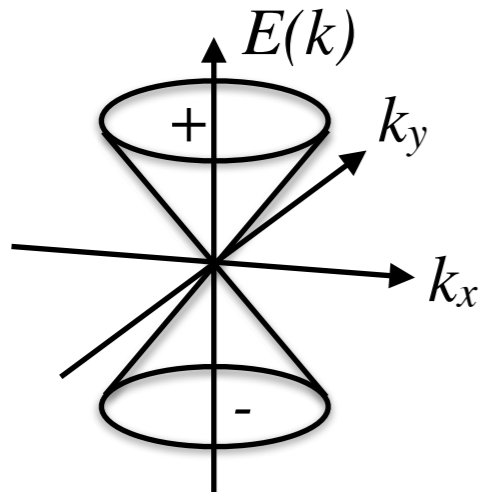
- **Minimal description** of i.a. electrons in graphene @ low energies?
- **Symmetries** of continuum interacting theory?
- What kinds of order expected? Nature of **phase transitions**?

• Start with linearized non-interacting Lagrangian for Dirac electrons:  $L_0 = \bar{\Psi}(\vec{x}, \tau) \gamma_\mu \partial_\mu \Psi(\vec{x}, \tau)$

• **Interactions** (quartic, local, ... and spinless):  $L_{int} = (\Psi^\dagger(\vec{x}, \tau) M_1 \Psi(\vec{x}, \tau)) (\Psi^\dagger(\vec{x}, \tau) M_2 \Psi(\vec{x}, \tau))$

4d hermitian matrices

# Symmetries & Fierz transformations



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4d hermitian matrices

➡ **136 independent couplings in spinless case!**

• Honeycomb lattice symmetries help to **reduce to 15 independent couplings**

• Further reduction by Fierz identities!

➡ **6 independent couplings in spinless case!**

# Symmetries, Fierz transformations & order parameters

- Reduction of independent couplings can be pushed further
  - ▶ near quantum critical point additional symmetries from non-interacting theory restored...
  - ▶ minimal low-energy description of i.a. spinless honeycomb electrons:

## Symmetries:

- ▶ Reflection
- ▶ Translation
- ▶ Time-reversal
- ▶ Rotations
- ▶ Lorentz
- ▶ Chiral

$$L = L_0 + g_{D2}(\bar{\Psi}\gamma_{35}\Psi)^2 + g_{C1}(\bar{\Psi}\Psi)^2$$

▶ order parameter: QAH

▶ order parameter: CDW



# FRG phase transitions & critical behavior

- Situation with spin works similarly - e.g. obtain AF-SDW order and CDW order
- Introduce **order parameter fields** via Hubbard-Stratonovich transformation

$$\Gamma_k = \int d^D x \left[ Z_{\Psi,k} \bar{\Psi} (\mathbb{1}_2 \otimes \gamma_\mu) \partial_\mu \Psi - \frac{1}{2} Z_{\phi,k} \phi_a \overset{\substack{\text{order parameter field } Z_2 \text{ or } SO(3)}}{\downarrow} \partial_\mu^2 \phi_a \right. \\ \left. + U_k(\rho) + \bar{g}_k \phi_a \bar{\Psi} (\sigma_a \otimes \mathbb{1}_4) \Psi \right],$$

# FRG phase transitions & critical behavior

- Situation with spin works similarly - e.g. obtain AF-SDW order and CDW order
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order parameter field  $Z_2$  or  $SO(3)$

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& have fun:

$$\partial_t \Gamma_k = \frac{1}{2} \text{STr} \left[ \partial_t R_k \left( \Gamma_k^{(2)} + R_k \right)^{-1} \right]$$

► quantify critical behaviour at quantum phase transition ( $Z_2$ ):

	$1/\nu$	$\eta_\phi$	$\eta_\Psi$
FRG [LPA', $\mathcal{O}(\tilde{\rho}^6)$ , $R_k^{\text{lin}}$ ]	0.982	0.760	0.032
FRG [LPA', $\mathcal{O}(\tilde{\rho}^6)$ , $R_k^{\text{sc}}$ ]	0.978	0.767	0.033
FRG [LPA', full $u(\tilde{\rho})$ , $R_k^{\text{lin}}$ ] <sup>21</sup>	0.982	0.756	0.032
$1/N_f$ -expansion (2nd/3rd order) <sup>34,35</sup>	0.962*	0.776	0.044
$(2 + \epsilon)$ -expansion (3rd order) <sup>28</sup>	0.764	0.602	0.081
$(4 - \epsilon)$ -expansion (2nd order) <sup>13,33</sup>	1.055	0.695	0.065
Polynomial interpolation $P_{2,2}$	0.995	0.753	0.034
Polynomial interpolation $P_{3,2}$	0.949	0.716	0.041
Monte-Carlo simulations <sup>33†</sup>	1.00(4)	0.754(8)	—

► Graphene: AFM-Mott transition

► “chiral Heisenberg universality class”

➡ talk by L. Janssen

Janssen & Herbut, Phys. Rev. B **89**, 205403 (2014)

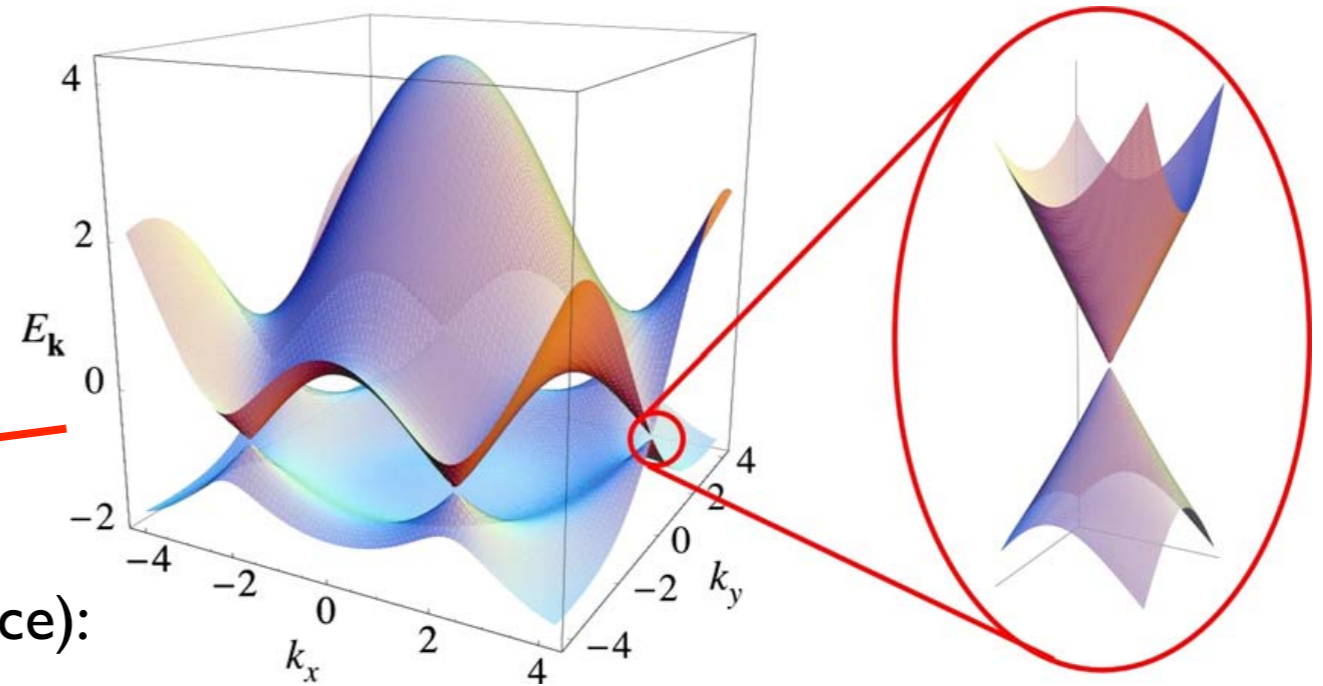
Janssen & Gies, Phys. Rev. D **86**, 105007 (2012)

Mesterhazy et al., Phys. Rev. B **86**, 245431 (2012)

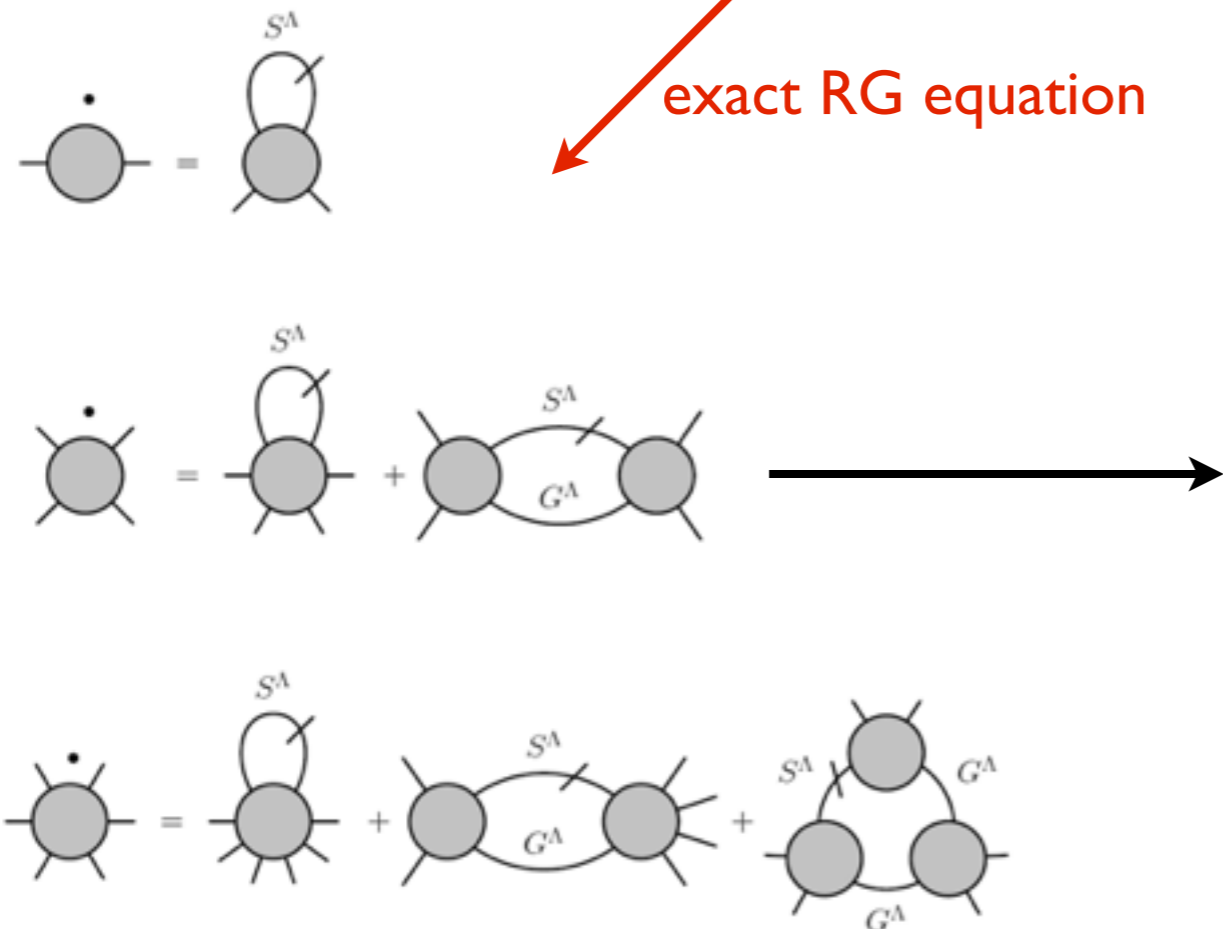
# N-Patch FRG scheme

- Take into account full band structure
- Modified bare propagator with IR cutoff:

$$G_0^\Lambda(k_0, \mathbf{k}) = \frac{\theta^\Lambda(\mathbf{k})}{ik_0 - \xi_{\mathbf{k}}}$$



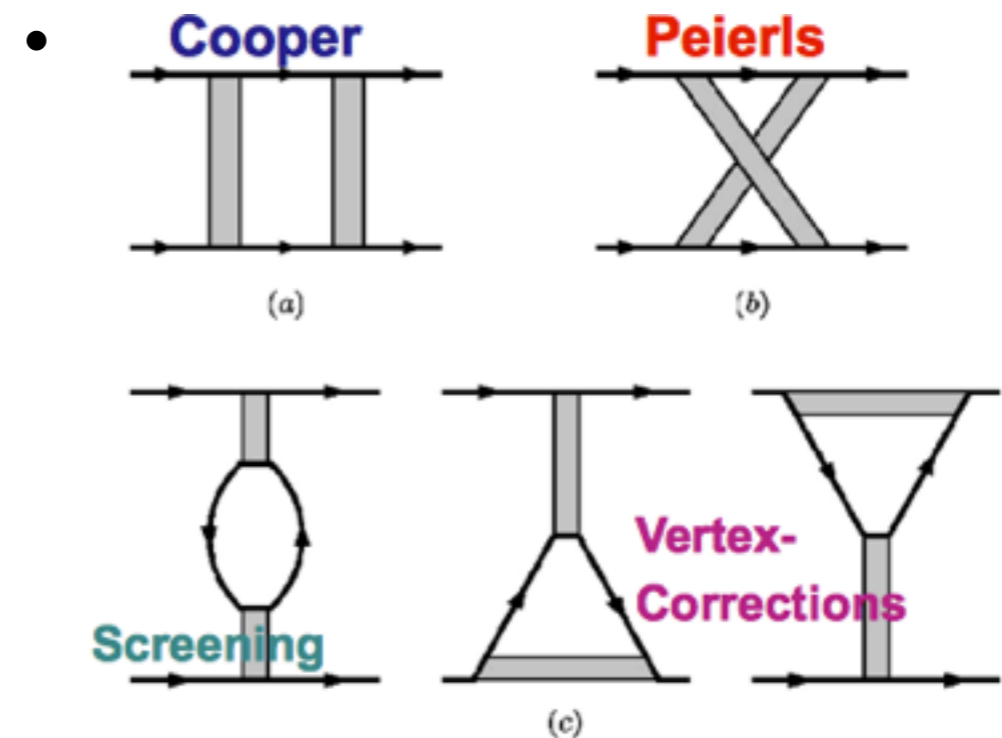
- Vertex expansion (with momentum dependence):



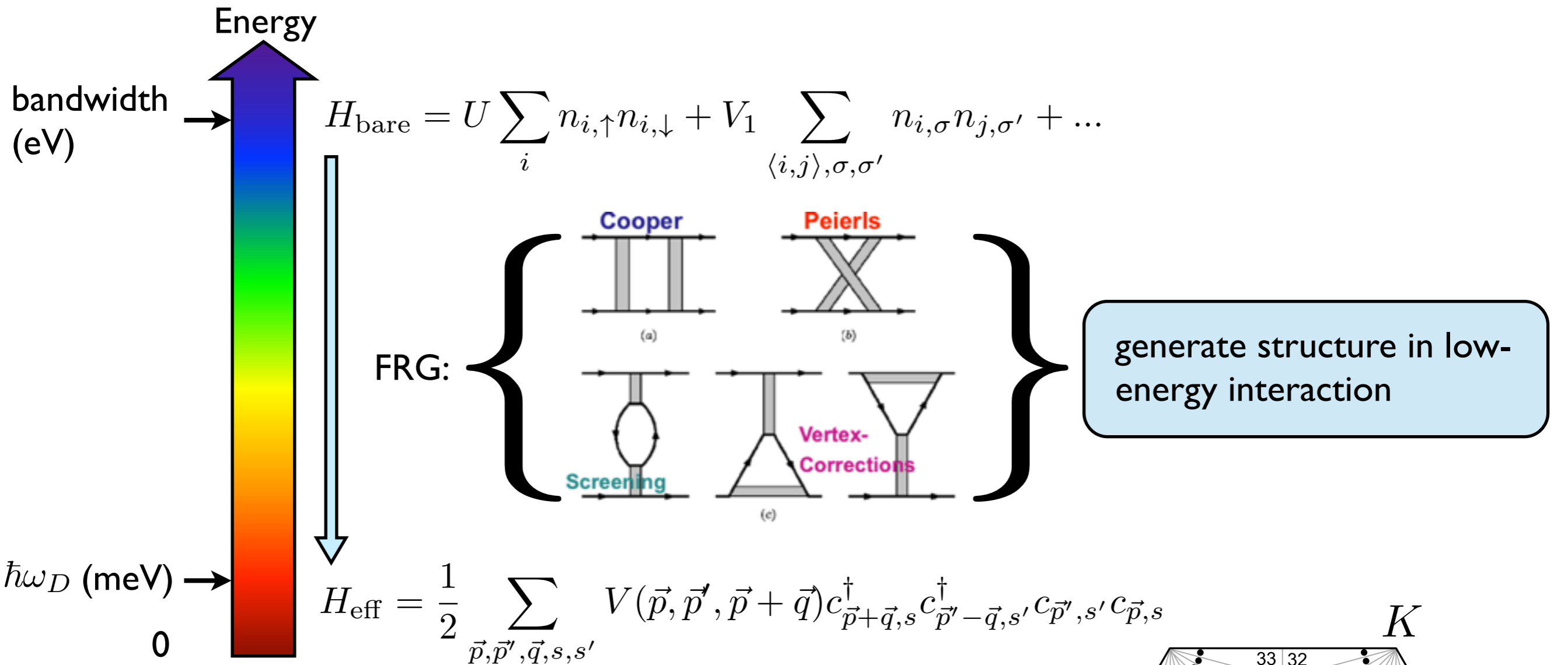
... Salmhofer & Honerkamp, Prog.Theor. Phys. 105 (2001)

Metzner et al., Rev. Mod. Phys. 84, 299 (2012)

- Neglect 6-point and higher vertices
- No self-energy feedback

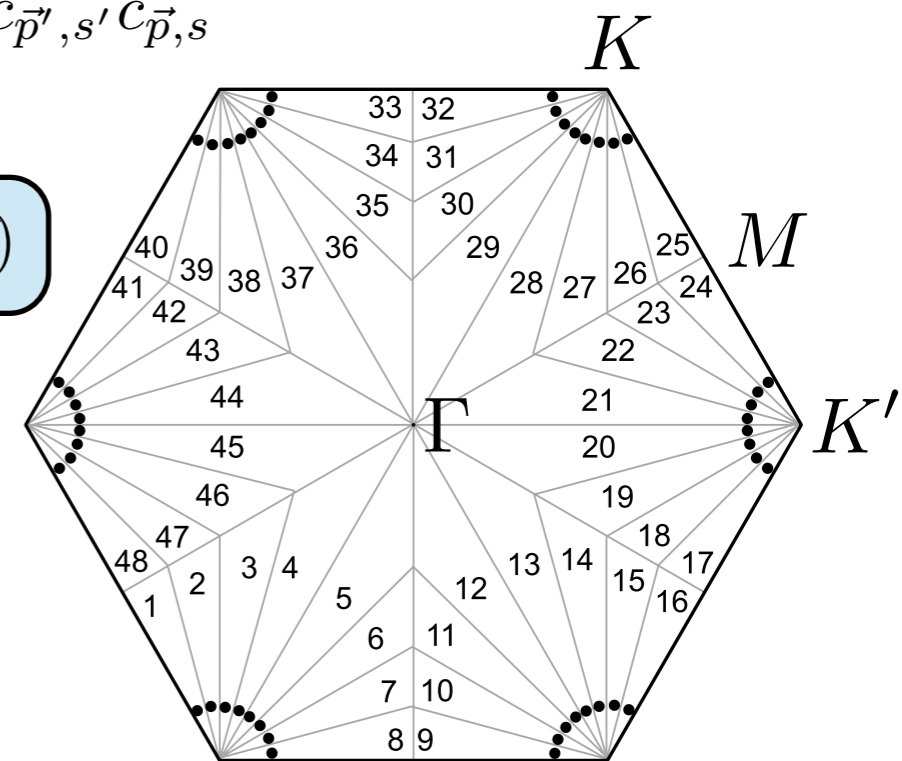


# N-Patch FRG scheme

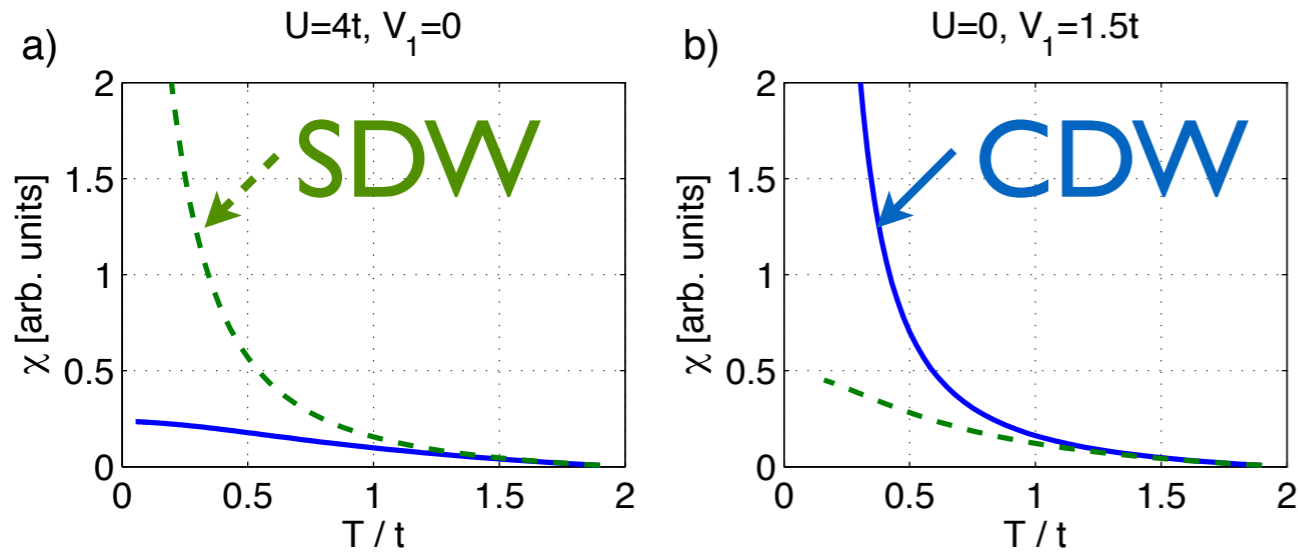


► low-energy effective action & momentum structure  $V(\vec{p}, \vec{p}', \vec{p} + \vec{q})$

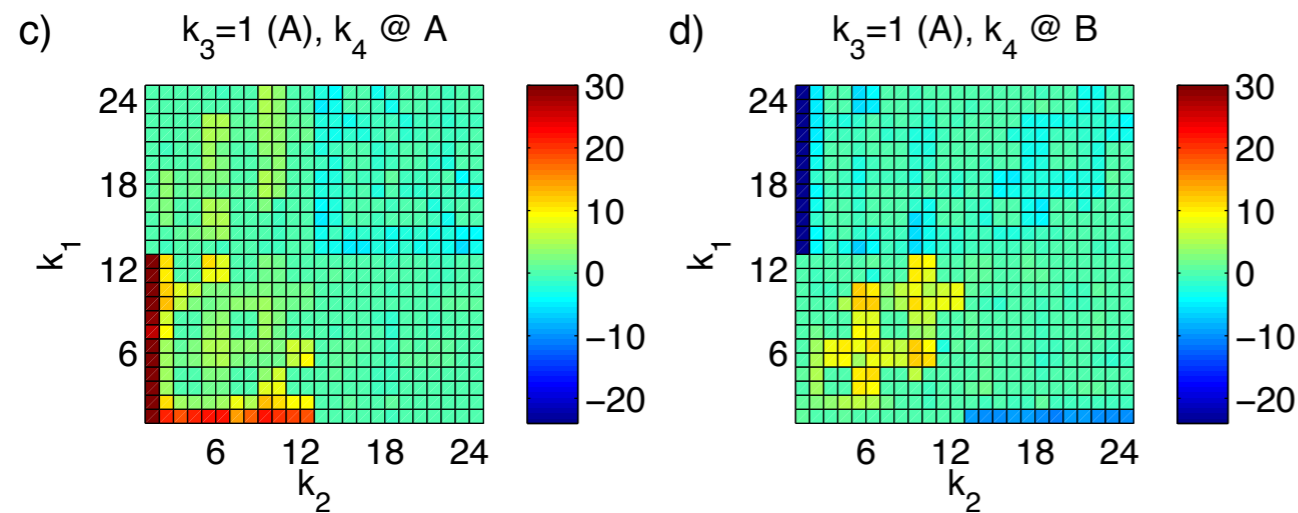
- Wavevector: discretized by patching of Brillouin zone
- Interaction constant within one patch



# Density waves on the honeycomb lattice

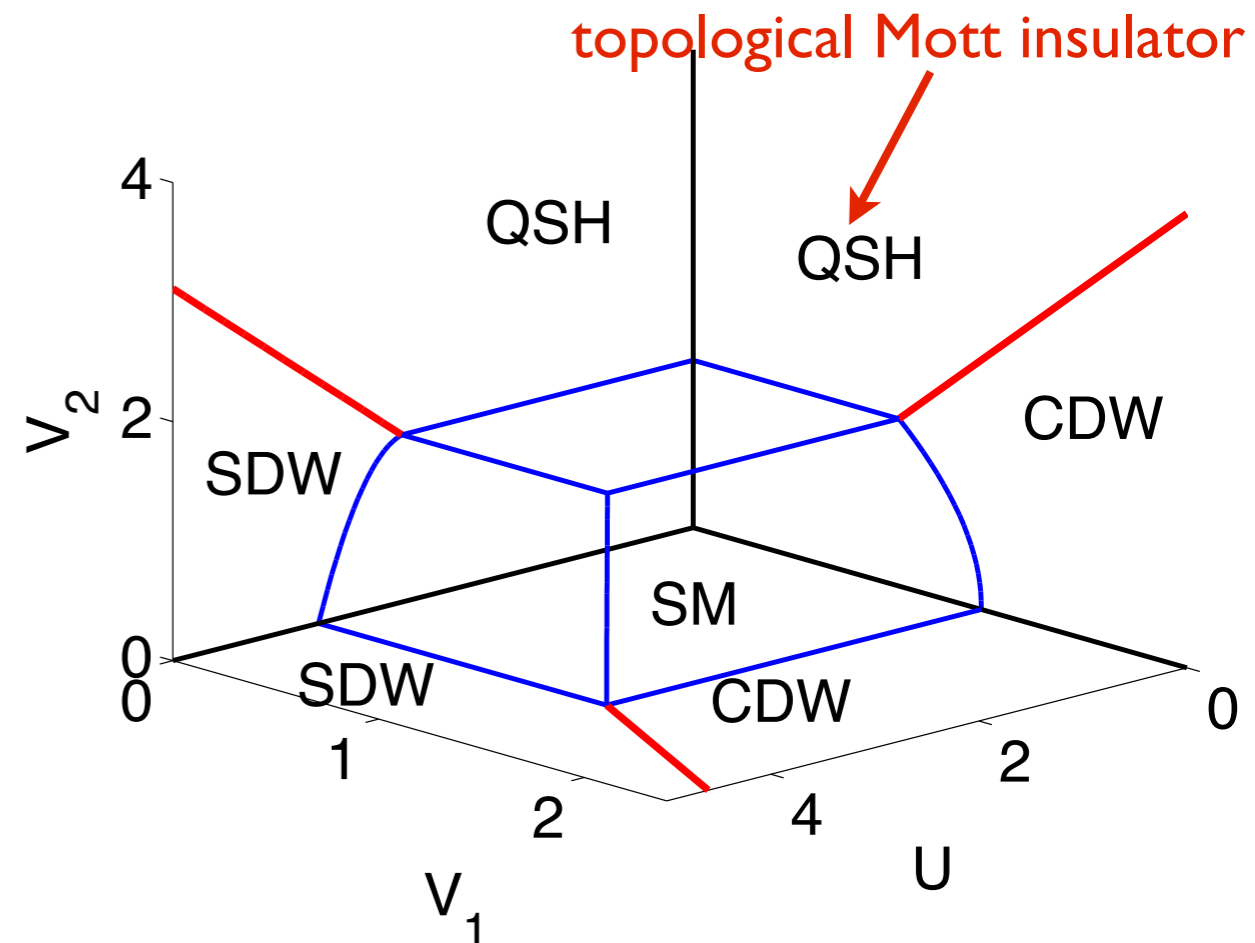


$$H_{\text{int}} = U \sum_i n_{i,\uparrow} n_{i,\downarrow} + V_1 \sum_{\langle i,j \rangle, \sigma, \sigma'} n_{i,\sigma} n_{j,\sigma'} + V_2 \sum_{\langle\langle i,j \rangle\rangle, \sigma, \sigma'} n_{i,\sigma} n_{j,\sigma'}$$



$$\sim \sum_{\vec{q}} J_{\vec{q}}^{b,b'} \vec{S}_{\vec{q}}^b \cdot \vec{S}_{-\vec{q}}^{b'}$$

## • Phase diagram:

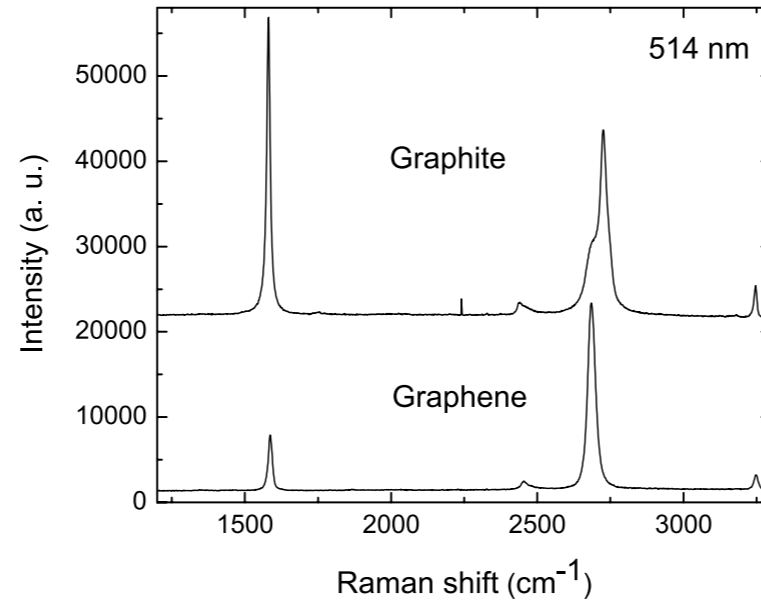


Honerkamp, Phys. Rev. Lett. 100, 146404 (2008)

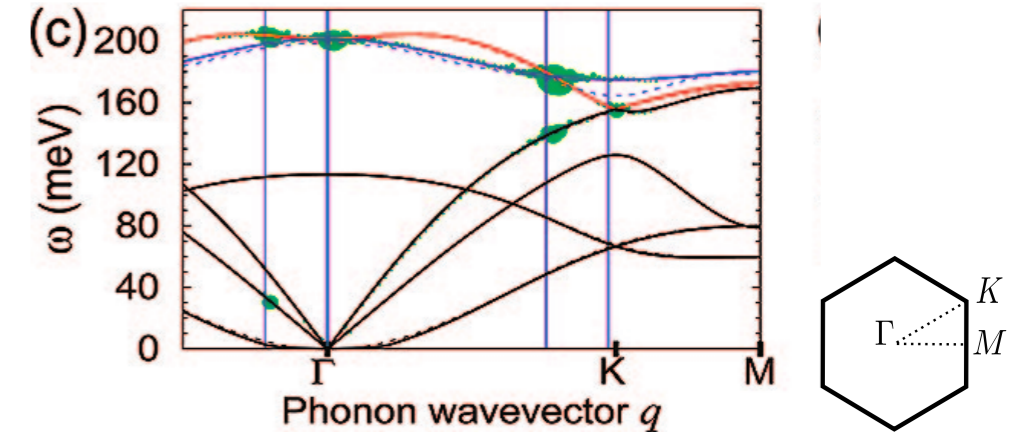
Raghu et al., Phys. Rev. Lett. 100, 156401 (2008)

# Phonon-induced electron-electron interactions

- Lattice vibrations (Raman spectroscopy)

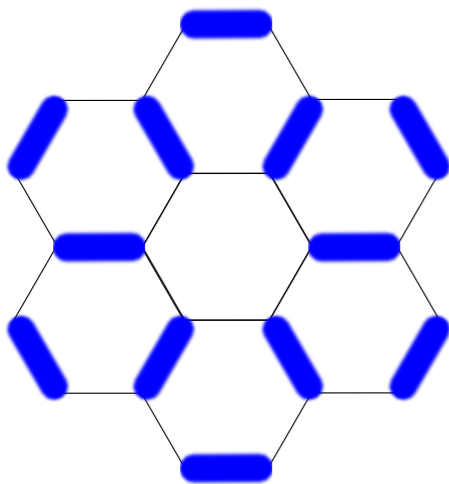


A. C. Ferrari *et al.* PRL **97** (2006)

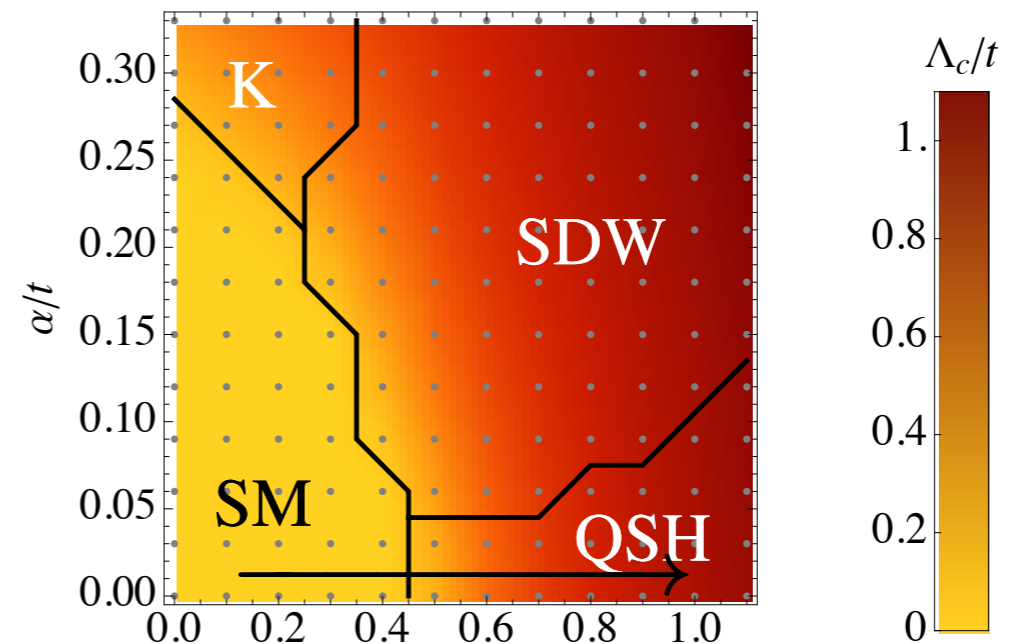


C. Park *et al.* NL **8** (2008)

- Modified hopping  $t \rightarrow t+dt$   
 $\rightarrow$  EPC coupling induces Kekule instability:



- Phase diagram with short-ranged Coulomb i.a.:

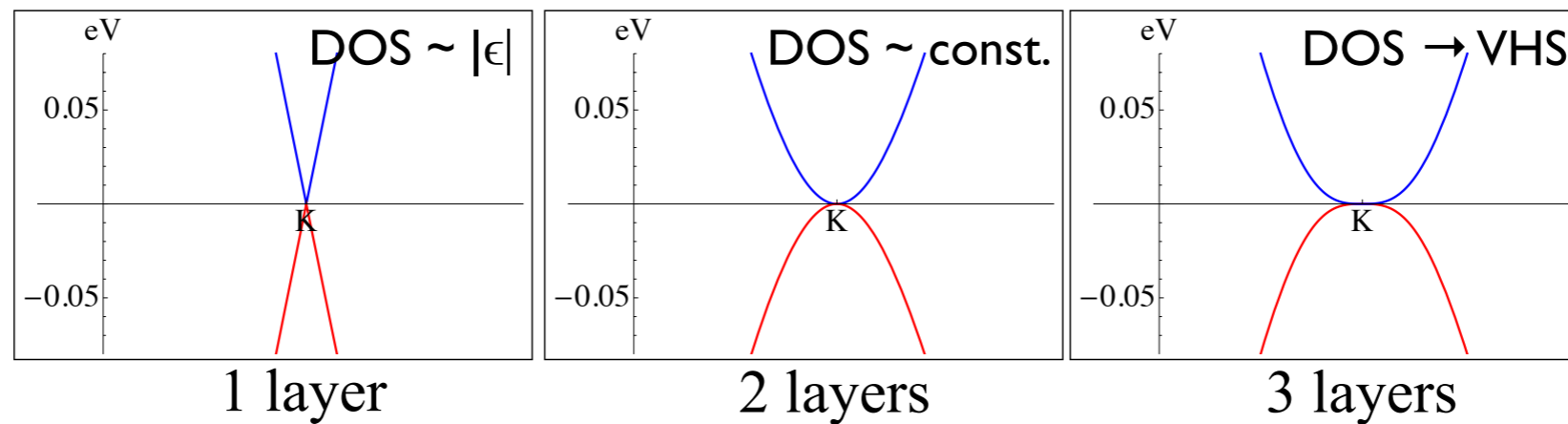


$$\{U/t, V_1/t, V_2/t\} \approx \{3.3, 2.0, 1.5\}$$

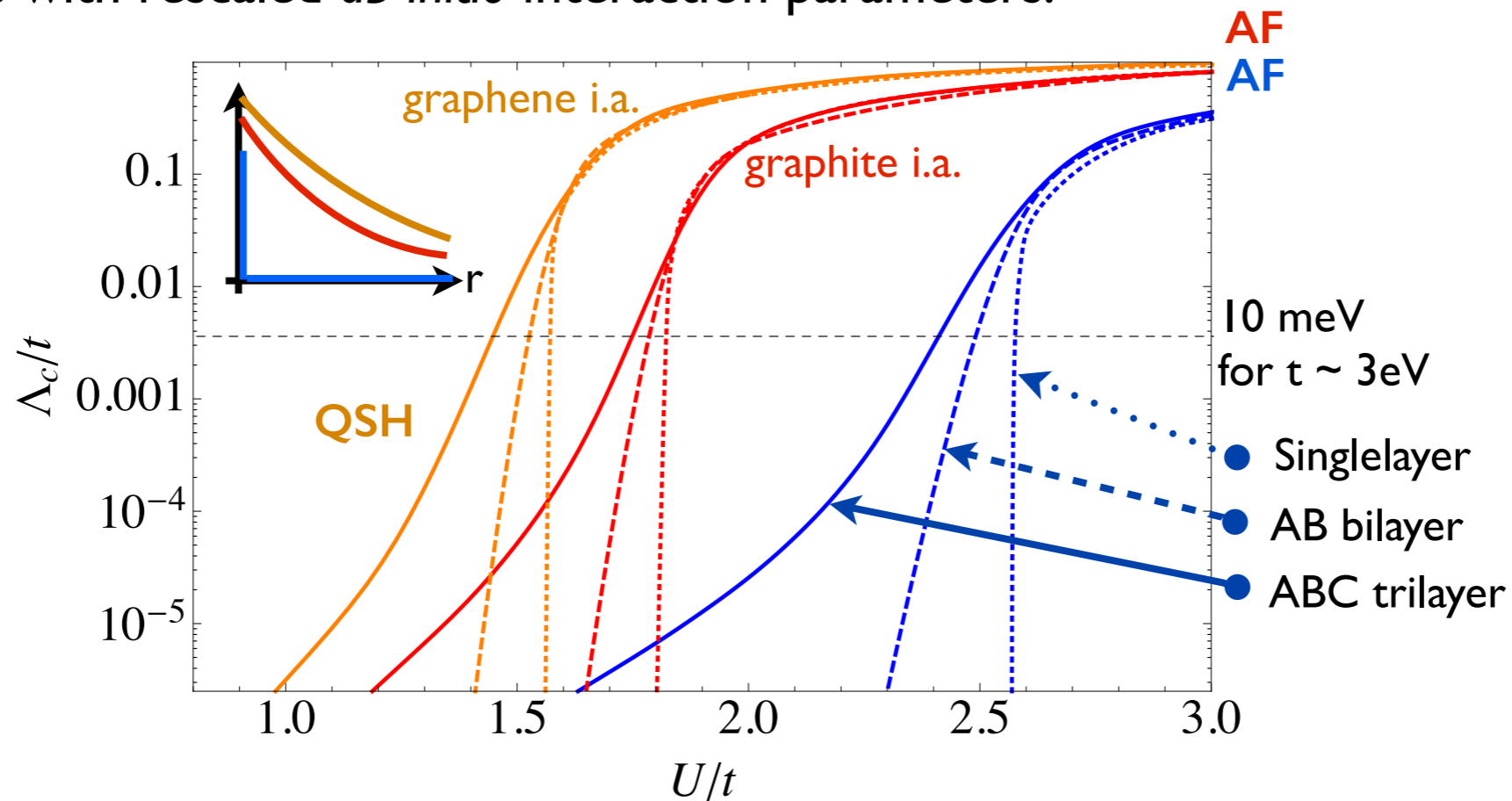
$$\rightarrow c \{U/t, V_1/t, V_2/t\}$$

$\rightarrow$  talk by L. Classen

# Critical scales of honeycomb stacks

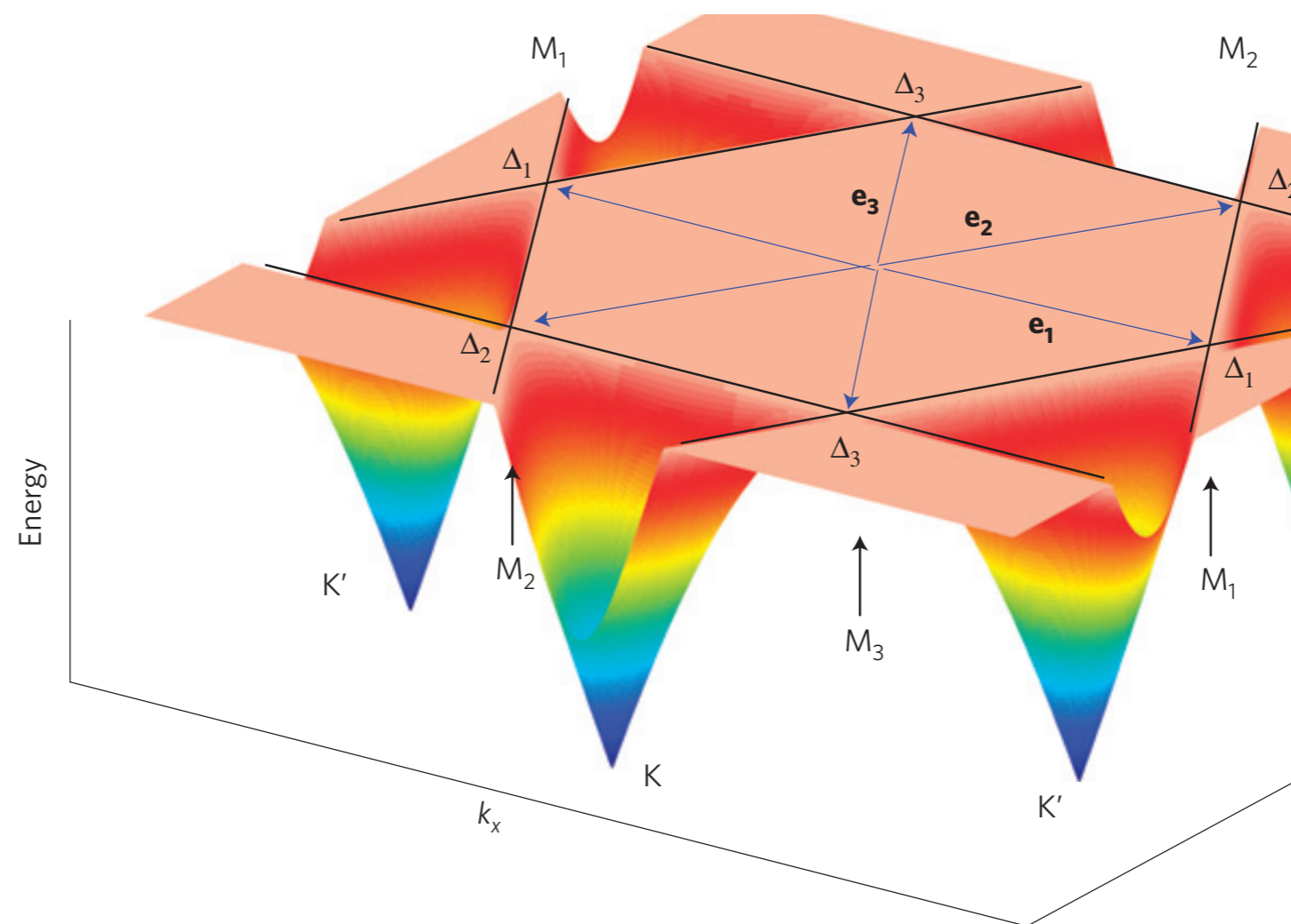


- Critical scales with rescaled *ab initio* interaction parameters:



- ABC trilayer most prone to instabilities, critical scale drops quickly when  $U < U_{c, \text{Singlelayer}}$
- *Ab initio* parameters put system close to QSH/AFM phase boundary

# (F)RG aspects of graphene @ VHS





# RG for graphene @VHS - g-ology

$$\mathcal{L} = \sum_{\alpha=1}^3 \psi_{\alpha}^{\dagger} (\partial_{\tau} - \epsilon_{\mathbf{k}} + \mu) \psi_{\alpha} - \frac{1}{2} g_4 \psi_{\alpha}^{\dagger} \psi_{\alpha}^{\dagger} \psi_{\alpha} \psi_{\alpha}$$

$$- \sum_{\alpha \neq \beta} \frac{1}{2} [g_1 \psi_{\alpha}^{\dagger} \psi_{\beta}^{\dagger} \psi_{\alpha} \psi_{\beta} + g_2 \psi_{\alpha}^{\dagger} \psi_{\beta}^{\dagger} \psi_{\beta} \psi_{\alpha} + g_3 \psi_{\alpha}^{\dagger} \psi_{\alpha}^{\dagger} \psi_{\beta} \psi_{\beta}]$$

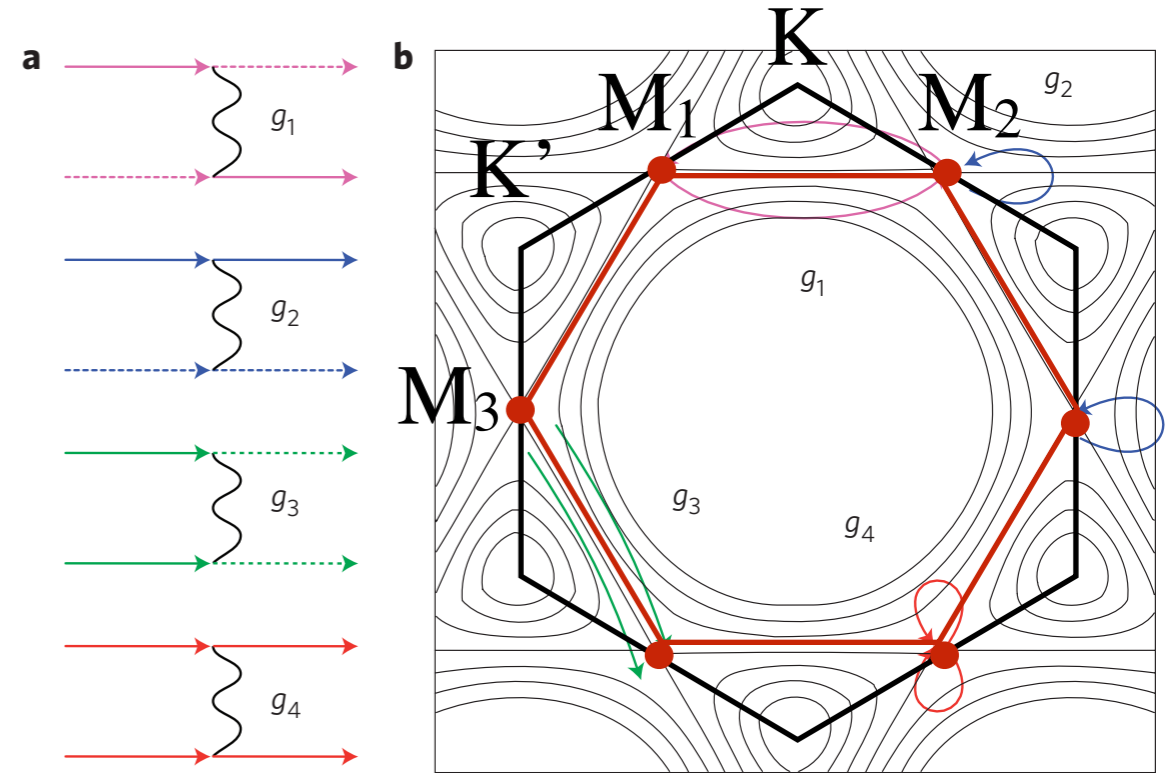
- RG equations:

$$\frac{dg_1}{dy} = 2d_1 g_1 (g_2 - g_1), \quad \frac{dg_2}{dy} = d_1 (g_2^2 + g_3^2)$$

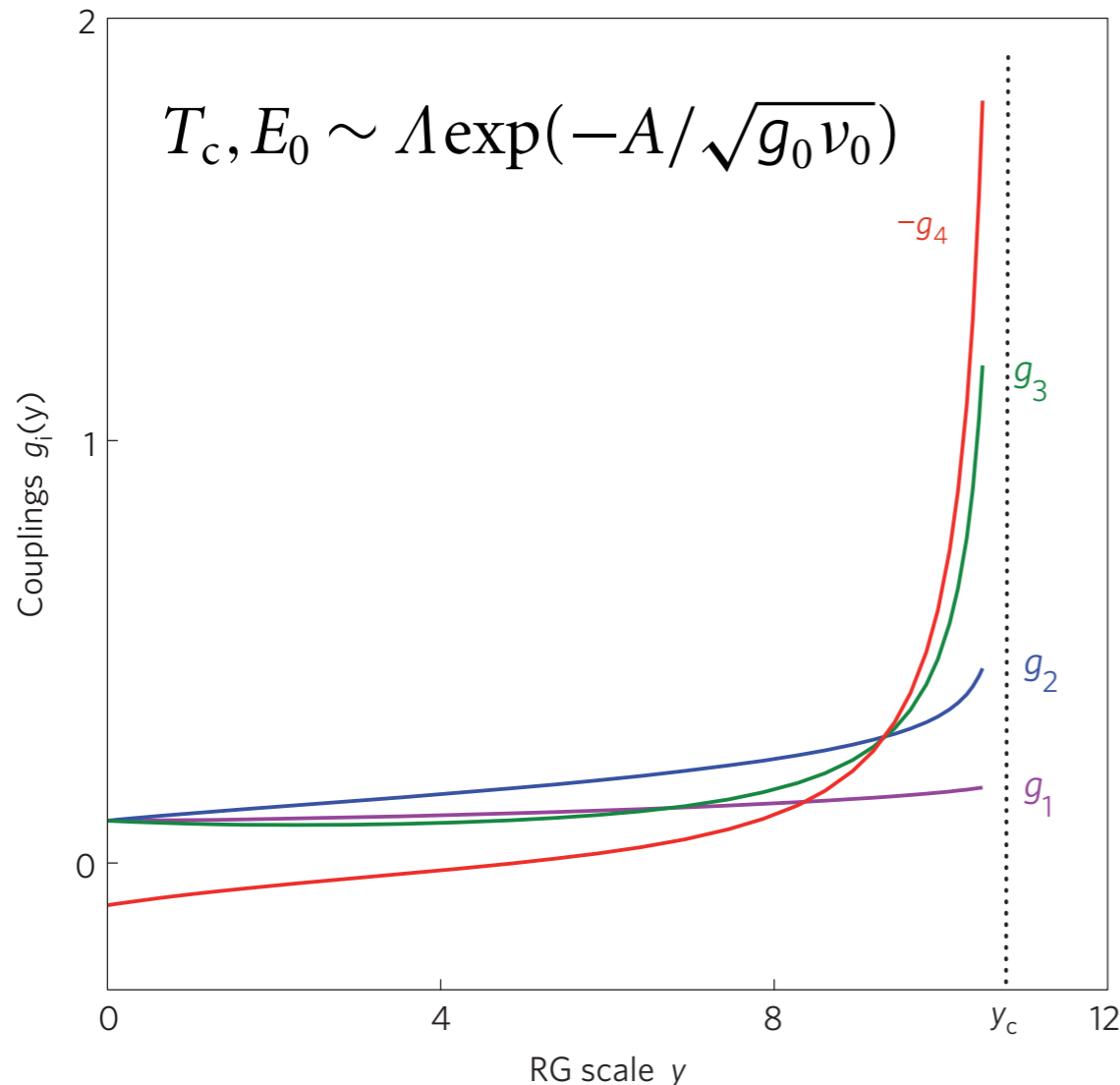
$$\frac{dg_3}{dy} = -(n-2)g_3^2 - 2g_3 g_4 + 2d_1 g_3 (2g_2 - g_1),$$

$$\frac{dg_4}{dy} = -(n-1)g_3^2 - g_4^2$$

- ▶ reproduce two-patch RG for n=2
- ▶ Graphene at VHS needs n=3



# RG for graphene @VHS - g-ology



- Susceptibilities (introduce test vertices)

$$\delta\mathcal{L} = \sum_{\alpha=1}^3 \tilde{\Delta}_{\alpha} \psi_{\alpha,\uparrow}^{\dagger} \psi_{\alpha,\downarrow}^{\dagger} + \tilde{\Delta}_{\alpha}^* \psi_{\alpha,\uparrow} \psi_{\alpha,\downarrow}$$

- d-wave SC susceptibility

$$\chi_{\text{dSC}}(y) \sim (y_c - y)^{-1.5}$$

- AF-SDW susceptibility @perfect nesting

$$\chi_{\text{SDW}}(y) \sim (y_c - y)^{-1}$$

- ▶ Close competition between SDW and dSC
- ▶ dSC is leading instability for all values of nesting
  - ➡ in contrast to square lattice (SDW@perfect nesting)

# FRG for graphene @VHS - N-Patch FRG

- Full band structure & realistic model parameters:

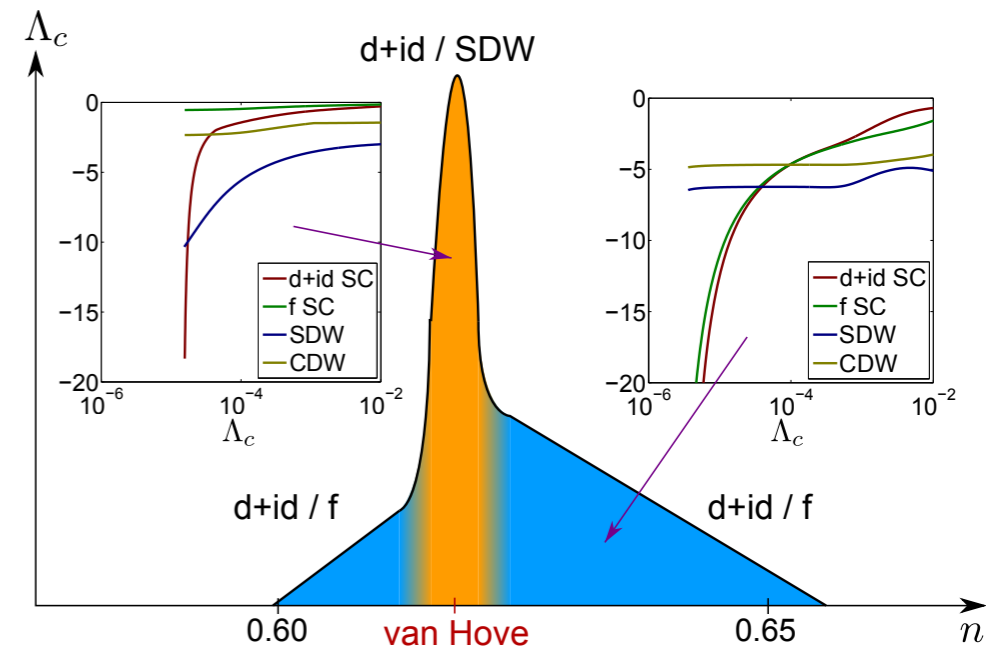
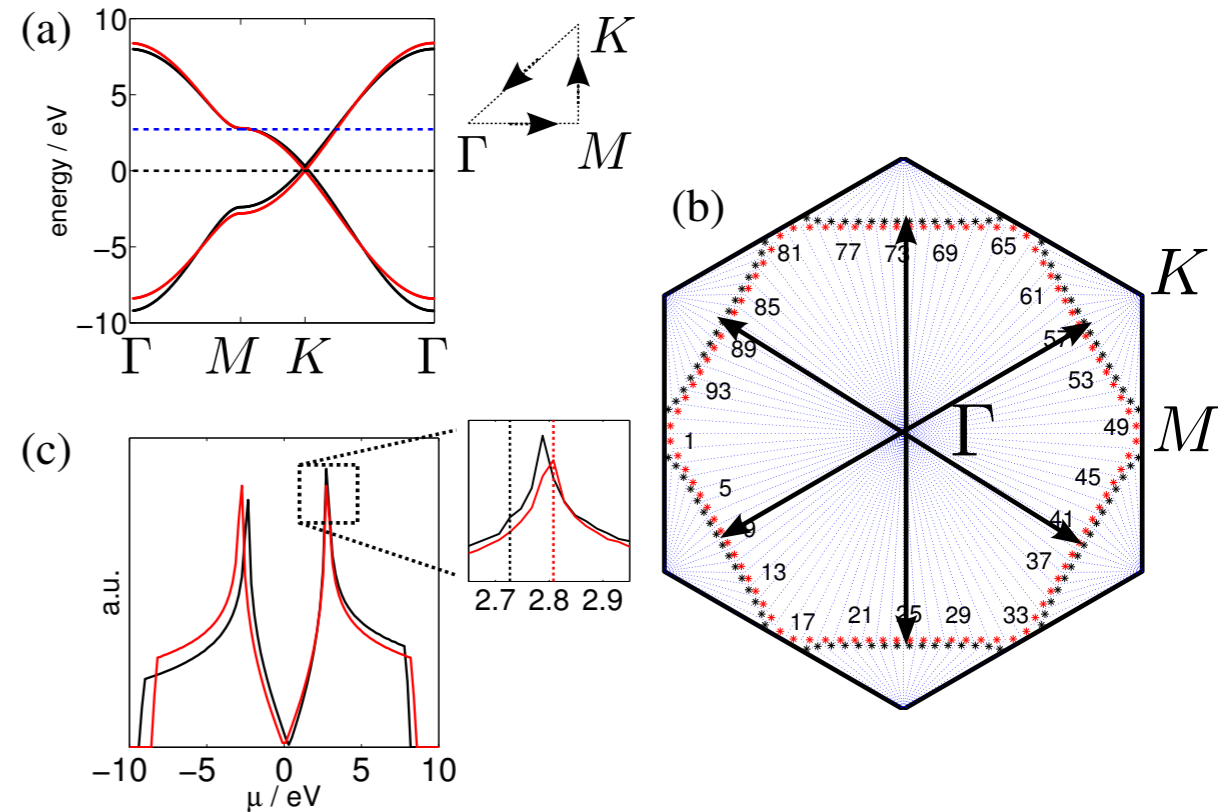
$$H_0 = \left[ t_1 \sum_{\langle i,j \rangle} \sum_{\sigma} c_{i,\sigma}^{\dagger} c_{j,\sigma} + t_2 \sum_{\langle\langle i,j \rangle\rangle} \sum_{\sigma} c_{i,\sigma}^{\dagger} c_{j,\sigma} + t_3 \sum_{\langle\langle\langle i,j \rangle\rangle\rangle} \sum_{\sigma} c_{i,\sigma}^{\dagger} c_{j,\sigma} + \text{h.c.} \right] - \mu n,$$

- Interaction terms:

	Graphene		Graphite	
	Bare	cRPA	Bare	cRPA
$U_{00}^{A \text{ or } B}$ (eV)	17.0	9.3	17.5, 17.7	8.0, 8.1
$U_{01}$ (eV)	8.5	5.5	8.6	3.9
$U_{02}^{A \text{ or } B}$ (eV)	5.4	4.1	5.4, 5.4	2.4, 2.4
$U_{03}$ (eV)	4.7	3.6	4.7	1.9

ⓘ Wehling et al, Phys. Rev. Lett. **106**, 236805 (2011)

- longer-range hoppings decrease degree of nesting
- dSC wins
- with realistic parameters:  $T_c \sim$  a few K



ⓘ Kiesel, Platt, Hanke, Abanin, Thomale, Phys. Rev. B **86**, 020507 (2012)

ⓘ Wang et al, Phys. Rev. B **85**, 035414 (2012)

# Conclusions & Outlook

**Graphene allows for beautiful/useful/complex/  
unprecedented/exotic theory!**

- Phase transitions and criticality @Dirac point
  - ▶ precision estimates - dynamical bosonization, higher-derivative terms? multicriticality?
- Phase transitions and criticality @VHS - realization of chiral d-wave superconductor?
  - ▶ coupling to the lattice
  - ▶ self-energy effects
  - ▶ van Hove situation difficult to assess by other methods (finite-size, sign problems)
  - ▶ collective fluctuations, cf. Krahl, Friederich, Wetterich on Hubbard model
- Methods & related materials:
  - ▶ spin-SU(2) - extend scheme to include SOC ➡ **talks by D. Scherer & G. Schober**
  - ▶ multiorbital models ➡ **talk by C. Platt**