

# Yang-Mills correlation functions from Dyson-Schwinger equations

UNI  
GRAZ



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7th International Conference on the Exact Renormalization Group, Lefkada

**HIC** for **FAIR**  
Helmholtz International Center



Investigate:

- ▶ Confinement
- ▶ Dynamical symmetry breaking
- ▶ Bound states

(Non-perturbative) Determination of correlation functions:

- ▶ Lattice
- ▶ Effective theories, e.g., (refined) Gribov-Zwanziger, massive Yang-Mills
- ▶ Functional equations
- ▶ ...

Functional equations...

- ▶ ... are exact.
- ▶ ... form an infinite system of equations.

Truncating the system: How close/far from exact solution?

- ▶ Calculate physical quantities.
- ▶ Compare to other methods.
- ▶ Modify the truncation.

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- ▶ Modify the truncation.

Recent results indicate that **primitively divergent correlation functions** of Landau gauge Yang-Mills theory provide reasonable truncation for a **quantitative, self-consistent and self-contained description**.

# Comparison: DSEs and flow equations

Dyson-Schwinger equations (DSEs)	Functional RG equations (FRGEs)
'integrated flow equations'	'differential DSEs'
effective action $\Gamma[\phi]$	effective average action $\Gamma^k[\phi]$
—	regulator
n-loop structure ( $n \text{ const.}$ )	1-loop structure
exactly only bare vertex per diagram	no bare vertices

$$\frac{\partial}{\partial \phi} \Gamma[\phi] = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \text{Diagram 4} + \text{Diagram 5} + \dots$$
$$k \frac{\partial}{\partial k} \Gamma^k[\phi] = \text{Diagram 1}$$

- Both **systems of equations** are **exact**
- Both contain infinitely many equations.

# Outline

- ▶ Vacuum
  - ▶ Propagators: Introduction
  - ▶ Four-gluon vertex
- ▶ Non-vanishing temperature
  - ▶ Ghost-gluon vertex
  - ▶ Three-gluon vertex

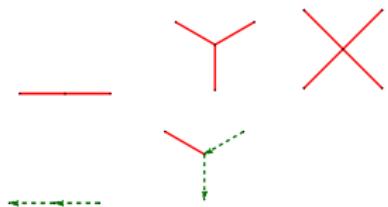
Gluonic sector of quantum chromodynamics: Yang-Mills theory

$$\mathcal{L} = \frac{1}{2} F^2 + \mathcal{L}_{gf} + \mathcal{L}_{gh}$$

$$F_{\mu\nu} = \partial_\mu \mathbf{A}_\nu - \partial_\nu \mathbf{A}_\mu + i g [\mathbf{A}_\mu, \mathbf{A}_\nu]$$

## Landau gauge

- ▶ simplest one for functional equations
- ▶  $\partial_\mu \mathbf{A}_\mu = 0$ :  $\mathcal{L}_{gf} = \frac{1}{2\xi} (\partial_\mu \mathbf{A}_\mu)^2$ ,  $\xi \rightarrow 0$
- ▶ requires ghost fields:  $\mathcal{L}_{gh} = \bar{\mathbf{c}} (-\square + g \mathbf{A} \times) \mathbf{c}$



# Propagators

$$\begin{aligned} i & \text{---} \bullet \text{---} j^{-1} = + \quad i & \text{---} \text{---} j^{-1} - \frac{1}{2} \quad i & \text{---} \textcircled{1} \text{---} j^{-1} - \frac{1}{2} \quad i & \text{---} \textcircled{1} \text{---} i + \quad i & \text{---} \textcircled{1} \text{---} i \\ & \qquad \qquad \qquad - \frac{1}{6} \quad i & \text{---} \textcircled{1} \text{---} i - \frac{1}{2} \quad i & \text{---} \textcircled{1} \text{---} j \\ j & \text{---} \bullet \text{---} i^{-1} = + \quad j & \text{---} \text{---} i^{-1} - \quad j & \text{---} \textcircled{1} \text{---} j \end{aligned}$$

Models or results for vertices required.

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Models or results for vertices required.

(Tadpole vanishes perturbatively, but can contribute non-perturbatively [MQH, von Smekal '14].)

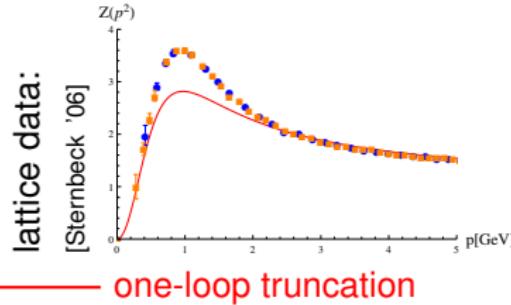
# Propagators

$$i \text{---} j^{-1} = i \text{---} j^{-1} - \frac{1}{2} i \text{---} \text{loop} + i \text{---} \text{ghost-loop}$$

$$j \text{---} i^{-1} = j \text{---} i^{-1} - j \text{---} \text{loop}$$

Typical truncation:

no four-gluon vertex, bare ghost-gluon vertex, model for three-gluon vertex



Comparison with lattice results  
→ missing strength in mid-momentum regime; attributed to

- ▶ neglected two-loop diagrams?
- ▶ vertices?

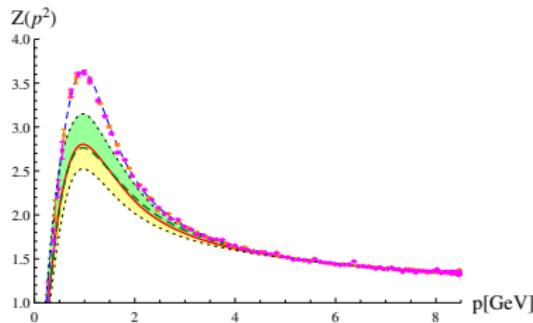
# Propagators

$$\text{Diagram 1: } i \text{ (red line)} - j \text{ (black dot)}^{-1} = + \text{ (red line)} - j \text{ (black dot)}^{-1} - \frac{1}{2} \text{ (red line)} - j \text{ (black dot)} \text{ (red loop)}^{-1} + \text{ (red line)} - j \text{ (black dot)} \text{ (green dashed loop)}^{-1}$$

$$\text{Diagram 2: } j \text{ (green dashed line)} - i \text{ (black dot)}^{-1} = + \text{ (green dashed line)} - j \text{ (black dot)}^{-1} - \text{ (green dashed line)} - j \text{ (black dot)} \text{ (red loop)}^{-1}$$

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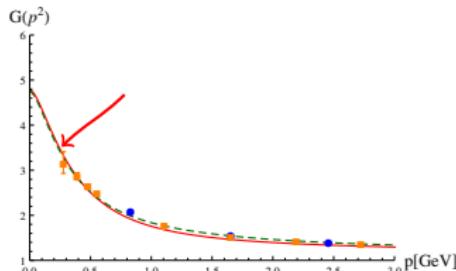
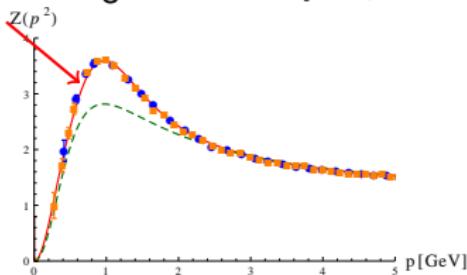


→ Using results from a three-gluon vertex calculation, importance of two-loop diagrams shown.

[Blum, MQH, Mitter, von Smekal '13]

# Propagator results

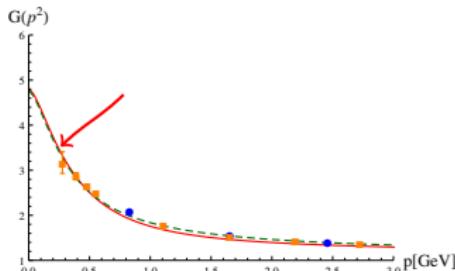
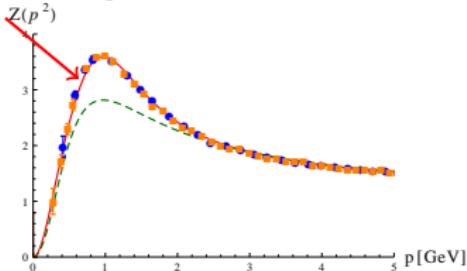
Dynamic ghost-gluon vertex, opt. eff.  
three-gluon vertex [MQH, von Smekal '13]



Good quantitative agreement for ghost *and* gluon dressings.  $\Rightarrow$  Input for further calculations.

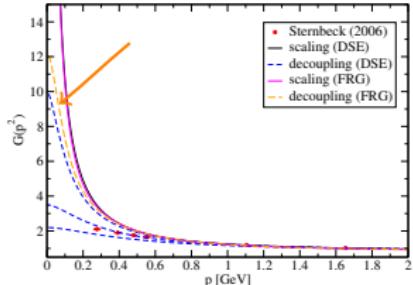
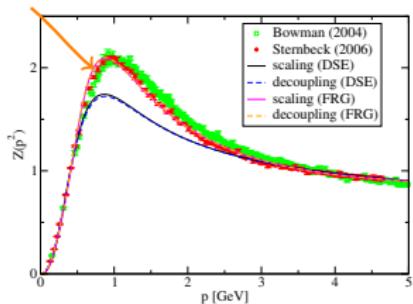
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FRG results

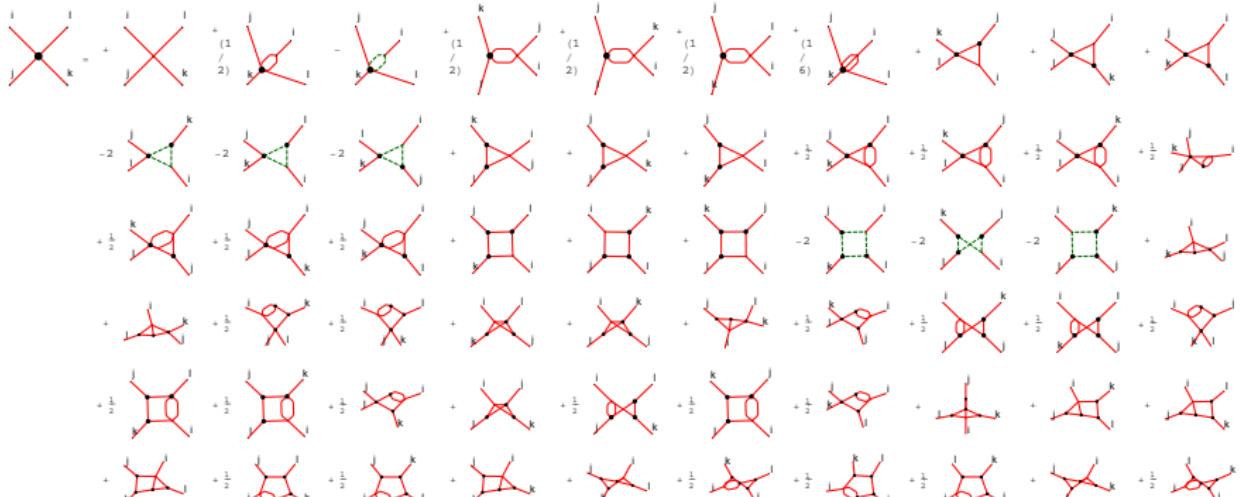
[Fischer, Maas, Pawłowski '08]



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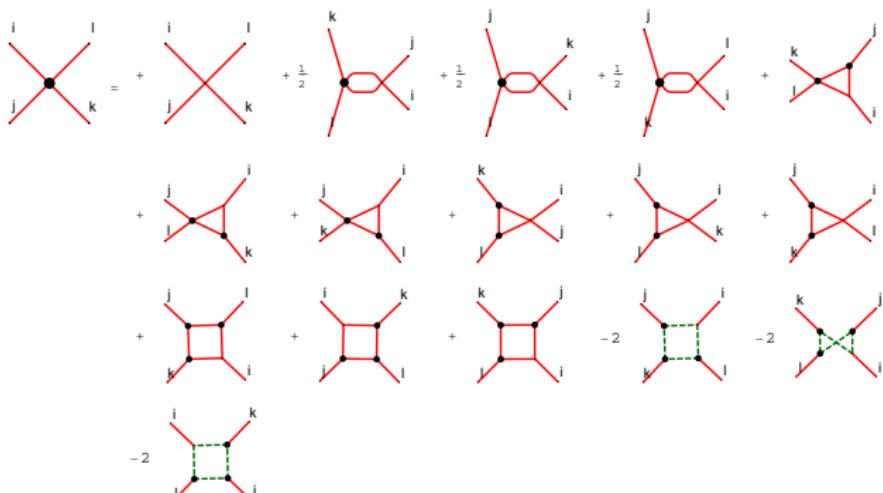
# Four-gluon vertex

- ▶ 20 one-loop, 39 two-loop diagrams



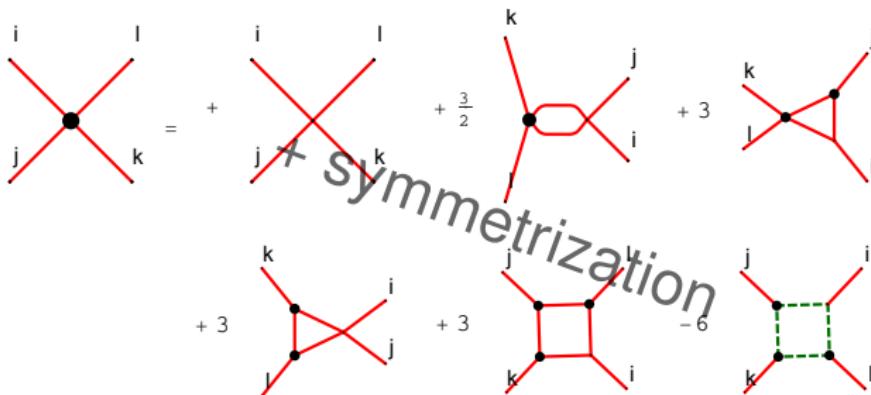
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- 16 diagrams



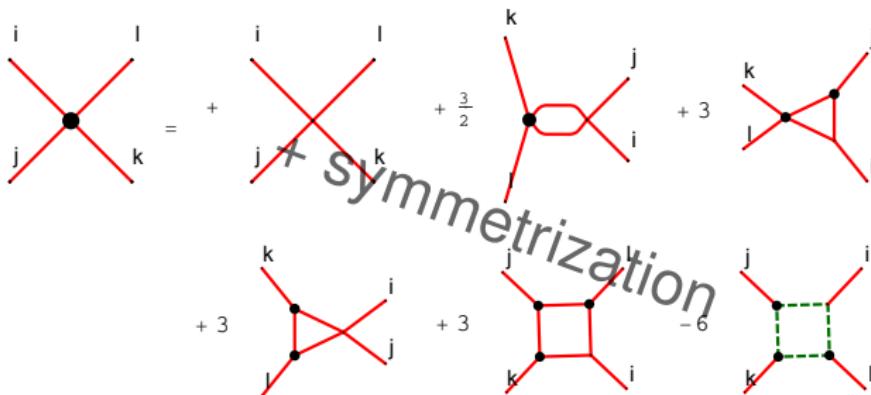
# Four-gluon vertex

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- ▶ Calculate **full** momentum dependence.
  - Access to all permutations of this diagram. → 6 diagrams



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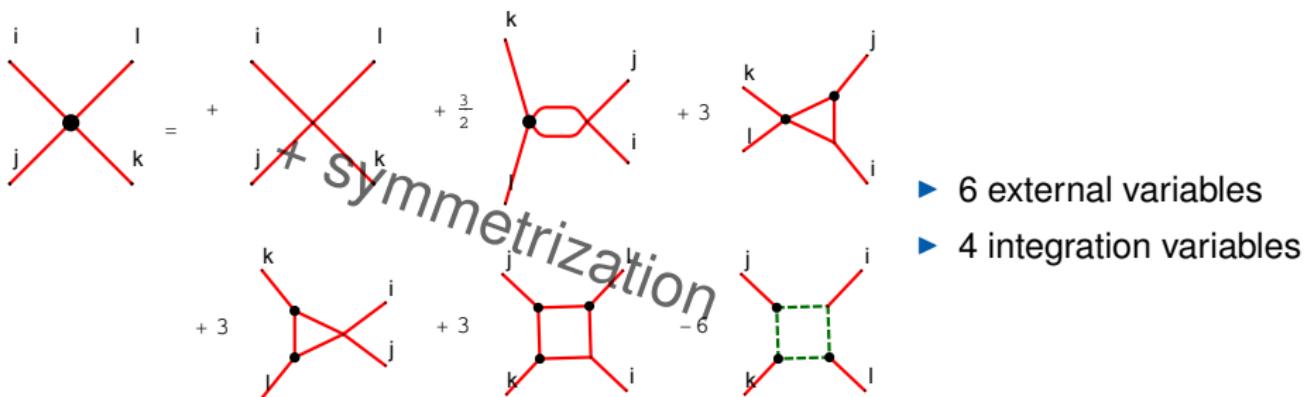
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No dependence on unknown Green functions! → 'Truncation closes.'

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Full calculation: Restriction to tree-level tensor:

$$\Gamma_{\mu\nu\rho\sigma}^{abcd}(p, q, r, s) = \Gamma_{\mu\nu\rho\sigma}^{(0),abcd} D^{4g}(p, q, r, s).$$

Non-perturbative information in  $D^{4g}(p, q, r, s)$ .

Derivation of DSE using *DoFun* [Braun, MQH '11]:

*Mathematica* package for derivation of functional equations

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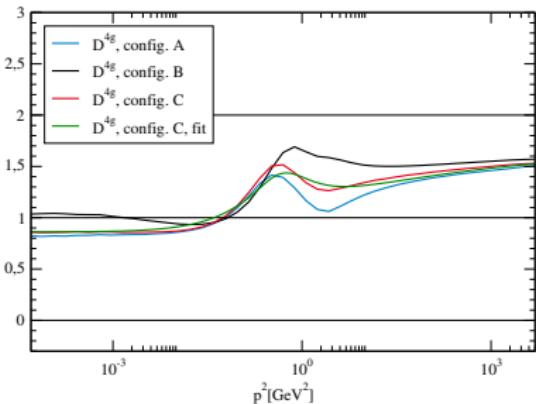
Derivation of DSE using *DoFun* [Braun, MQH '11]:

*Mathematica* package for derivation of functional equations

Other DSE results (configuration **A** only):

- ▶ Box-only truncation: [Kellermann, Fischer '08]
- ▶ Same truncation as here, semi-perturbative approximation (1 iteration):  
[Binosi, Ibáñez, Papavassiliou '14]

# Four-gluon vertex results

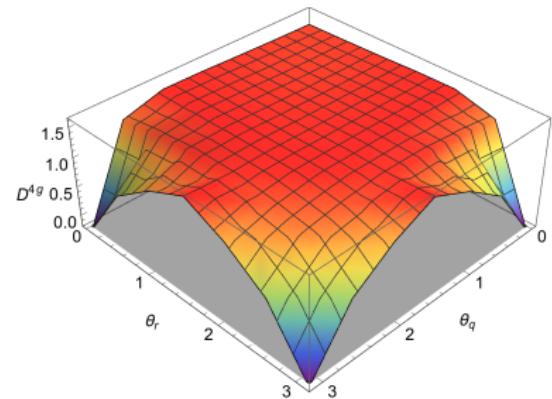
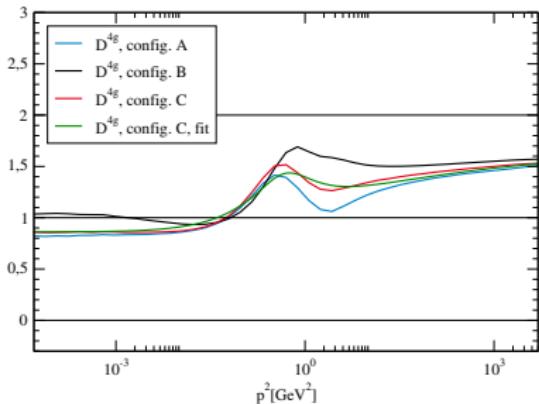


[Cyrol, MQH, von Smekal '14]

2-parameter fit:

$$D_{\text{model}}^{4g, \text{ dec}}(p, q, r, s) = (a \tanh(b/\bar{p}^2) + 1) D_{\text{RG}}^{4g}(p, q, r, s)$$

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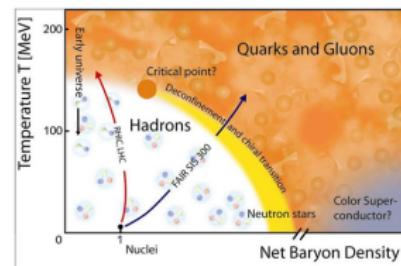
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# Non-vanishing temperature

- ▶ Phase diagram: (dual) quark condensates, Polyakov loop potential, ...
- ▶ Can be calculated from correlation functions!

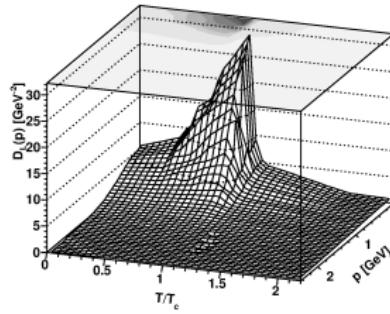
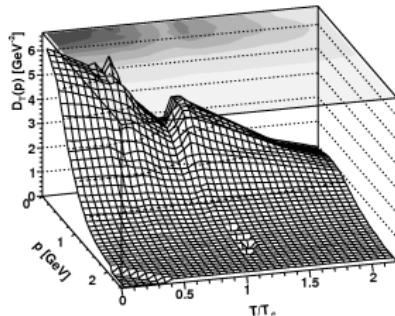
Related talks: Fister, Juricic,  
Lücker, Kamikado, Khan, Mitter,  
Müller, Pawłowski, Rechenberger,  
Rennecke, Schaefer, Springer,  
Strodthoff, Weise, Yamada, ...



- ▶ For now: Yang-Mills theory
- ▶ Goal: Calculate quantities that are difficult to obtain on the lattice.

# Propagators

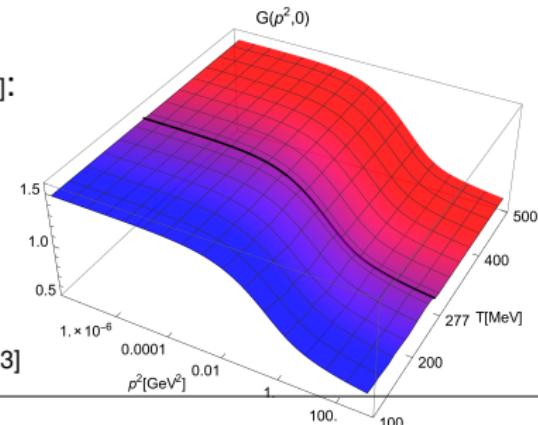
Chromomagnetic and chromoelectric gluons from the lattice [Fischer, Maas, Müller '10]:



→ Input for DSEs

Ghost dressing  $G(p^2)$  from DSE [MQH, von Smekal '13]:

$$\begin{array}{c} \text{---} \\ \bullet \end{array} -1 + \begin{array}{c} \text{---} \\ \bullet \end{array} -1 - \begin{array}{c} \text{---} \\ \bullet \end{array} - \begin{array}{c} \text{---} \\ \bullet \end{array}$$



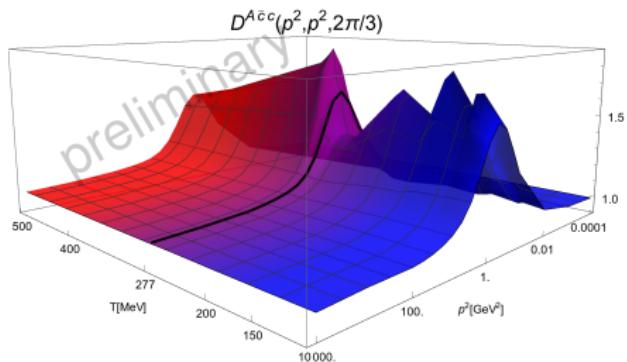
Propagators from

FRG: [Fister, Pawłowski '11]

Massive Yang-Mills: [Reinosa, Serreau, Tissier, Wschebor '13]

# Three-point warm-up: Ghost-gluon vertex

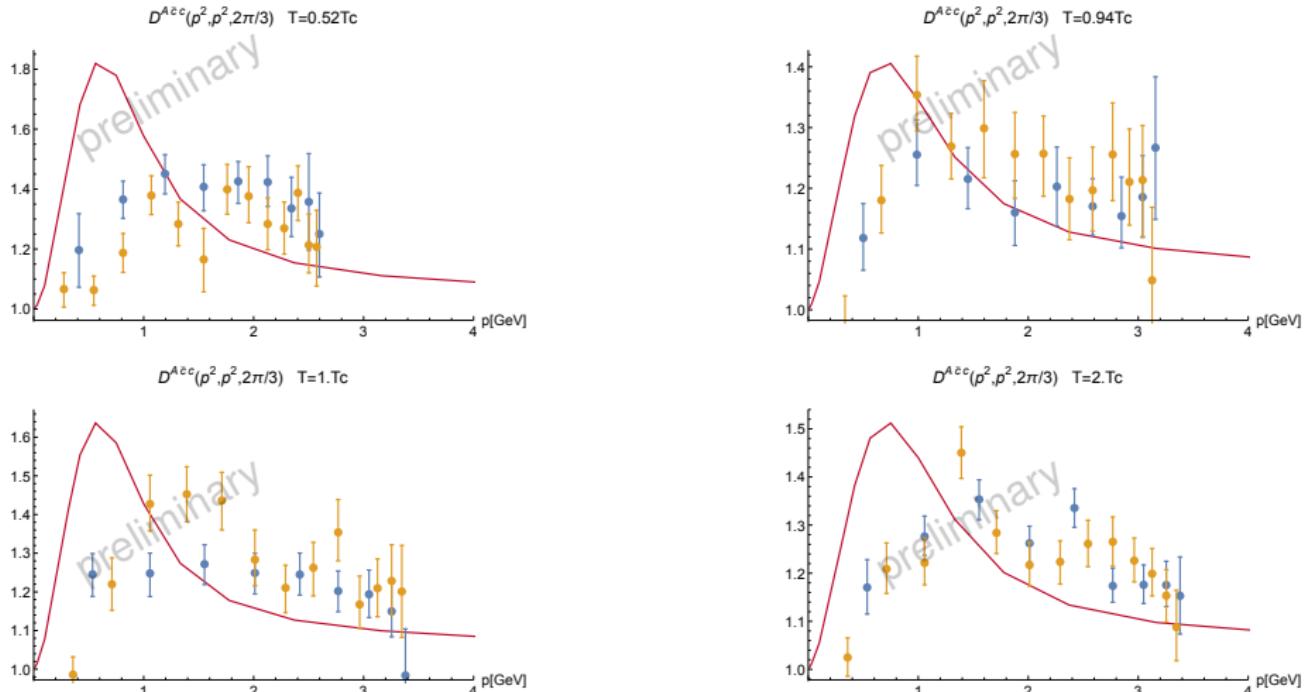
DSE calculation: self-consistent solution of truncated DSE, zeroth Matsubara frequency only



- ▶ Vertices quite expensive on lattice.
- ▶ Full momentum dependence from functional equations.

Vertex from FRG: [Fister, Pawlowski '11]

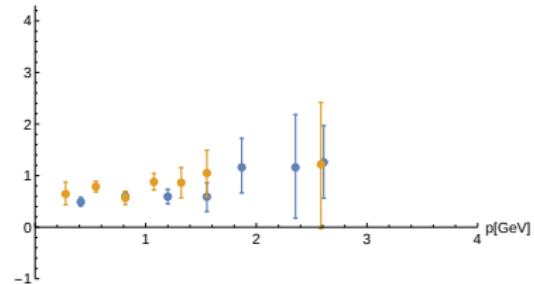
# Ghost-gluon vertex: Continuum and lattice



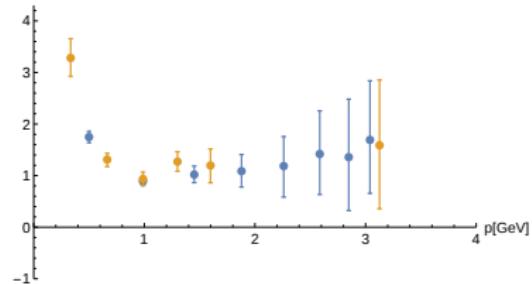
Lattice: [Fister, Maas '14]

# Three-gluon vertex: Continuum and lattice

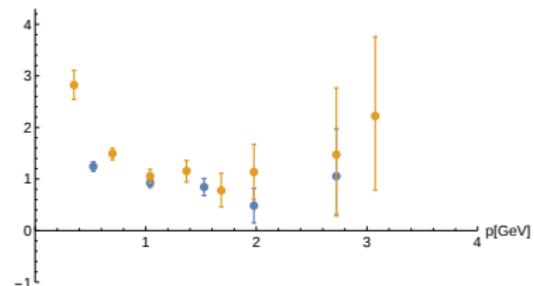
$D^{\text{AAA}}(p^2, p^2, 2\pi/3)$   $T=0.52T_c$



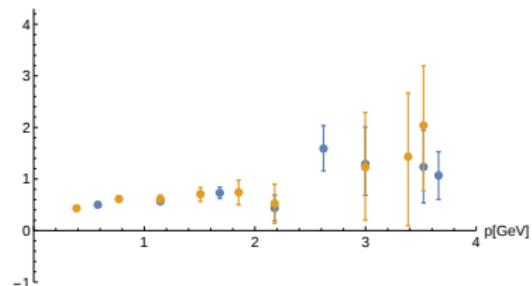
$D^{\text{AAA}}(p^2, p^2, 2\pi/3)$   $T=0.94T_c$



$D^{\text{AAA}}(p^2, p^2, 2\pi/3)$   $T=0.98T_c$



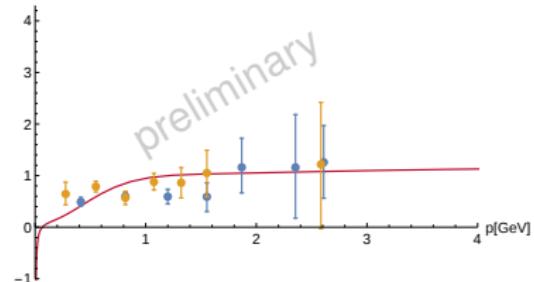
$D^{\text{AAA}}(p^2, p^2, 2\pi/3)$   $T=1.08T_c$



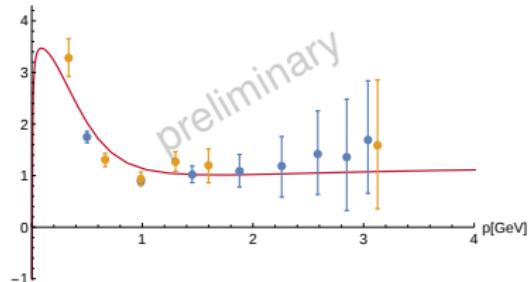
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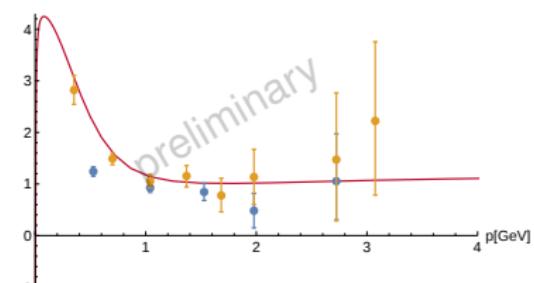
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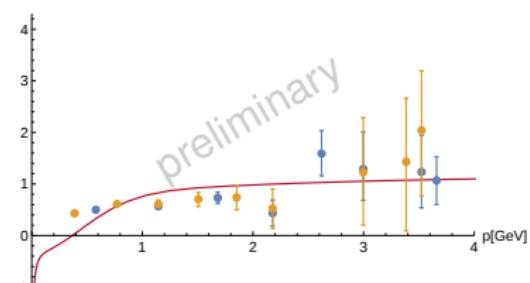
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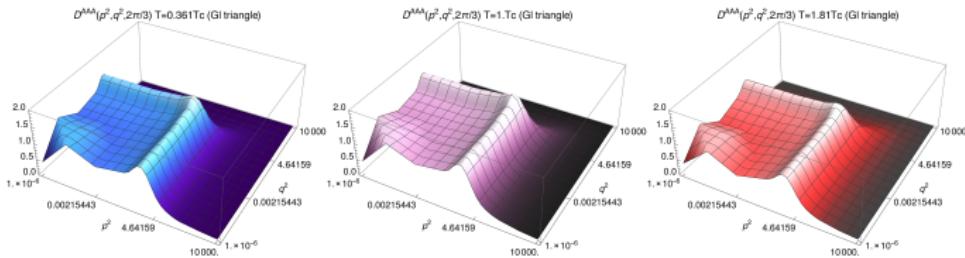


Lattice: [Fister, Maas '14]

# Three-gluon vertex: Diagram contributions

Diagram contributions ( $T = 0$ : 6,  $T > 0$ : 22 diagrams):

- ▶ Only chromomagnetic propagators: dominant



# Three-gluon vertex: Diagram contributions

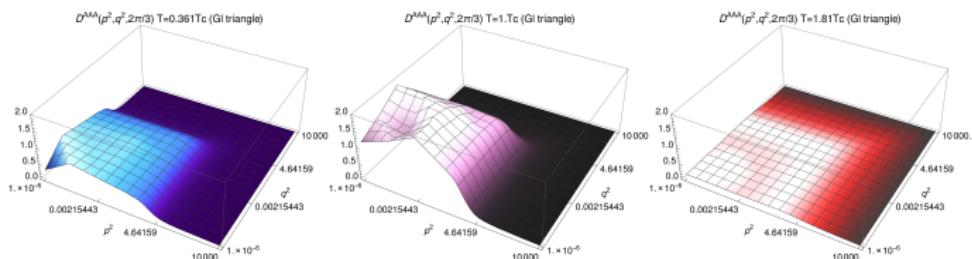
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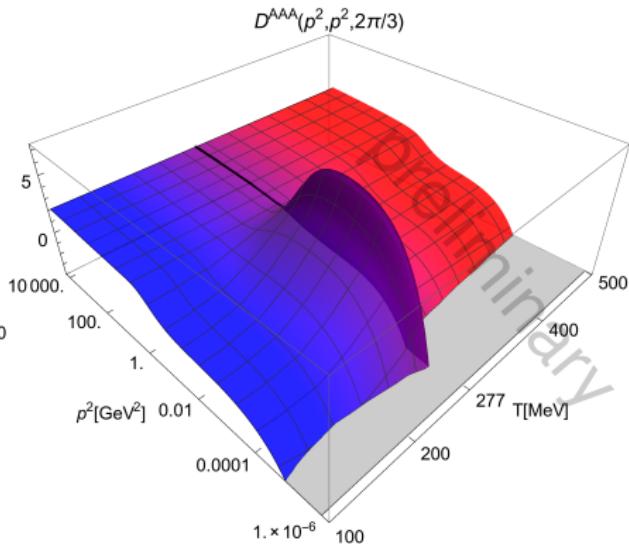
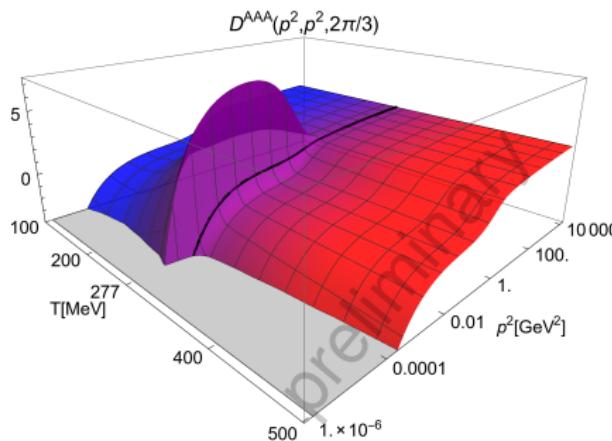
Diagram contributions ( $T = 0$ : 6,  $T > 0$ : 22 diagrams):

- ▶ Only chromomagnetic propagators: dominant
- ▶ Both propagators contained: negligible
- ▶ Only chromoelectric propagators: Importance depends on temperature, especially for  $T < T_c$ .



# Three-gluon vertex

DSE calculation: semi-perturbative approximation (first iteration only)



- ▶ Conjecture:  
Primitively divergent correlation functions sufficient for a  
**quantitative, self-consistent and self-contained description.**
- ▶ Truncation naturally 'closes' with four-gluon vertex.
- ▶ Two-point functions: Quantitative effects understood.
- ▶ Three-point functions: Good agreement with lattice.
- ▶ Four-point functions: Test four-gluon vertex in DSEs of lower n-point functions.

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Vertices at non-vanishing temperature:

- ▶ Qualitative agreement with lattice results
- ▶ Source of behavior of **three-gluon vertex  $\sim T_c$**  elucidated
- ▶ Basis for model building
- ▶ Effects of dressed vertices, e.g., in Polyakov loop potential?
- ▶ Basis for extension to QCD and  $\mu > 0$

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Thank you for your attention.

# One-momentum configurations

6 independent variables: 3 momentum squares, 3 angles

Configuration

A

B

C

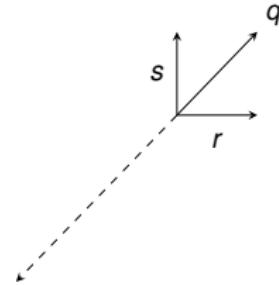
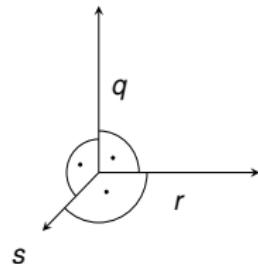
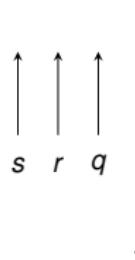
Definition

$$S^2 = R^2 = Q^2 = p^2 \\ \theta_r = \theta_q = \psi_q = 0$$

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$$S^2 = R^2 = p^2, \quad Q^2 = 2p^2 \\ \theta_r = \frac{\pi}{2}, \quad \theta_q = \frac{\pi}{4}, \quad \psi_q = 0$$

Visualization



# Other dressing functions

By transverse projection and orthonormalization construct from

$$\tilde{V}_{1,\mu\nu\rho\sigma}^{abcd} = \Gamma_{\mu\nu\rho\sigma}^{(0),abcd}$$

$$\tilde{V}_{2,\mu\nu\rho\sigma}^{abcd} = \delta^{ab}\delta^{cd}\delta_{\mu\nu}\delta_{\rho\sigma} + \delta^{ac}\delta^{bd}\delta_{\mu\rho}\delta_{\nu\sigma} + \delta^{ad}\delta^{bc}\delta_{\mu\sigma}\delta_{\nu\rho}$$

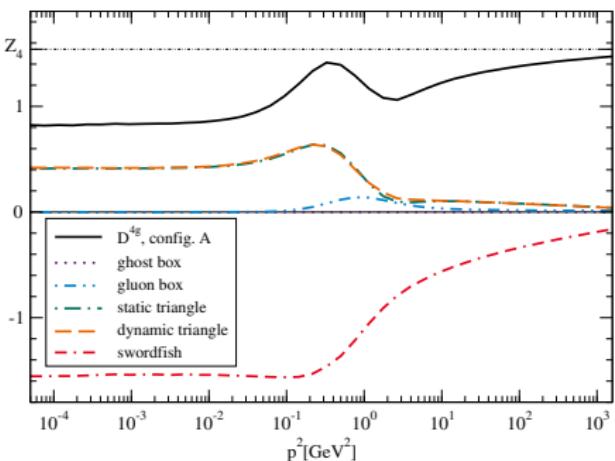
$$\tilde{V}_{3,\mu\nu\rho\sigma}^{abcd} = (\delta^{ac}\delta^{bd} + \delta^{ad}\delta^{bc})\delta_{\mu\nu}\delta_{\rho\sigma} + (\delta^{ab}\delta^{cd} + \delta^{ad}\delta^{bc})\delta_{\mu\rho}\delta_{\nu\sigma} + (\delta^{ab}\delta^{cd} + \delta^{ac}\delta^{bd})\delta_{\mu\sigma}\delta_{\nu\rho}$$

the tensors  $V_1$ ,  $V_2$ ,  $V_3$ .

Another class (4 momenta):

$$\tilde{P}_{\mu\nu\rho\sigma}^{abcd} = (\delta^{ab}\delta^{cd} + \delta^{ac}\delta^{bd} + \delta^{ad}\delta^{bc}) \frac{s_\mu r_\nu q_\rho p_\sigma + r_\mu s_\nu p_\rho q_\sigma + q_\mu p_\nu s_\rho r_\sigma}{\sqrt{p^2 q^2 r^2 p^2}}$$

# Four-gluon vertex: Individual diagrams



- ▶ Swordfish
- ▶ Dynamic triangle
- ▶ Static triangle
- ▶ Gluon box
- ▶ Ghost box

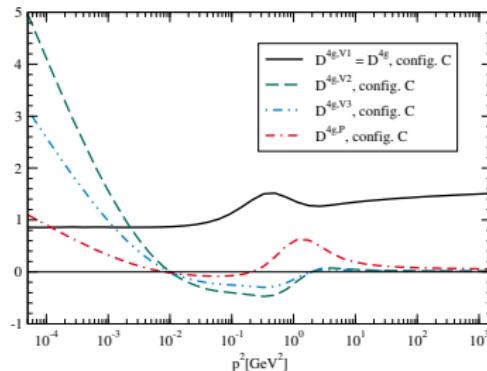


In general the non-perturbative four-gluon vertex is very close to its tree-level.

Note: Renormalization important to get the correct weights of all diagrams.

# Results for other dressing functions

- ▶  $V_1$  close to tree-level.
- ▶ Logarithmic IR divergence
- ▶ For  $V_2$  and  $V_3$  individual diagrams almost cancel each other.
- ▶  $P$  small, but richer structure
- ▶ Not calculated self-consistently



[Cyrol, MQH, von Smekal '14]

$V_1$ : tree-level

$V_2, V_3$ : only metric tensors,  $\sim \delta_{\mu\nu}\delta_{\rho\sigma}$

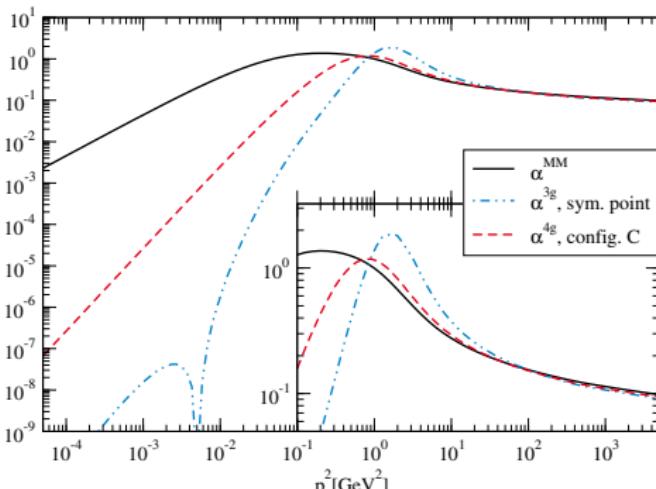
$P$ : four momenta,  $\sim s_\mu r_\nu q_\rho p_\sigma$

Note: Renormalization important to get the correct weights of all diagrams.

# Running couplings

Can be determined from

- ▶ ghost-gluon vertex  $\rightarrow \alpha^{MM}$ , cf. MiniMOM scheme  
[von Smekal, Alkofer, Hauck '97, von Smekal, Maltman, Sternbeck '09]
- ▶ all other vertices  $\rightarrow \alpha^{3g}, \alpha^{4g}, \dots$  [Alkofer, Llanes-Estrada, Fischer '05]



$$\alpha^{MM}(p^2) = \alpha(\mu^2) G(p^2)^2 Z(p^2)$$

$$\alpha^{3g}(p^2) = \alpha(\mu^2) \frac{[D^{4g}(p^2)]^2 Z^3(p^2)}{[D^{4g}(\mu^2)]^2 Z^3(\mu^2)}$$

$$\alpha^{4g}(p^2) = \alpha(\mu^2) \frac{D^{4g}(p^2) Z^2(p^2)}{D^{4g}(\mu^2) Z^2(\mu^2)}$$

[Blum, MQH, Mitter, von Smekal '14,  
Cyrol, MQH, von Smekal '14]