

Probing phase transitions between discrete and continuum quantum spacetime

Astrid Eichhorn

Perimeter Institute, Waterloo, Canada

with Tim Koslowski, University of New Brunswick, Canada

ERG 2014, Lefkada

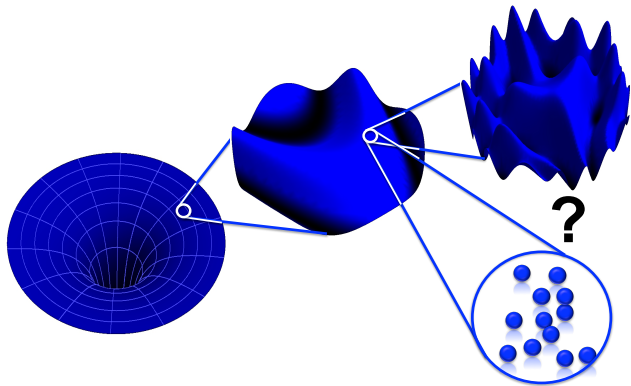


Outline

- Matrix models for 2-d quantum gravity
- Continuum limit as RG fixed point
- New FRG tools and results [with T. Koslowski, 2013/14]
- Relation to asymptotic safety?

What is the fundamental nature of spacetime?

Quantum gravity: spacetime fluctuations at the Planck scale



Goal: $\int_{\text{spacetimes}} e^{iS} \rightarrow \int_{\text{spacetimes}} e^{-S}$

Path-integral for quantum gravity

$$\int_{\text{spacetimes}} e^{iS} \rightarrow \int_{\text{spacetimes}} e^{-S}$$

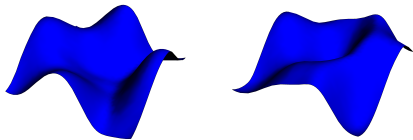
Path-integral for quantum gravity

$$\int_{\text{spacetimes}} e^{iS} \rightarrow \int_{\text{spacetimes}} e^{-S} \quad ? \text{ What fluctuates ?}$$

Path-integral for quantum gravity

$$\int_{\text{spacetimes}} e^{iS} \rightarrow \int_{\text{spacetimes}} e^{-S} \quad ? \text{ What fluctuates ?}$$

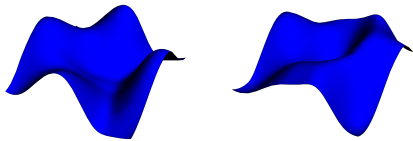
metric/curvature



Path-integral for quantum gravity

$$\int_{\text{spacetimes}} e^{iS} \rightarrow \int_{\text{spacetimes}} e^{-S} \quad ? \text{ What fluctuates ?}$$

metric/curvature

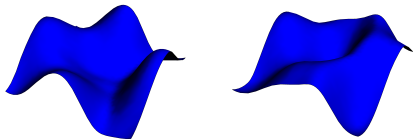


dimensionality

Path-integral for quantum gravity

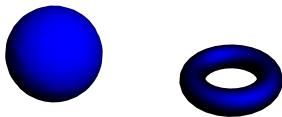
$$\int_{\text{spacetimes}} e^{iS} \rightarrow \int_{\text{spacetimes}} e^{-S} \quad ? \text{ What fluctuates ?}$$

metric/curvature



dimensionality

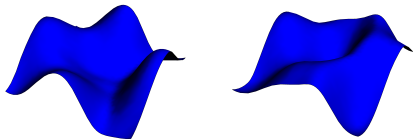
topology



Path-integral for quantum gravity

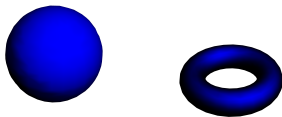
$$\int_{\text{spacetimes}} e^{iS} \rightarrow \int_{\text{spacetimes}} e^{-S} \quad ? \text{ What fluctuates ?}$$

metric/curvature



dimensionality

topology



matrix/tensor models [Weingarten, Ambjorn, Durhuus, Fröhlich, Kazakov, Migdal, Boulatov, 1980's]

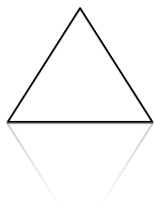
(group field theories/ tensor track [Freidel, Gurau, Oriti, Rivasseau, 2000's]):

metric & topology fluctuate

evaluate path integral by discretization

Building quantum spacetime

building blocks of spacetime:



2d

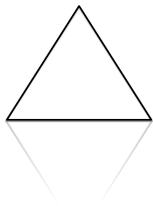


3d

...

Building quantum spacetime

building blocks of spacetime:

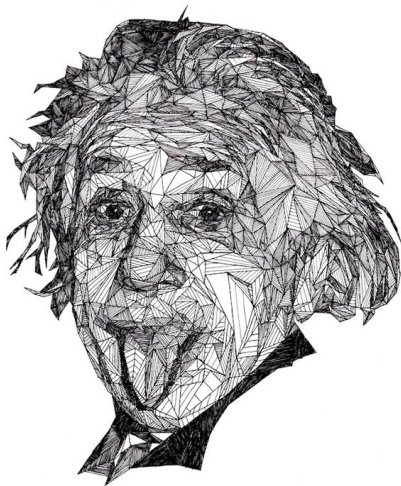


2d



3d

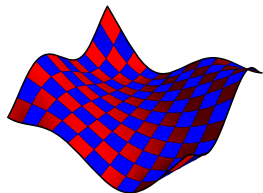
...



Building quantum spacetime

$$\sum_{\text{top.}} \int \mathcal{D}g_{\mu\nu} e^{-S[g_{\mu\nu}]} \longrightarrow \sum_{\mathcal{T}} e^{-S[\mathcal{T}]}$$

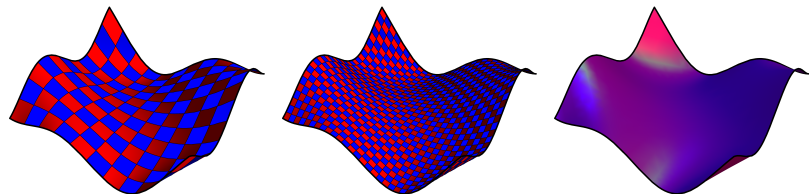
superposition of quantum spacetimes \rightarrow sum over simplicial complexes



Building quantum spacetime

$$\sum_{\text{top.}} \int \mathcal{D}g_{\mu\nu} e^{-S[g_{\mu\nu}]} \longrightarrow \sum_{\mathcal{T}} e^{-S[\mathcal{T}]}$$

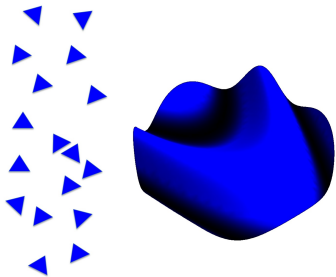
superposition of quantum spacetimes \rightarrow sum over simplicial complexes



Main challenge: What is the continuum limit?

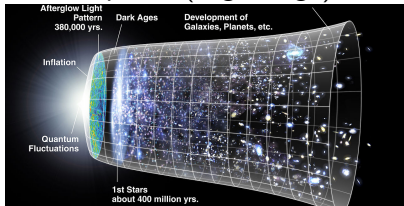
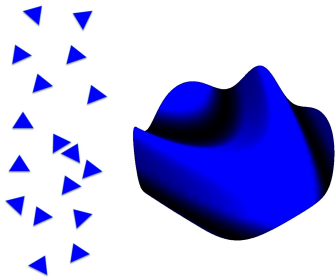
(Tensor models/ group field theories [Ben Geloun, Freidel, Gurau, Oriti, Rivasseau...])

Spacetime as a “condensate” of building blocks



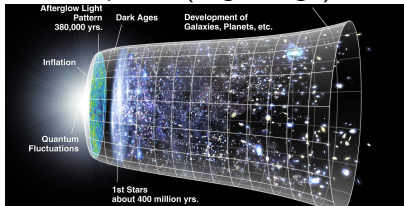
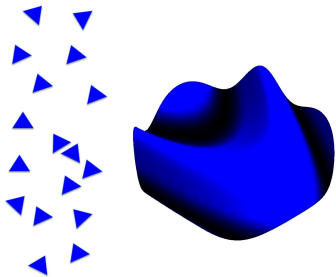
Spacetime as a “condensate” of building blocks

physical phase transition from pregeometric phase (Big Bang?)



Spacetime as a “condensate” of building blocks

physical phase transition from pregeometric phase (Big Bang?)



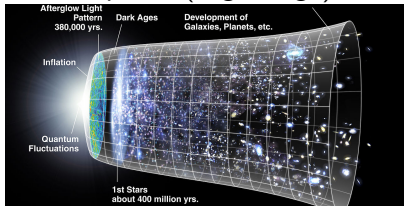
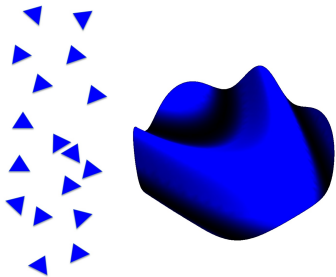
or

mathematical transition to the only physical phase



Spacetime as a “condensate” of building blocks

physical phase transition from pregeometric phase (Big Bang?)



or

mathematical transition to the only physical phase



technically: Renormalization Group fixed point

two-dimensional quantum gravity

Matrix model for two-dimensional quantum gravity

$$2 \text{ d: } \int d^2x \sqrt{g} R = 4\pi\chi$$

$$\chi = 2 - 2h:$$



$$\chi = 2$$



$$\chi = 0$$



$$\chi = -2$$



$$\chi = -4$$

Matrix model for two-dimensional quantum gravity

$$2 \text{ d: } \int d^2x \sqrt{g} R = 4\pi\chi$$

$$\chi = 2 - 2h:$$



$$\chi = 2$$



$$\chi = 0$$



$$\chi = -2$$



$$\chi = -4$$

assume Einstein-Hilbert action:

$$S = \frac{1}{16\pi G} \int d^2x \sqrt{g} (R - 2\Lambda) = \frac{1}{4G} \chi - 2\Lambda A$$

Matrix model for two-dimensional quantum gravity

$$2 \text{ d: } \int d^2x \sqrt{g} R = 4\pi\chi$$

$$\chi = 2 - 2h:$$



$$\chi = 2$$



$$\chi = 0$$



$$\chi = -2$$



$$\chi = -4$$

assume Einstein-Hilbert action:

$$S = \frac{1}{16\pi G} \int d^2x \sqrt{g} (R - 2\Lambda) = \frac{1}{4G} \chi - 2\Lambda A$$

$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi}$$

Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi}$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

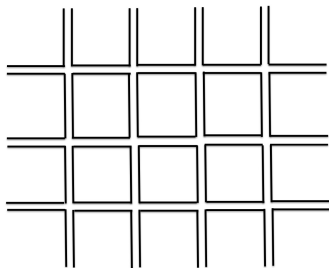
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi}$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

Feynman diagrams:



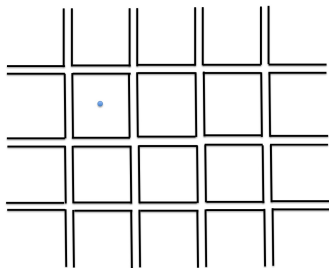
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi}$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

Feynman diagrams:



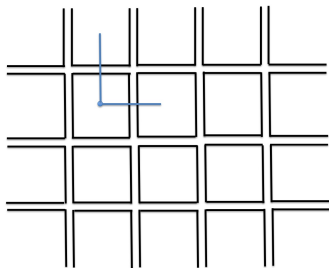
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi}$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

Feynman diagrams:



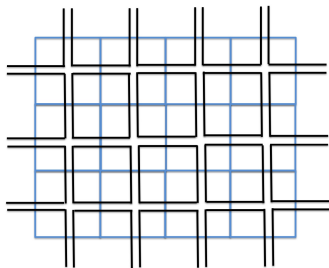
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi},$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

Feynman diagrams:
dual is “squarulation”



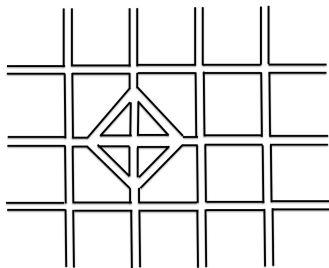
Matrix model for two-dimensional quantum gravity

$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi},$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

Feynman diagrams:
dual is "squarulation"



Matrix model for two-dimensional quantum gravity

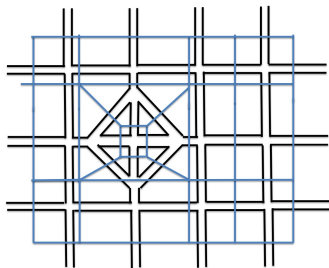
$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi},$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

Feynman diagrams:
dual is “squarulation”

Feynman diagram expansion
yields sum over all “squarulations”



Matrix model for two-dimensional quantum gravity

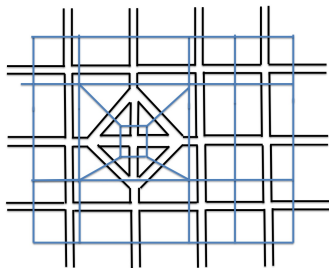
$$Z_{grav} \sim \sum_{\text{topologies}} \int \mathcal{D}g_{\mu\nu} e^{-2\Lambda A + \frac{1}{4G} \chi},$$

reformulate as hermitian matrix model: ϕ is $N \times N$ matrix

$$Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$$

Feynman diagrams:
dual is “squarulation”

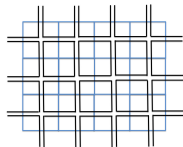
Feynman diagram expansion
yields sum over all “squarulations”



similar: d -dimensional gravity \leftrightarrow rank- d tensor model

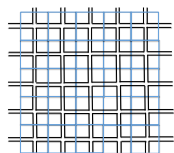
Matrix model for two-dimensional quantum gravity

continuum limit: $Z_{grav} \sim \ln \left(\int d\phi e^{-\text{tr}\phi^2 - g\text{tr}\phi^4} \right)$



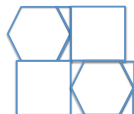
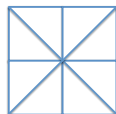
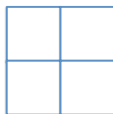
$$g \rightarrow g_c$$

$$\langle A \rangle \rightarrow \infty$$



(cf. phase transition)

→ universal behavior!



Large N limit

$$Z_{grav} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h = 0 \text{ dominates}$$

Large N limit

$$Z_{grav} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h = 0 \text{ dominates}$$

Double-scaling limit: [Brézin & Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_s/2} = \text{const.}, \quad \gamma_s = -1/2.$$

Large N limit

$$Z_{grav} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h = 0 \text{ dominates}$$

Double-scaling limit: [Brézin & Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_s/2} = \text{const.}, \quad \gamma_s = -1/2.$$

$$\rightarrow g(N) = g_c + \left(\frac{N}{c}\right)^{-2/(2-\gamma_s)}$$

Large N limit

$$Z_{grav} = \sum_h N^{2(1-h)} Z_h \quad \Rightarrow \quad N \rightarrow \infty: h = 0 \text{ dominates}$$

Double-scaling limit: [Brézin& Kazakov, Gross & Migdal, Douglas & Shenker, 1990]

$N \rightarrow \infty$ and $g \rightarrow g_c$ such that N^{-2h} is compensated

$$N(g - g_c)^{1-\gamma_s/2} = \text{const.}, \quad \gamma_s = -1/2.$$

$$\rightarrow g(N) = g_c + \left(\frac{N}{c}\right)^{-2/(2-\gamma_s)}$$



\rightarrow looks like RG fixed-point behavior: $g(k) = g_* + c \left(\frac{k}{k_0}\right)^{-\theta}$

Double-scaling (continuum) limit \leftrightarrow Renormalization Group fixed point

$$g(N) = g_c + \left(\frac{N}{c}\right)^{-2/(2-\gamma_s)} \leftrightarrow g(k) = g_* + c \left(\frac{k}{k_0}\right)^{-\theta}$$

$$g_c \leftrightarrow g_*$$

$$N \leftrightarrow k$$

$$\frac{2}{2-\gamma_s} \leftrightarrow \theta$$

Double-scaling (continuum) limit \leftrightarrow Renormalization Group fixed point

$$g(N) = g_c + \left(\frac{N}{c}\right)^{-2/(2-\gamma_s)} \leftrightarrow g(k) = g_* + c \left(\frac{k}{k_0}\right)^{-\theta}$$

$$g_c \leftrightarrow g_*$$

$$N \leftrightarrow k$$

$$\frac{2}{2-\gamma_s} \leftrightarrow \theta$$

perturbative approach [Brézin, Zinn-Justin, 1992]

FRG approach [A.E., Tim Koslowski, 2013]

Double-scaling (continuum) limit \leftrightarrow Renormalization Group fixed point

$$g(N) = g_c + \left(\frac{N}{c}\right)^{-2/(2-\gamma_s)} \leftrightarrow g(k) = g_* + c \left(\frac{k}{k_0}\right)^{-\theta}$$

$$g_c \leftrightarrow g_*$$

$$N \leftrightarrow k$$

$$\frac{2}{2-\gamma_s} \leftrightarrow \theta$$

perturbative approach [Brézin, Zinn-Justin, 1992]

FRG approach [A.E., Tim Koslowski, 2013]

main advantage:
extension to higher-dimensional models conceptually clear
(technically: larger theory space [Rivasseau, 2014])

Renormalization group scale

basic RG requirement: distinguish “low-momentum” from “high-momentum” quantum fluctuations

quantum gravity: metric fluctuates

discrete approaches: based on “pre-geometric” building blocks
⇒ No notion of momentum

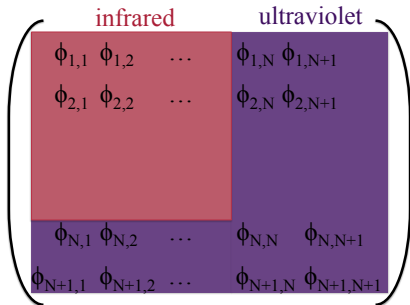
Renormalization group scale

basic RG requirement: distinguish “low-momentum” from “high-momentum” quantum fluctuations

quantum gravity: metric fluctuates

discrete approaches: based on “pre-geometric” building blocks
⇒ No notion of momentum

use N to set a scale [Zinn-Justin, Brezin, 1993; Oriti, Rivasseau, Gurau..., since 2010]



$N \rightarrow \infty$ no degrees of freedom
integrated out (ultraviolet)

$N \rightarrow 0$ all degrees of freedom
integrated out (infrared)

FRG for matrix models

regulator $\phi_{a,b} R_N(a,b)_{abcd} \phi_{cd}$ with $R_N(a,b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right)$

inspired by "Litim-cutoff"

FRG for matrix models

regulator $\phi_{a,b} R_N(a,b)_{abcd} \phi_{cd}$ with $R_N(a,b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right)$

inspired by "Litim-cutoff"

effective action Γ_N from $Z_N = \int d\phi e^{-S[\phi] - \frac{1}{2} \text{tr} \phi R_N \phi}$

FRG for matrix models

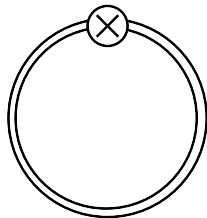
regulator $\phi_{a,b} R_N(a,b)_{abcd} \phi_{cd}$ with $R_N(a,b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right)$

inspired by "Litim-cutoff"

effective action Γ_N from $Z_N = \int d\phi e^{-S[\phi] - \frac{1}{2} \text{tr} \phi R_N \phi}$

Wetterich-equation adapted for discrete case:

$$N \partial_N \Gamma_N = \frac{1}{2} \text{tr} \left(\Gamma_N^2 + R_N \right)^{-1} N \partial_N R_N$$



FRG for matrix models

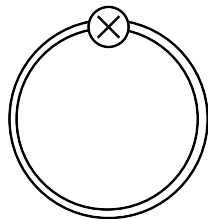
regulator $\phi_{a,b} R_N(a,b)_{abcd} \phi_{cd}$ with $R_N(a,b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right)$

inspired by "Litim-cutoff"

effective action Γ_N from $Z_N = \int d\phi e^{-S[\phi] - \frac{1}{2} \text{tr} \phi R_N \phi}$

Wetterich-equation adapted for discrete case:

$$N \partial_N \Gamma_N = \frac{1}{2} \text{tr} \left(\Gamma_N^2 + R_N \right)^{-1} N \partial_N R_N$$



theory space: $\Gamma_N = \sum_i \frac{g_{2i}}{2} \text{tr} \phi^{2i} + \sum_{i,j; i+j \text{ even}} \frac{g_{i,j}}{2} \text{tr} \phi^i \text{tr} \phi^j + \dots$

assignment of canonical dimensionality: $\bar{g}_i = g_i N^{-(i-2)/2}$

SU(N) symmetry breaking and tadpole approximation

RG approach: cutoff at $N \Rightarrow$ symmetry-breaking

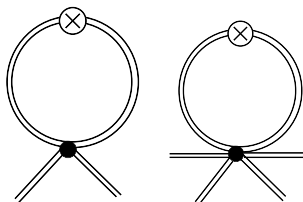
cutoff function: $R_N(a, b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right) \Rightarrow$ completely broken symmetry $\Rightarrow \phi_{ab} f(a, b) \phi_{ba}$ etc are generated

SU(N) symmetry breaking and tadpole approximation

RG approach: cutoff at $N \Rightarrow$ symmetry-breaking

cutoff function: $R_N(a, b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right) \Rightarrow$ completely broken symmetry $\Rightarrow \phi_{ab} f(a, b) \phi_{ba}$ etc are generated

approximation: restrict to tadpole diagrams

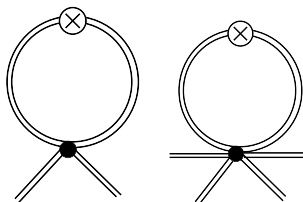


SU(N) symmetry breaking and tadpole approximation

RG approach: cutoff at $N \Rightarrow$ symmetry-breaking

cutoff function: $R_N(a, b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right) \Rightarrow$ completely broken symmetry $\Rightarrow \phi_{ab} f(a, b) \phi_{ba}$ etc are generated

approximation: restrict to tadpole diagrams



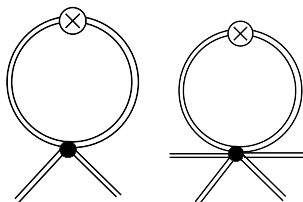
✓ $SU(N)$ symmetry unbroken

SU(N) symmetry breaking and tadpole approximation

RG approach: cutoff at $N \Rightarrow$ symmetry-breaking

cutoff function: $R_N(a, b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right) \Rightarrow$ completely broken symmetry $\Rightarrow \phi_{ab} f(a, b) \phi_{ba}$ etc are generated

approximation: restrict to tadpole diagrams



✓ $SU(N)$ symmetry unbroken

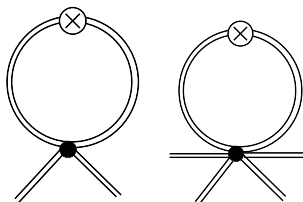
✓ computational simplicity (important at multi-trace level)

SU(N) symmetry breaking and tadpole approximation

RG approach: cutoff at $N \Rightarrow$ symmetry-breaking

cutoff function: $R_N(a, b) = \left(\frac{2N}{a+b} - 1 \right) \theta \left(1 - \frac{a+b}{2N} \right) \Rightarrow$ completely broken symmetry $\Rightarrow \phi_{ab} f(a, b) \phi_{ba}$ etc are generated

approximation: restrict to tadpole diagrams



- ✓ $SU(N)$ symmetry unbroken
- ✓ computational simplicity (important at multi-trace level)
- ✓ consistent approximation if fixed-point couplings small

Results: Double-scaling limit

search for: fixed point at $g_4 < 0$ with $\theta_1 = 0.8$ and $\theta_n < 0$ ($n \geq 2$)

Results: Double-scaling limit

search for: fixed point at $g_4 < 0$ with $\theta_1 = 0.8$ and $\theta_n < 0$ ($n \geq 2$)

for $\dot{R}\mathcal{P}^{-2} = x$

Results: Double-scaling limit

search for: fixed point at $g_4 < 0$ with $\theta_1 = 0.8$ and $\theta_n < 0$ ($n \geq 2$)

for $\dot{R}\mathcal{P}^{-2} = x$

$$\eta = -\partial_t \ln Z = 2g_4 x$$

$$\beta_{g_{2n}} = ((n-1) + \eta n) g_{2n} - 2n x g_{2(n+1)} \quad [\text{A.E., T. Koslowski, 2013, 2014}]$$

Results: Double-scaling limit

search for: fixed point at $g_4 < 0$ with $\theta_1 = 0.8$ and $\theta_n < 0$ ($n \geq 2$)

for $\dot{R}\mathcal{P}^{-2} = x$

$$\eta = -\partial_t \ln Z = 2g_4 x$$

$$\beta_{g_{2n}} = ((n-1) + \eta n) g_{2n} - 2n x g_{2(n+1)} \quad [\text{A.E., T. Koslowski, 2013, 2014}]$$

admits fixed point at $g_4 = -\frac{1}{4x}$ and $g_{2n=0}$ ($n > 2$)

Results: Double-scaling limit

search for: fixed point at $g_4 < 0$ with $\theta_1 = 0.8$ and $\theta_n < 0$ ($n \geq 2$)

for $\dot{R}\mathcal{P}^{-2} = x$

$$\eta = -\partial_t \ln Z = 2g_4 x$$

$$\beta_{g_{2n}} = ((n-1) + \eta n) g_{2n} - 2n x g_{2(n+1)} \quad [\text{A.E., T. Koslowski, 2013, 2014}]$$

admits fixed point at $g_4 = -\frac{1}{4x}$ and $g_{2n=0}$ ($n > 2$)

$$\theta_1 = 1, \theta_n < 0, n \geq 2$$

Results: Double-scaling limit

search for: fixed point at $g_4 < 0$ with $\theta_1 = 0.8$ and $\theta_n < 0$ ($n \geq 2$)

for $\dot{R}\mathcal{P}^{-2} = x$

$$\eta = -\partial_t \ln Z = 2g_4 x$$

$$\beta_{g_{2n}} = ((n-1) + \eta n) g_{2n} - 2n x g_{2(n+1)} \quad [\text{A.E., T. Koslowski, 2013, 2014}]$$

admits fixed point at $g_4 = -\frac{1}{4x}$ and $g_{2n}=0$ ($n > 2$) (**nonuniversal**)

$\theta_1 = 1$, $\theta_n < 0$, $n \geq 2$ (**universal**)

Results: Double-scaling limit

search for: fixed point at $g_4 < 0$ with $\theta_1 = 0.8$ and $\theta_n < 0$ ($n \geq 2$)

for $\dot{R}\mathcal{P}^{-2} = x$

$$\eta = -\partial_t \ln Z = 2g_4 x$$

$$\beta_{g_{2n}} = ((n-1) + \eta n) g_{2n} - 2n x g_{2(n+1)} \quad [\text{A.E., T. Koslowski, 2013, 2014}]$$

admits fixed point at $g_4 = -\frac{1}{4x}$ and $g_{2n=0}$ ($n > 2$) (**nonuniversal**)

$\theta_1 = 1$, $\theta_n < 0$, $n \geq 2$ (**universal**)

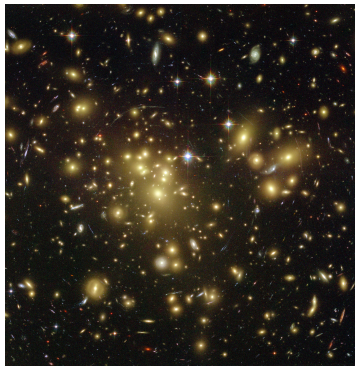
Double-scaling limit recovered with Functional RG (multi-trace truncation confirms result)

Multicritical points

Experimental fact:

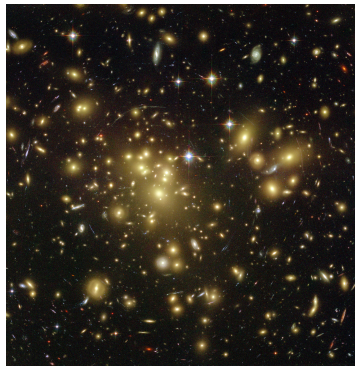
Multicritical points

Experimental fact: Universe contains gravity & matter



Multicritical points

Experimental fact: Universe contains gravity & matter



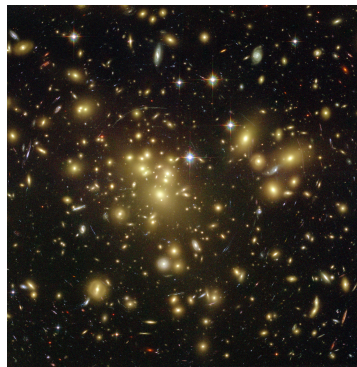
- matter and gravitational d.o.f. exist fundamentally

or

- both “emerge” from discrete pregeometric building blocks

Multicritical points

Experimental fact: Universe contains gravity & matter



- matter and gravitational d.o.f. exist fundamentally

or

- both “emerge” from discrete pregeometric building blocks

Here: matter \leftrightarrow collective excitations of matrix model

Multicritical points

fixed point $V[\phi] = \frac{1}{2}\text{tr}\phi^2 + g_4\text{tr}\phi^4$ with $g_4 < 0 \rightarrow$ 2d gravity

fixed points $V[\phi] = \frac{1}{2}\text{tr}\phi^2 + g_4\text{tr}\phi^4 + g_6\text{tr}\phi^6 + \dots$ with $g_4 < 0, g_6 > 0, \dots$



$g_4 < 0$

pure-gravity case



$g_4 < 0$



$g_6 > 0$

multicritical case



$g_4 < 0$



$g_6 > 0$



$g_8 < 0$

Multicritical points

fixed point $V[\phi] = \frac{1}{2}\text{tr}\phi^2 + g_4\text{tr}\phi^4$ with $g_4 < 0 \rightarrow$ 2d gravity

fixed points $V[\phi] = \frac{1}{2}\text{tr}\phi^2 + g_4\text{tr}\phi^4 + g_6\text{tr}\phi^6 + \dots$ with $g_4 < 0, g_6 > 0, \dots$



$g_4 < 0$

pure-gravity case



$g_4 < 0$



$g_6 > 0$

multicritical case



$g_4 < 0$



$g_6 > 0$



$g_8 < 0$

some geometric d.o.f. contribute with negative weight \Rightarrow extra d.o.f.?

\rightarrow tower of multicritical points \leftrightarrow 2d gravity + conformal matter [Kazakov,

1989, Staudacher, 1990]

Multicritical points

fixed point $V[\phi] = \frac{1}{2}\text{tr}\phi^2 + g_4\text{tr}\phi^4$ with $g_4 < 0 \rightarrow$ 2d gravity

fixed points $V[\phi] = \frac{1}{2}\text{tr}\phi^2 + g_4\text{tr}\phi^4 + g_6\text{tr}\phi^6 + \dots$ with $g_4 < 0, g_6 > 0, \dots$



$g_4 < 0$

pure-gravity case



$g_4 < 0$



$g_6 > 0$

multicritical case



$g_4 < 0$



$g_6 > 0$



$g_8 < 0$

some geometric d.o.f. contribute with negative weight \Rightarrow extra d.o.f.?

\rightarrow tower of multicritical points \leftrightarrow 2d gravity + conformal matter [Kazakov,

1989, Staudacher, 1990]

\rightarrow gravity critical exponent: $\gamma_s = 3/2 - m$ [$\theta = 2/(m + 1/2)$]

for $m = 2$ (pure gravity), $m > 2$ (gravity + conformal matter)

Results: Multicritical points

analytic results for fixed points: [A.E., T. Koslowski, 2014]

g_4	g_6	g_8	g_{10}	g_{12}	g_{14}	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6
0	0	0	0	0	0	-1	-2	-3	-4	-5	-6
$-\frac{1}{4x}$	0	0	0	0	0	1	$-\frac{1}{2}$	-1	$-\frac{3}{2}$	-2	$-\frac{5}{2}$
$-\frac{1}{3x}$	$\frac{1}{36x^2}$	0	0	0	0	1	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{2}{3}$	-1	$-\frac{4}{3}$
$-\frac{3}{8x}$	$\frac{64x^2}{3}$	$-\frac{1}{512x^3}$	0	0	0	1	$\frac{3}{4}$	$\frac{1}{2}$	$-\frac{1}{4}$	$-\frac{1}{2}$	$-\frac{3}{4}$
$-\frac{2}{5x}$	$\frac{3}{50x^2}$	$-\frac{1}{250x^3}$	$\frac{1}{10^4x^4}$	0	0	1	$\frac{4}{5}$	$\frac{3}{5}$	$\frac{2}{5}$	$-\frac{1}{5}$	$-\frac{2}{5}$
$-\frac{5}{12x}$	$\frac{5}{72x^2}$	$-\frac{5}{864x^3}$	$\frac{5}{20736x^4}$	$-\frac{1}{248832x^5}$	0	1	$\frac{5}{6}$	$\frac{2}{3}$	$\frac{1}{2}$	$\frac{1}{3}$	$-\frac{1}{6}$
$-\frac{3}{7x}$	$\frac{15}{196x^2}$	$-\frac{5}{686x^3}$	$\frac{15}{38416x^4}$	$-\frac{3}{268912x^5}$	$\frac{1}{7529537x^6}$	1	$\frac{6}{7}$	$\frac{5}{7}$	$\frac{4}{7}$	$\frac{3}{7}$	$\frac{2}{7}$

Results: Multicritical points

analytic results for fixed points: [A.E., T. Koslowski, 2014]

g_4	g_6	g_8	g_{10}	g_{12}	g_{14}	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6
0	0	0	0	0	0	-1	-2	-3	-4	-5	-6
$-\frac{1}{4x}$	0	0	0	0	0	1	$-\frac{1}{2}$	-1	$-\frac{3}{2}$	-2	$-\frac{5}{2}$
$-\frac{1}{3x}$	$\frac{1}{36x^2}$	0	0	0	0	1	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{2}{3}$	-1	$-\frac{4}{3}$
$-\frac{3}{8x}$	$\frac{64x^2}{3}$	$-\frac{1}{512x^3}$	0	0	0	1	$\frac{3}{4}$	$\frac{1}{2}$	$-\frac{1}{4}$	$-\frac{1}{2}$	$-\frac{3}{4}$
$-\frac{2}{5x}$	$\frac{3}{50x^2}$	$-\frac{1}{250x^3}$	$\frac{1}{10^4x^4}$	0	0	1	$\frac{4}{5}$	$\frac{3}{5}$	$\frac{2}{5}$	$-\frac{1}{5}$	$-\frac{2}{5}$
$-\frac{5}{12x}$	$\frac{5}{72x^2}$	$-\frac{864x^3}{5}$	$\frac{5}{20736x^4}$	$-\frac{1}{248832x^5}$	0	1	$\frac{5}{6}$	$\frac{2}{3}$	$\frac{1}{2}$	$\frac{1}{3}$	$-\frac{1}{6}$
$-\frac{3}{7x}$	$\frac{15}{196x^2}$	$-\frac{5}{686x^3}$	$\frac{15}{38416x^4}$	$-\frac{3}{268912x^5}$	$\frac{1}{7529537x^6}$	1	$\frac{6}{7}$	$\frac{5}{7}$	$\frac{4}{7}$	$\frac{3}{7}$	$\frac{2}{7}$

Gaußian fixed point

Results: Multicritical points

analytic results for fixed points: [A.E., T. Koslowski, 2014]

g_4	g_6	g_8	g_{10}	g_{12}	g_{14}	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6
0	0	0	0	0	0	-1	-2	-3	-4	-5	-6
$-\frac{1}{4x}$	0	0	0	0	0	1	$-\frac{1}{2}$	-1	$-\frac{3}{2}$	-2	$-\frac{5}{2}$
$-\frac{1}{3x}$	$\frac{1}{36x^2}$	0	0	0	0	1	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{2}{3}$	-1	$-\frac{4}{3}$
$-\frac{3}{8x}$	$\frac{64x^2}{3}$	$-\frac{1}{512x^3}$	0	0	0	1	$\frac{3}{4}$	$\frac{1}{2}$	$-\frac{1}{4}$	$-\frac{1}{2}$	$-\frac{3}{4}$
$-\frac{2}{5x}$	$\frac{3}{50x^2}$	$-\frac{1}{250x^3}$	$\frac{1}{10^4x^4}$	0	0	1	$\frac{4}{5}$	$\frac{3}{5}$	$\frac{2}{5}$	$-\frac{1}{5}$	$-\frac{2}{5}$
$-\frac{5}{12x}$	$\frac{5}{72x^2}$	$-\frac{864x^3}{5}$	$\frac{20736x^4}{15}$	$-\frac{1}{248832x^5}$	0	1	$\frac{5}{6}$	$\frac{2}{3}$	$\frac{1}{2}$	$\frac{1}{3}$	$-\frac{1}{6}$
$-\frac{3}{7x}$	$\frac{15}{196x^2}$	$-\frac{5}{686x^3}$	$\frac{38416x^4}{15}$	$-\frac{1}{268912x^5}$	$\frac{1}{7529537x^6}$	1	$\frac{6}{7}$	$\frac{5}{7}$	$\frac{4}{7}$	$\frac{3}{7}$	$\frac{2}{7}$

Gaussian fixed point

pure gravity fixed point

Results: Multicritical points

analytic results for fixed points: [A.E., T. Koslowski, 2014]

g_4	g_6	g_8	g_{10}	g_{12}	g_{14}	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6
0	0	0	0	0	0	-1	-2	-3	-4	-5	-6
$-\frac{1}{4x}$	0	0	0	0	0	1	$-\frac{1}{2}$	-1	$-\frac{3}{2}$	-2	$-\frac{5}{2}$
$-\frac{1}{3x}$	$\frac{1}{36x^2}$	0	0	0	0	1	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{2}{3}$	-1	$-\frac{4}{3}$
$-\frac{3}{8x}$	$\frac{64x^2}{3}$	$-\frac{1}{512x^3}$	0	0	0	1	$\frac{3}{4}$	$\frac{1}{2}$	$-\frac{1}{4}$	$-\frac{1}{2}$	$-\frac{3}{4}$
$-\frac{2}{5x}$	$\frac{3}{50x^2}$	$-\frac{1}{250x^3}$	$\frac{1}{10^4x^4}$	0	0	1	$\frac{4}{5}$	$\frac{3}{5}$	$\frac{2}{5}$	$-\frac{1}{5}$	$-\frac{2}{5}$
$-\frac{5}{12x}$	$\frac{5}{72x^2}$	$-\frac{864x^3}{5}$	$\frac{20736x^4}{15}$	$-\frac{1}{248832x^5}$	0	1	$\frac{5}{6}$	$\frac{2}{3}$	$\frac{1}{2}$	$\frac{1}{3}$	$-\frac{1}{6}$
$-\frac{3}{7x}$	$\frac{15}{196x^2}$	$-\frac{5}{686x^3}$	$\frac{15}{38416x^4}$	$-\frac{3}{268912x^5}$	$\frac{1}{7529537x^6}$	1	$\frac{6}{7}$	$\frac{5}{7}$	$\frac{4}{7}$	$\frac{3}{7}$	$\frac{2}{7}$

Gaußian fixed point

pure gravity fixed point

multicritical points

Results: Multicritical points

analytic results for fixed points: [A.E., T. Koslowski, 2014]

g_4	g_6	g_8	g_{10}	g_{12}	g_{14}	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6
0	0	0	0	0	0	-1	-2	-3	-4	-5	-6
$-\frac{1}{4x}$	0	0	0	0	0	1	$-\frac{1}{2}$	-1	$-\frac{3}{2}$	-2	$-\frac{5}{2}$
$-\frac{1}{3x}$	$\frac{1}{36x^2}$	0	0	0	0	1	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{2}{3}$	-1	$-\frac{4}{3}$
$-\frac{3}{8x}$	$\frac{64x^2}{3}$	$-\frac{1}{512x^3}$	0	0	0	1	$\frac{3}{4}$	$\frac{1}{2}$	$-\frac{1}{4}$	$-\frac{1}{2}$	$-\frac{3}{4}$
$-\frac{2}{5x}$	$\frac{3}{50x^2}$	$-\frac{1}{250x^3}$	$\frac{1}{10^4x^4}$	0	0	1	$\frac{4}{5}$	$\frac{3}{5}$	$\frac{2}{5}$	$-\frac{1}{5}$	$-\frac{2}{5}$
$-\frac{5}{12x}$	$\frac{5}{72x^2}$	$-\frac{864x^3}{5}$	$\frac{20736x^4}{15}$	$-\frac{1}{248832x^5}$	0	1	$\frac{5}{6}$	$\frac{2}{3}$	$\frac{1}{2}$	$\frac{1}{3}$	$-\frac{1}{6}$
$-\frac{3}{7x}$	$\frac{15}{196x^2}$	$-\frac{5}{686x^3}$	$\frac{38416x^4}{15}$	$-\frac{1}{268912x^5}$	$\frac{1}{7529537x^6}$	1	$\frac{6}{7}$	$\frac{5}{7}$	$\frac{4}{7}$	$\frac{3}{7}$	$\frac{2}{7}$

Gaussian fixed point

pure gravity fixed point

multicritical points

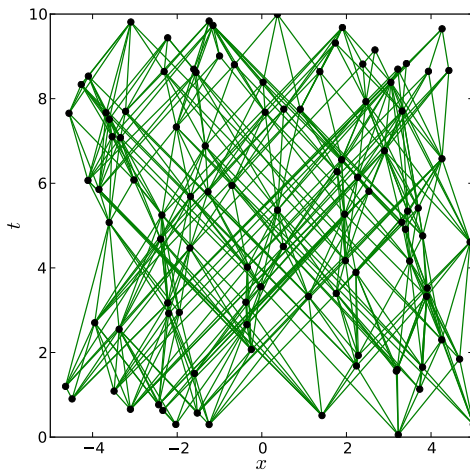
exact result: $\theta = 2/(m + 1/2)$, here: $\theta = 2/m$

Multicritical points recovered with Functional RG

Quantum spacetime is ...

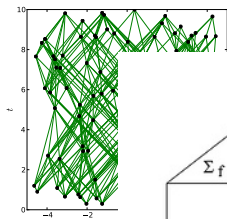
Quantum spacetime is ...

... a causal set

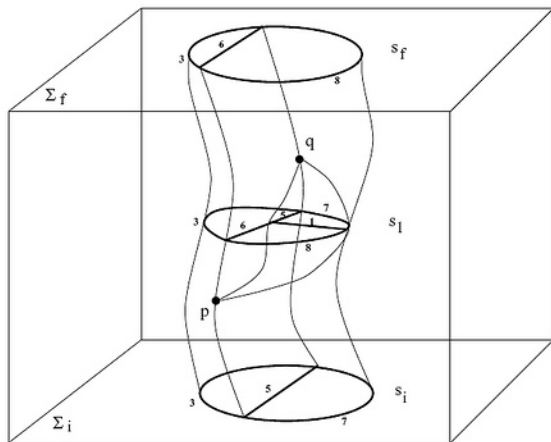


Quantum spacetime is ...

... a causal set

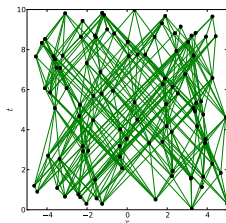


... a spinfoam

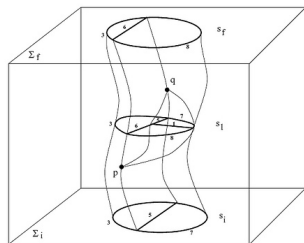


Quantum spacetime is ...

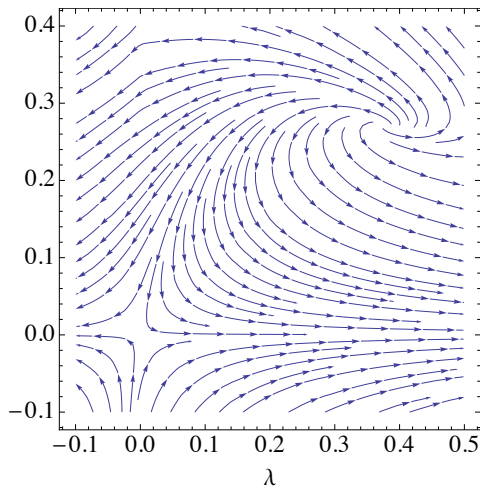
... a causal set



... a spinfoam

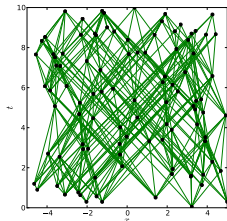


... a QFT described by an
RG fixed point



Quantum spacetime is ...

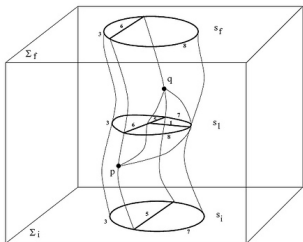
... a causal set



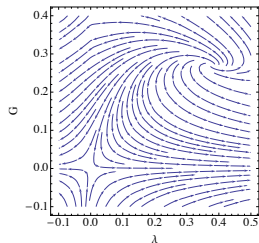
... noncommutative

$$[x_1, x_2] \sim \theta$$

... a spinfoam

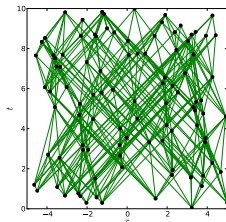


... q QFT described by an
RG fixed point

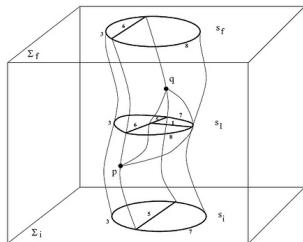


Quantum spacetime is ...

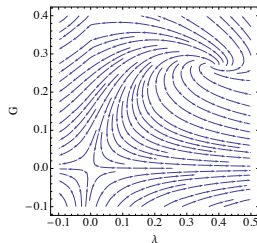
... a causal set



... a spinfoam



... described by an RG
fixed point

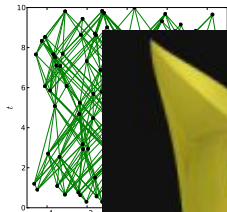


... noncommutative
 $[x_1, x_2] \sim \theta$

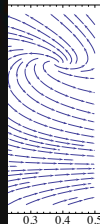
Quantum spacetime is ...

... a causal set

...stringy and braney

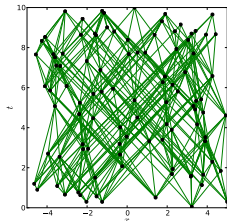


by an RG
point



Quantum spacetime is ...

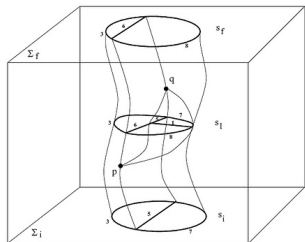
... a causal set



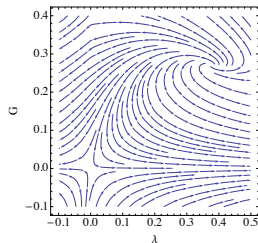
...stringy and braney



... a spinfoam



... a QFT described by an RG fixed point

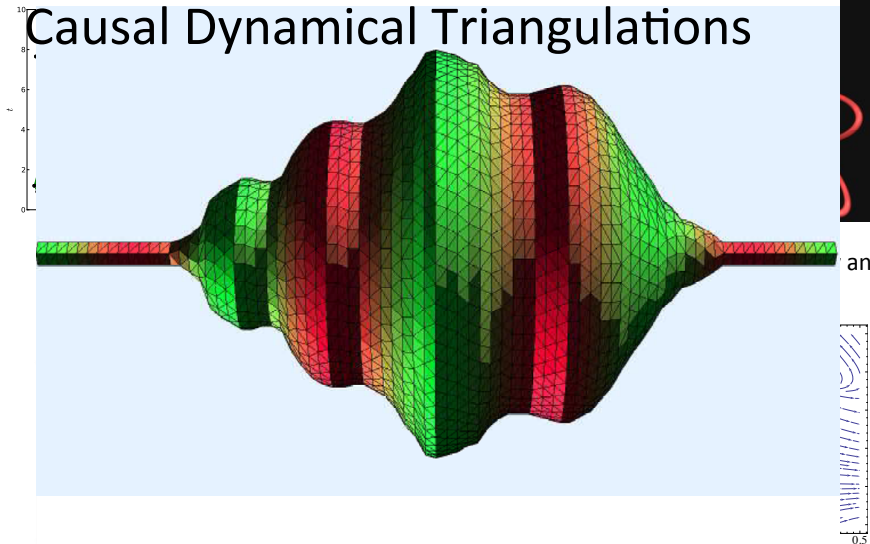


... noncommutative
 $[x_1, x_2] \sim \theta$

Quantum spacetime is ...

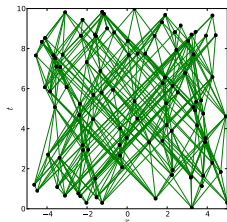
... the continuum limit of

Causal Dynamical Triangulations

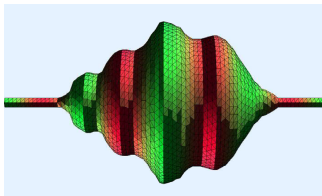


Quantum spacetime is ...

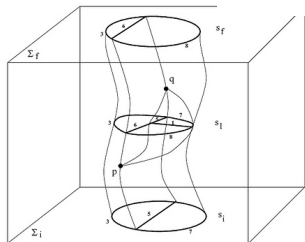
... a causal set



... the continuum limit of
Causal Dynamical Triangulations



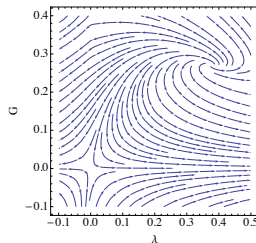
... a spinfoam



...stringy and braney



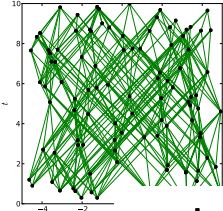
... a QFT described by an
RG fixed point



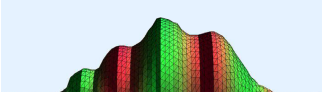
... noncommutative
 $[x_1, x_2] \sim \theta$

Quantum spacetime is ...

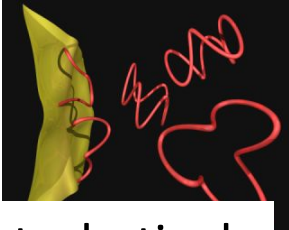
... a causal set



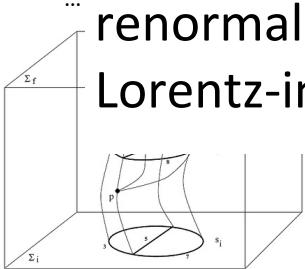
... the continuum limit of Causal Dynamical Triangulations



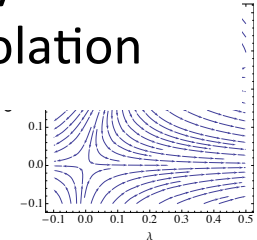
...stringy and braney



... described by a perturbatively renormalizable theory with Lorentz-invariance violation

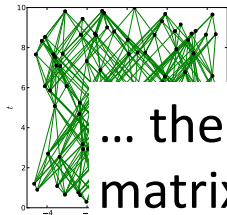


... noncommutative
 $[x_1, x_2] \sim \theta$



Quantum spacetime is ...

... a causal set



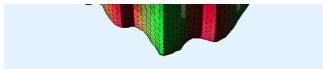
... the continuum limit of
Causal Dynamical Triangulations

...stringy and braney

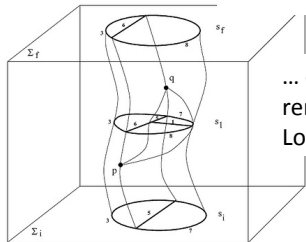


... the continuum limit of a matrix/ tensor model

... a spinfoam

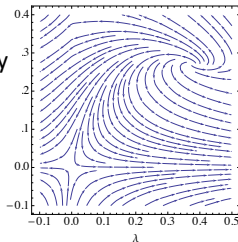


... a QFT described by an
RG fixed point



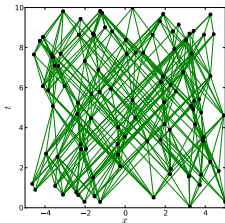
... described by a perturbatively
renormalizable theory with
Lorentz-invariance violation

... noncommutative
 $[x_1, x_2] \sim \theta$



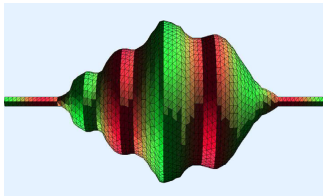
Quantum spacetime is ...

... a causal set

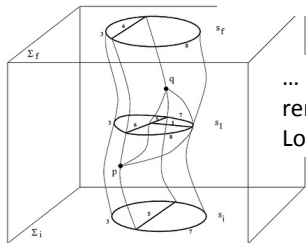


... the continuum limit of a matrix/ tensor model

... the continuum limit of Causal Dynamical Triangulations



... a spinfoam



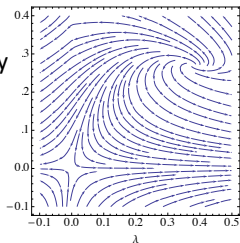
... described by a perturbatively renormalizable theory with Lorentz-invariance violation

... noncommutative
 $[x_1, x_2] \sim \theta$

...stringy and braney



... a QFT described by an RG fixed point



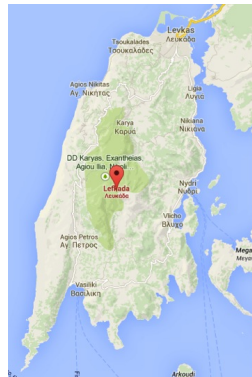
Which way to quantum gravity?



Which way to quantum gravity?



Which way to quantum gravity?



Are any of the quantum gravity models related?

Two sides of the same picture?

Two sides of the same picture?

Discrete approaches:

pregeometric phase

phase transition

to continuum spacetime

Asymptotic safety:

geometric phase: $\langle g_{\mu\nu} \rangle$ exists

interacting fixed point

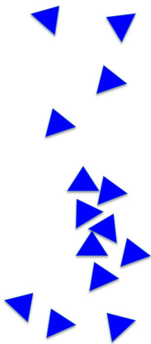
in the ultraviolet

Two sides of the same picture?

Discrete approaches:

pregeometric phase

phase transition
to continuum spacetime



Asymptotic safety:

geometric phase: $\langle g_{\mu\nu} \rangle$ exists

interacting fixed point
in the ultraviolet

Two sides of the same picture?

Discrete approaches:

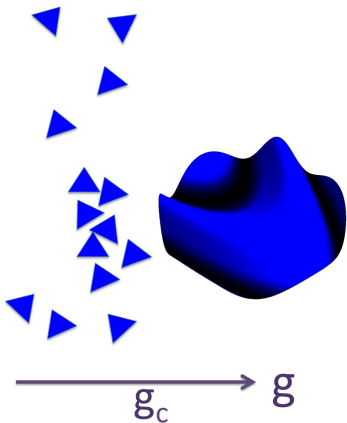
pregeometric phase

phase transition
to continuum spacetime

Asymptotic safety:

geometric phase: $\langle g_{\mu\nu} \rangle$ exists

interacting fixed point
in the ultraviolet



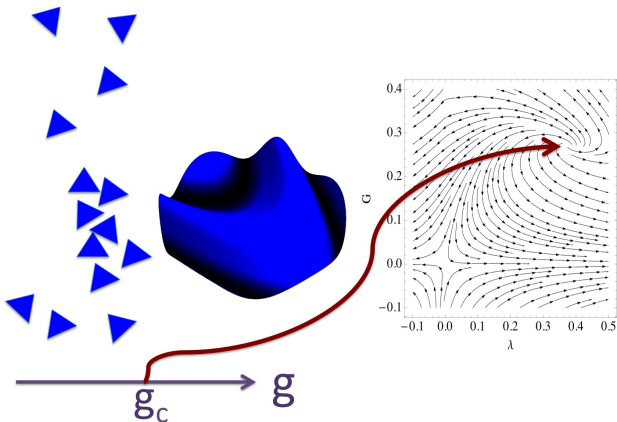
Two sides of the same picture?

Discrete approaches:
pregeometric phase

phase transition
to continuum spacetime

Asymptotic safety:
geometric phase: $\langle g_{\mu\nu} \rangle$ exists

interacting fixed point
in the ultraviolet



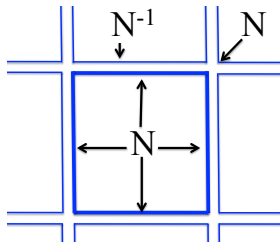
Double-scaling limit and continuum beta function

Double-scaling limit and continuum beta function

matrix model:

configurations scale with

$$N^{F-E+V}$$
$$= N^{\chi}$$

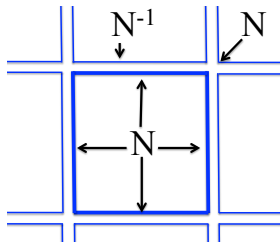


Double-scaling limit and continuum beta function

matrix model:

configurations scale with

$$N^{F-E+V}$$
$$= N^{\chi}$$



spacetime description:

$$e^{-S} = e^{-\Lambda A + \frac{1}{4G_0} \chi}$$

configurations scale with

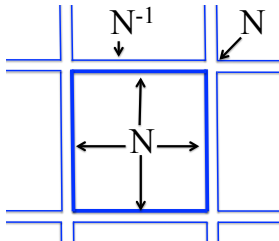
$$e^{\frac{1}{4G_0} \chi}$$

Double-scaling limit and continuum beta function

matrix model:

configurations scale with

N^χ



spacetime description:

$$e^{-S} = e^{-\Lambda A + \frac{1}{4G_0} \chi}$$

configurations scale with

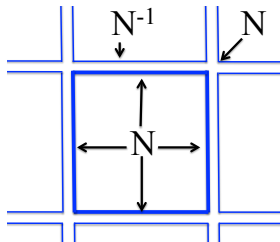
$$e^{\frac{1}{4G_0} \chi}$$

Double-scaling limit and continuum beta function

matrix model:

configurations scale with

N^χ



g_4^n (n: number of building blocks)

spacetime description:

$$e^{-S} = e^{-\Lambda A + \frac{1}{4G_0} \chi}$$

configurations scale with

$$e^{\frac{1}{4G_0} \chi}$$

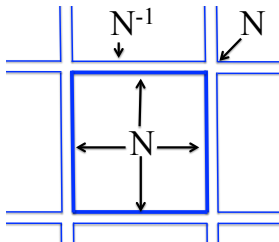
$$e^{-\Lambda A} = e^{-\Lambda a^2 n}$$

Double-scaling limit and continuum beta function

matrix model:

configurations scale with

N^χ



spacetime description:

$$e^{-S} = e^{-\Lambda A + \frac{1}{4G_0} \chi}$$

configurations scale with

$$e^{\frac{1}{4G_0} \chi}$$

g_4^n (n : number of building blocks)

$$e^{-\Lambda A} = e^{-\Lambda a^2 n}$$

$$(g_4 - g_{4c})^{5/4} N = \text{const} \Rightarrow (\Lambda_R a^2)^{5/4} e^{\frac{1}{4G_0} \chi} = \text{const} \rightarrow G_0(a) \text{ [J. Ambjorn]}$$

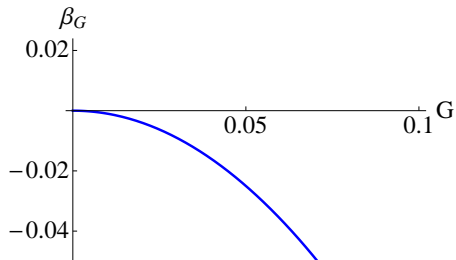
$$\Rightarrow \beta_G = -a \partial_a G = -10G^2$$

Double-scaling limit and continuum beta function

$$\beta_G = -10G^2$$

Double-scaling limit and continuum beta function

$$\beta_G = -10G^2$$



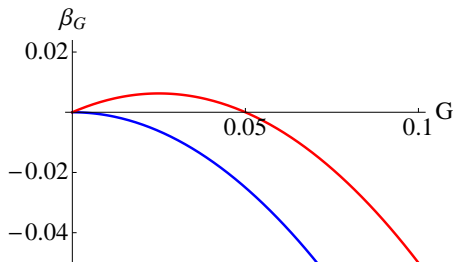
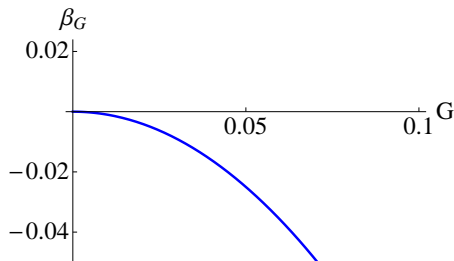
Double-scaling limit and continuum beta function

$$\beta_G = -10G^2$$

$$d = 2 + \epsilon$$

[Christensen, Duff, 1978; Gastmanns et al., 1978; Weinberg, 1979]

$$\beta_G = \epsilon G - cG^2 \quad (c > 0)$$

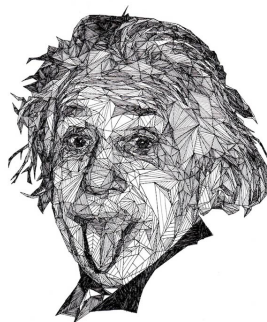


interacting fixed point (asymptotic safety) \leftrightarrow double-scaling limit

Summary

Matrix/tensor models for quantum gravity:
Sum over spacetimes \leftrightarrow Matrix/tensor path
integral

Continuum limit:
double scaling limit \leftrightarrow RG fixed point



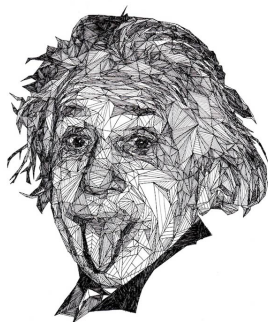
Summary

Matrix/tensor models for quantum gravity:
Sum over spacetimes \leftrightarrow Matrix/tensor path
integral

Continuum limit:
double scaling limit \leftrightarrow RG fixed point

new FRG tools: double-scaling limit (pure gravity and with matter) in
matrix models

connection between double-scaling limit and $d = 2 + \epsilon$ continuum beta
function showing asymptotic safety



Outlook:

Towards $d = 4$, quantitative precision, and new connections

- towards quantitative precision: control symmetry-breaking sector
- towards $d = 4$: similar strategy as in $d = 2$, larger theory space
ongoing work in $d = 3$ tensor models [Ben Geloun, Benedetti, Oriti]
- relation to continuum approaches/ asymptotic safety: compare critical exponents in $d = 4$

Outlook:

Towards $d = 4$, quantitative precision, and new connections

- towards quantitative precision: control symmetry-breaking sector
- towards $d = 4$: similar strategy as in $d = 2$, larger theory space
ongoing work in $d = 3$ tensor models [Ben Geloun, Benedetti, Oriti]
- relation to continuum approaches/ asymptotic safety: compare critical exponents in $d = 4$

Thank you for your attention!